

Note: Minimum Uncertainty States (MUS)

- ☒ **Heisenberg's Uncertainty Relation**
- ☒ **Minimum Uncertainty States (MUS)**
- ☒ **Gaussian States**
- ☒ **Free-particle expansion**
- ☐ **Squeezed States**
- **More on Uncertainty Relations**
 - ☐ Intelligent States
 - ☐ Robertson–Schrodinger uncertainty relations
 - ☐ Quantum entropic uncertainty principle
 - ☐ **Quantum Metrology**
 - ☐ **Heisenberg limit** → **Quantum Cramer-Rao bound**
 - ☐ Quantum Non-Demolition (QND) Measurement

Note: Phase Space

- ☒ **Quantum Distribution Theory**
- ☒ **Ordering**
- ☐ **Normally Ordering, P-representation** (Glauber-Sudarshan)
- ☐ **Anti-Normally Ordering, Q-representation** (Husimi)
- ☒ **Symmetrically Ordering, W-representation** (Wigner-Weyl)
- **More on Phase Space**
 - ☐ Fokker-Planck equation
 - ☐ Quantum Liouville equation
 - ☐ Moyal function
 - ☐ Homodyne detection
 - ☐ Wigner Flow

Note: Squeezed States

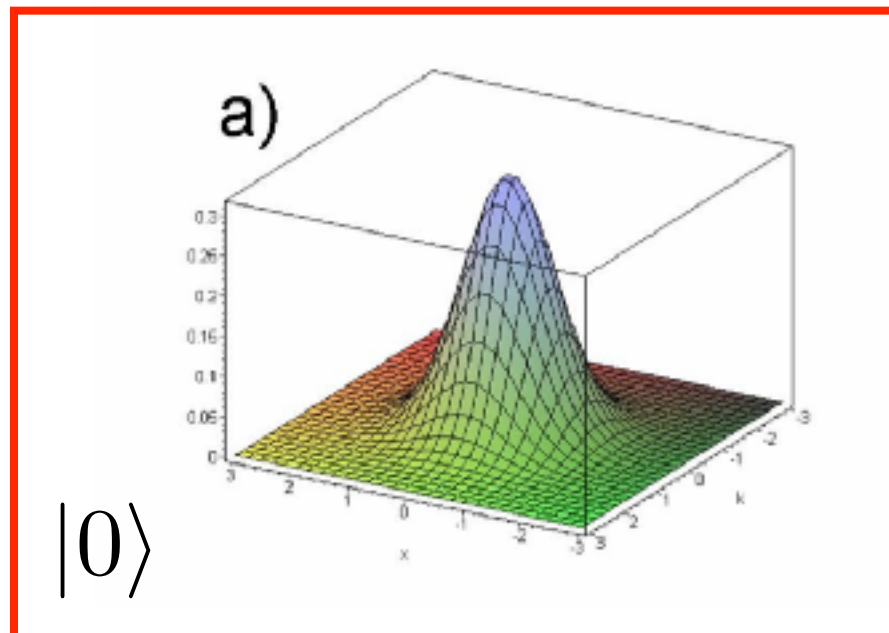
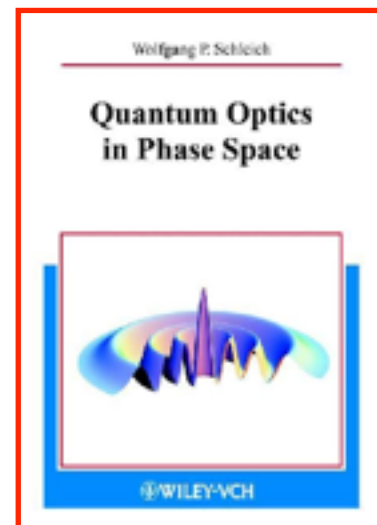
- ☐ Squeezed Operator
- ☐ Squeezed Vacuum
- ☐ Squeezed Coherent v.s. Coherent Squeezed states
- ☐ Squeezed states as the Minimum Uncertainty states
- ☐ Generation of Squeezed states by OPO
- ☐ Homodyne detections
- More on Squeezed States
 - ☐ OPA
 - ☐ FWM Squeezing
 - ☐ Soliton Squeezing
 - ☐ Squeezing spectrum
 - ☐ Applications to Gravitational Wave Detectors

Number (Fock) states

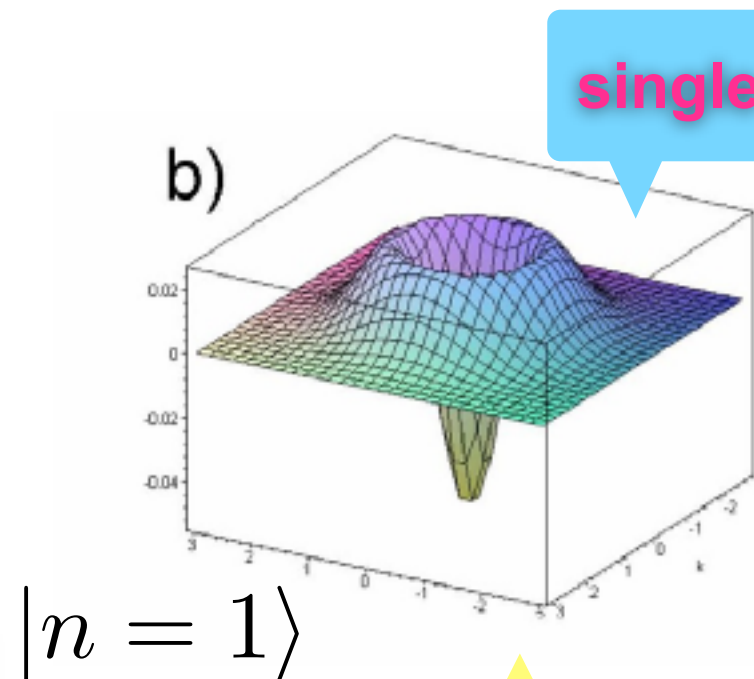
non-classical
states

- Wigner quasiprobability distribution

$$W(x, p) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\xi \psi^* \left(x - \frac{\xi}{2}\right) \psi \left(x + \frac{\xi}{2}\right) e^{-ik\xi}$$



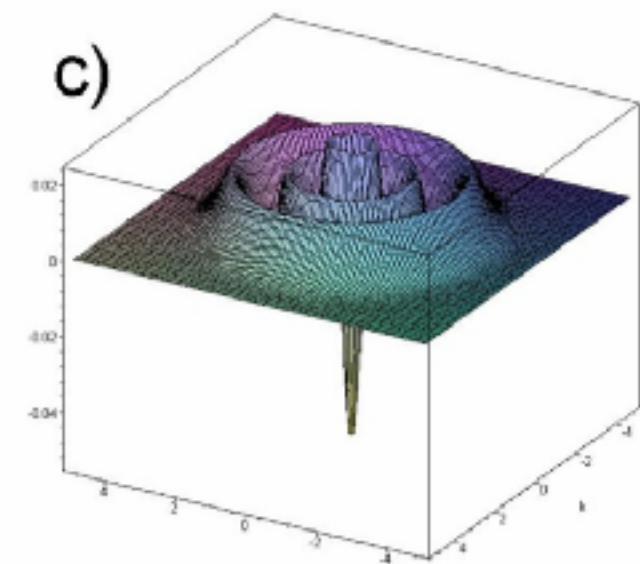
$|0\rangle$



single-photon

$|n = 1\rangle$

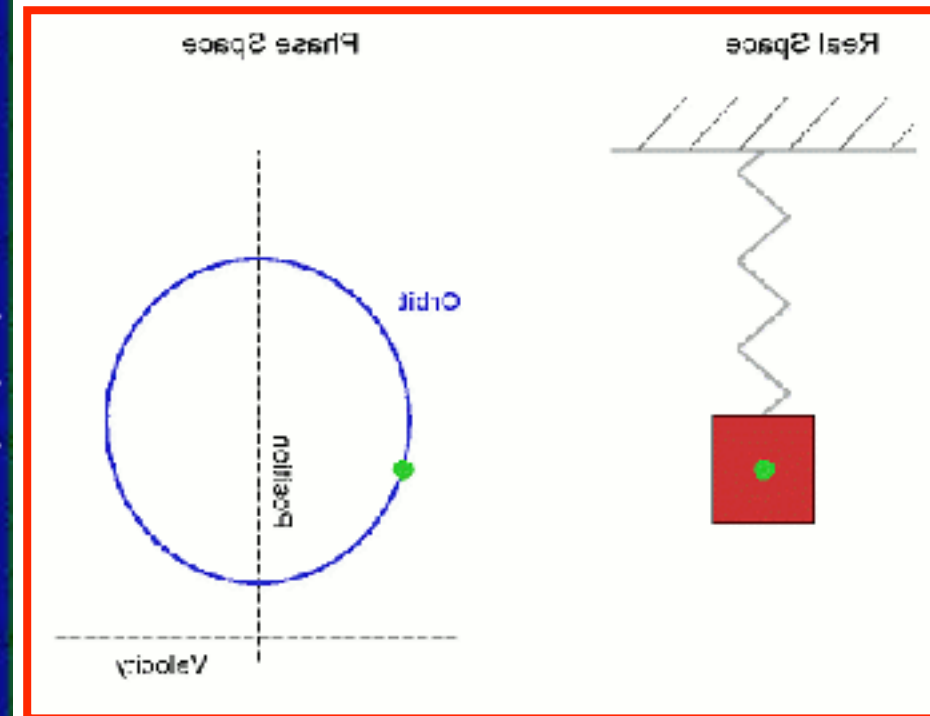
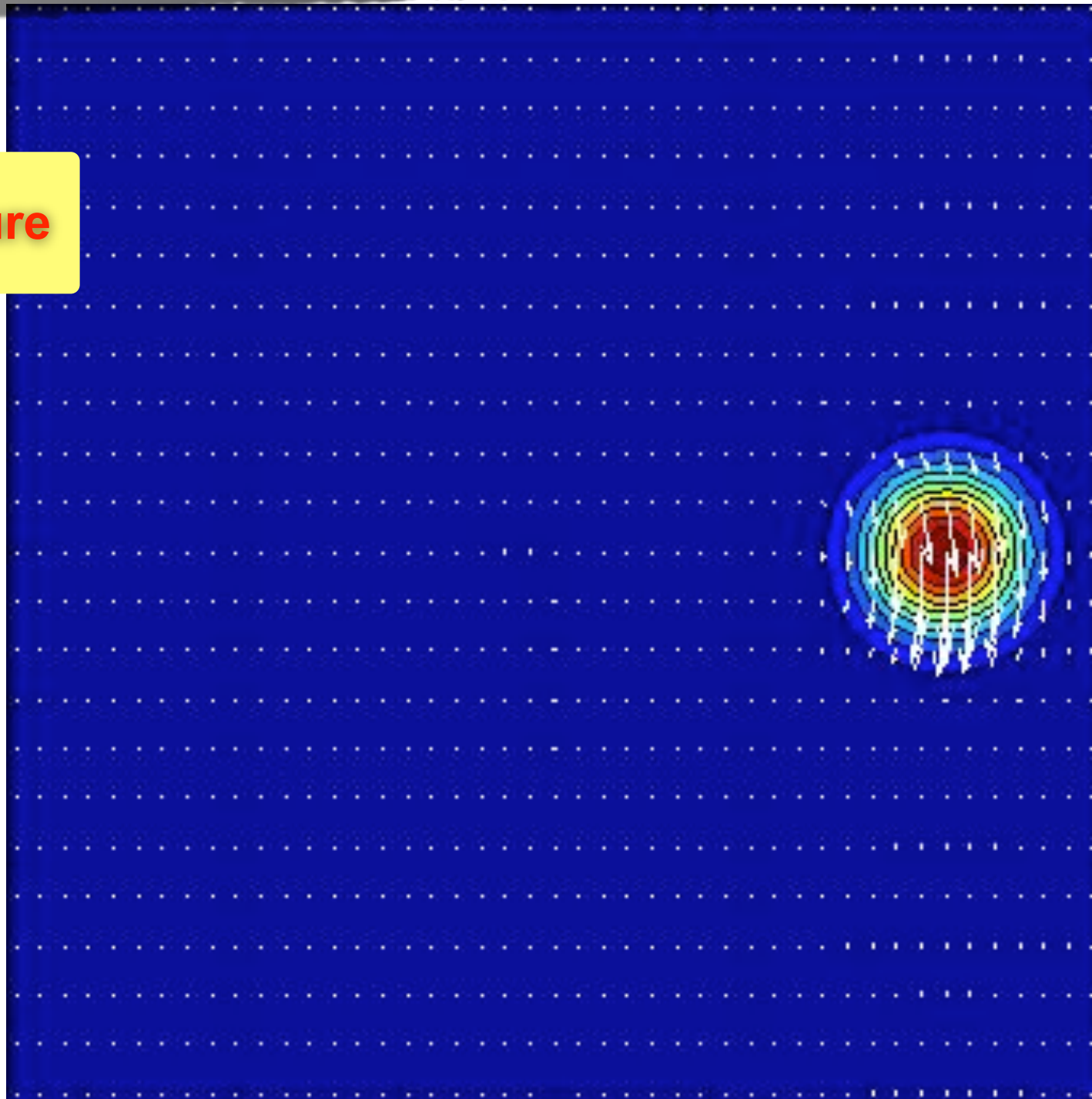
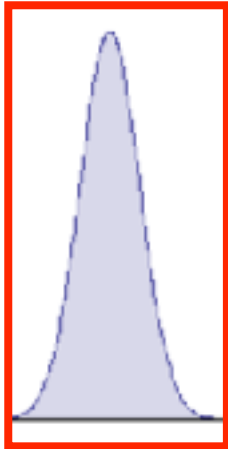
negative probability



with Ludmila Praxmeyer

Coherent states

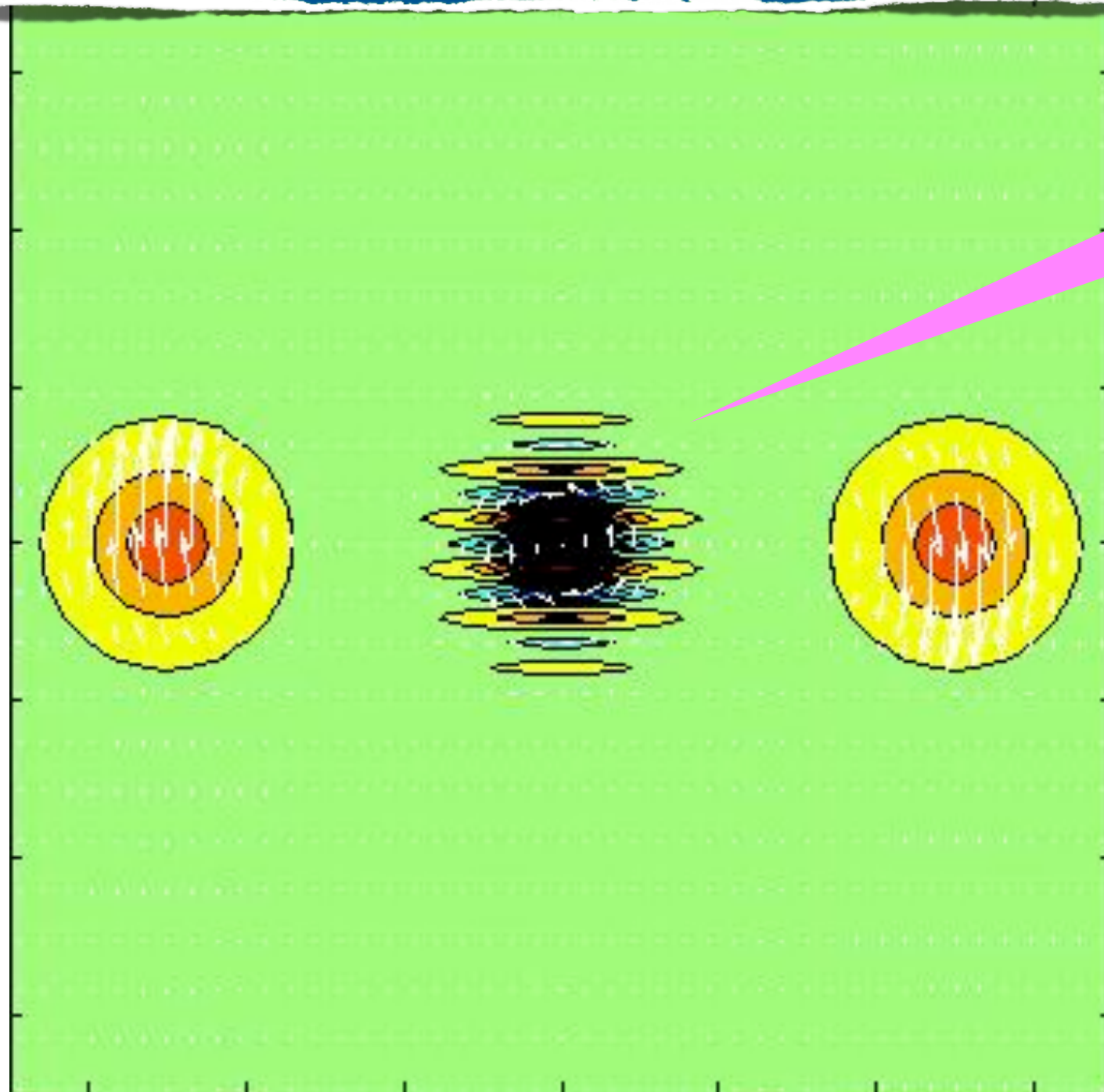
wave-nature



with Popo Yang

Popo Yang, Ivan F. Valtierra, Andrei B. Klimov, Shin-Tza Wu, RKL, Luis L. Sanchez-Soto, and Gerd Leuchs, Physica Scripta for the New Focus issue: [Quantum Optics and Beyond - in honour of Wolfgang Schleich](#).

Cat States in Phase Space



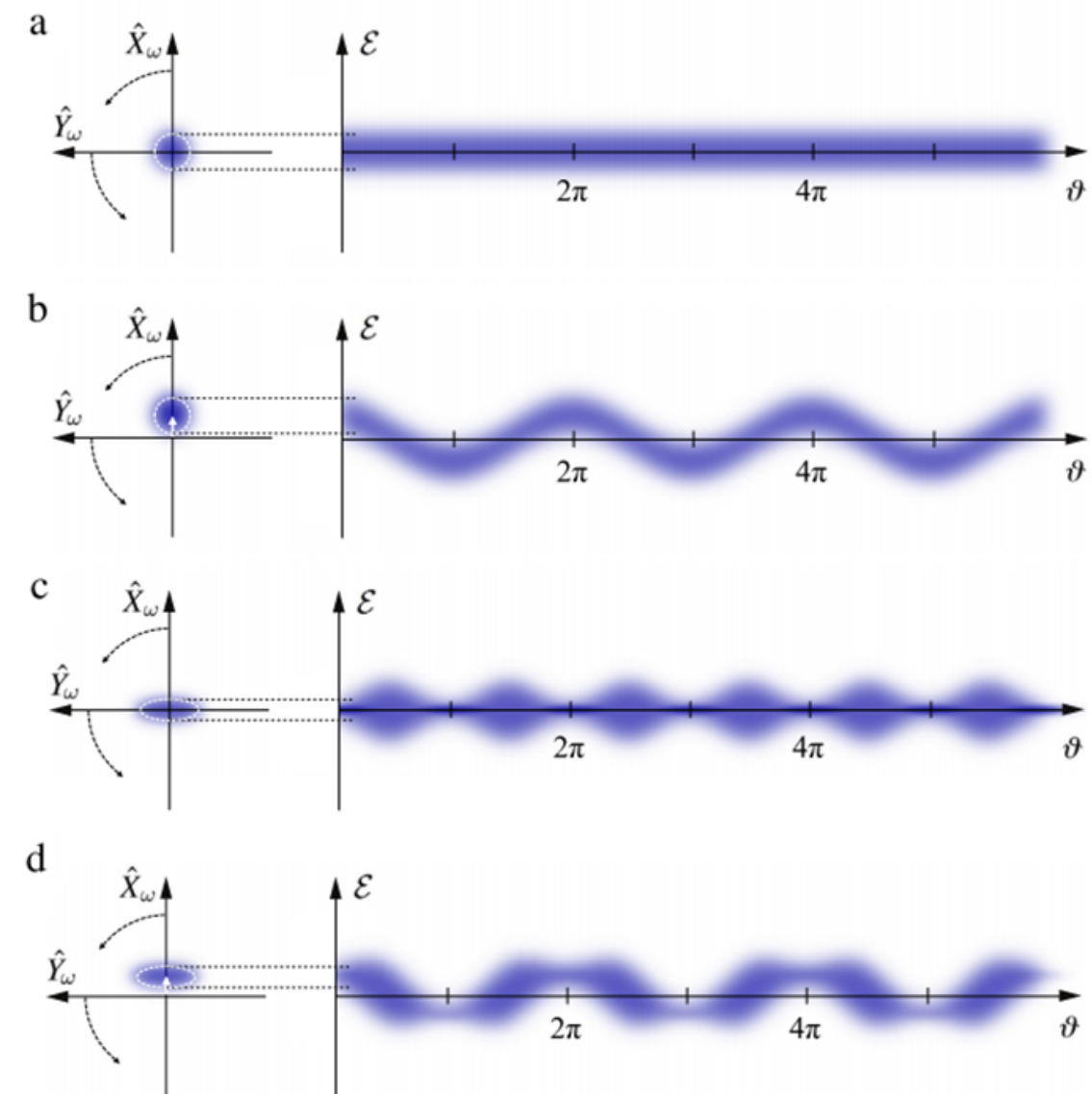
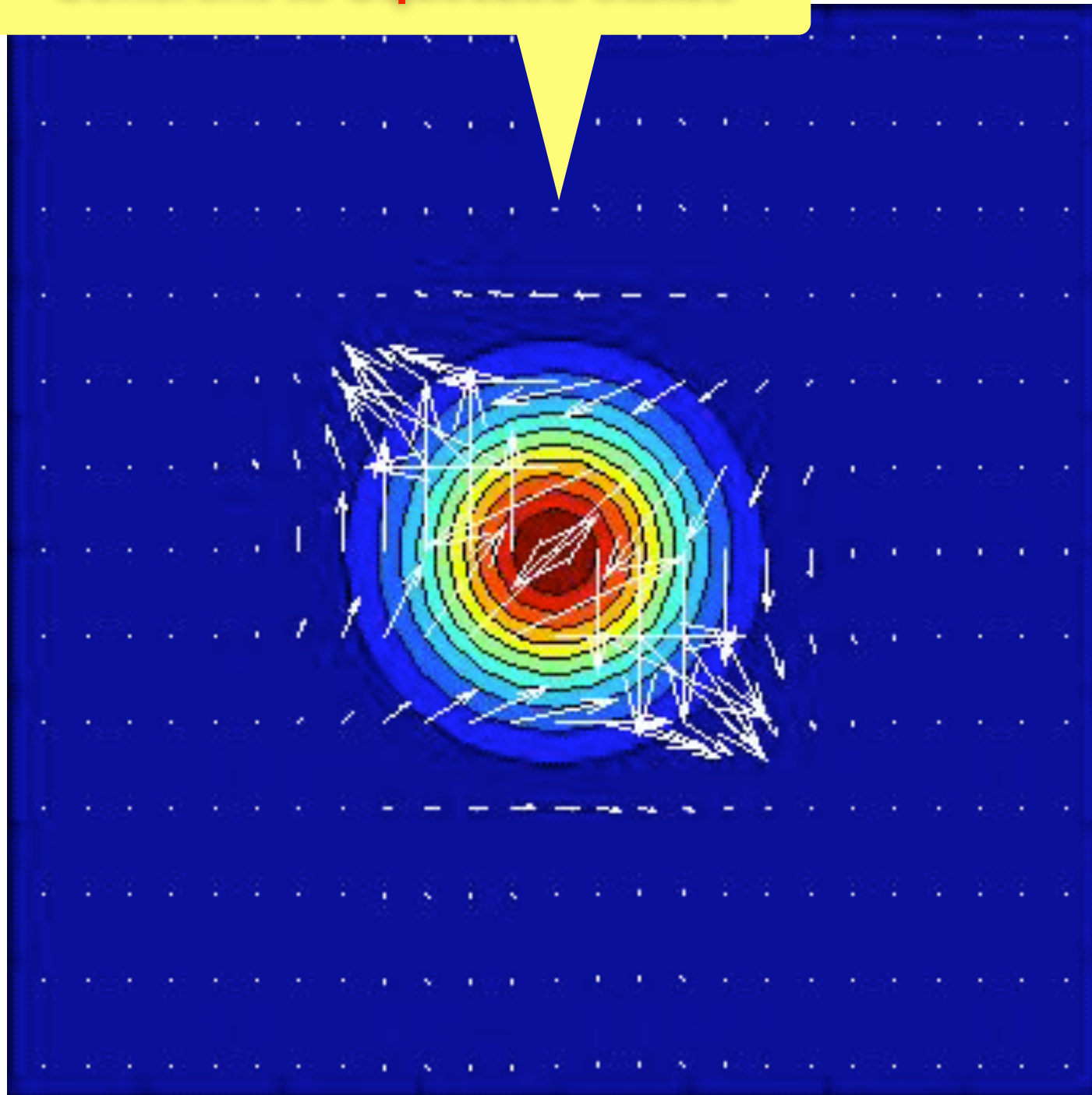
Non-classicality

with Popo Yang

Popo Yang, Ivan F. Valtierra, Andrei B. Klimov, Shin-Tza Wu, RKL, Luis L. Sanchez-Soto, and Gerd Leuchs, Physica Scripta for the New Focus issue: **Quantum Optics and Beyond - in honour of Wolfgang Schleich.**

Squeezed States

Coherent to Squeezed states

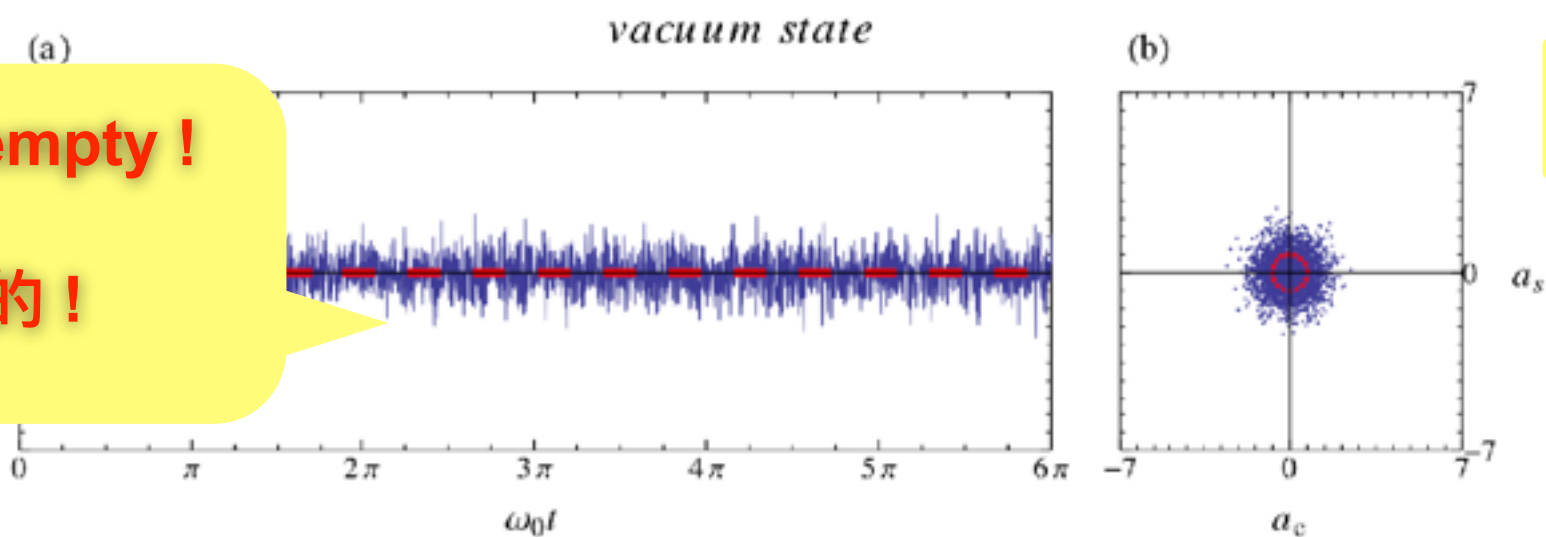


Courtesy:
Roman Schnabel (2017).

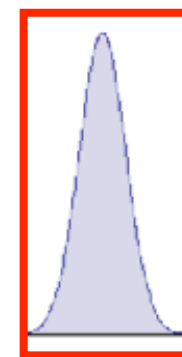
by Popo Yang

Vacuum is NOT empty !

真空並不是空的！



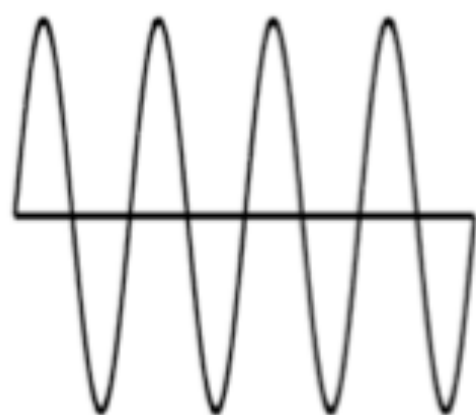
wave-nature



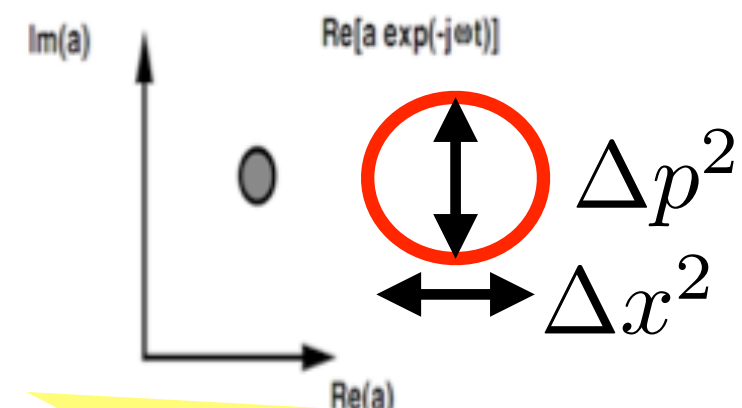
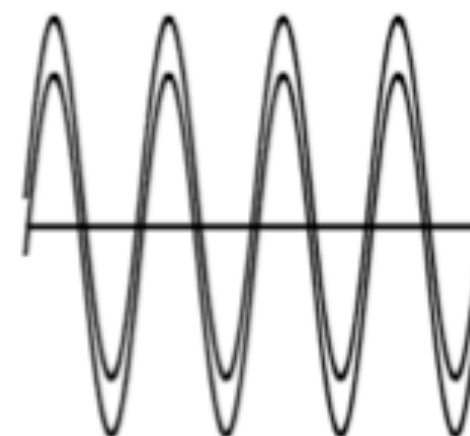
● coherent state
 $(\Delta \hat{X}_1)^2 = \frac{1}{4}, (\Delta \hat{X}_2)^2 = \frac{1}{4}, \Delta \hat{X}_1 \Delta \hat{X}_2 = \frac{1}{4}$

● squeezed state
 $(\Delta \hat{Y}_1)^2 = \frac{1}{4} e^{-2r}, (\Delta \hat{Y}_2)^2 = \frac{1}{4} e^{2r}, \Delta \hat{Y}_1 \Delta \hat{Y}_2 = \frac{1}{4}$

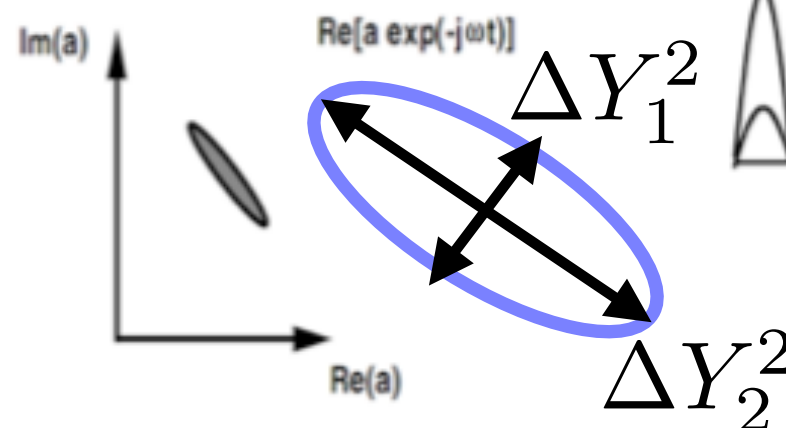
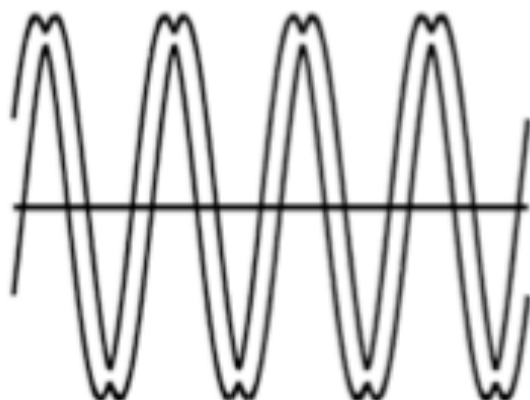
● classical field



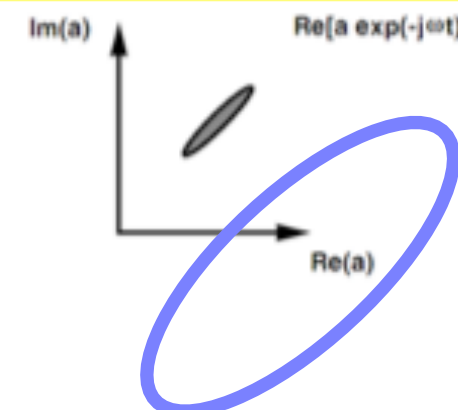
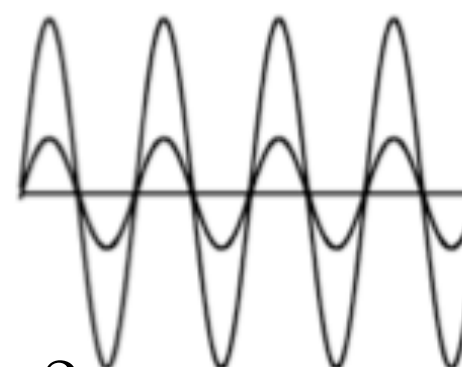
● coherent field



● amplitude-squeezed field

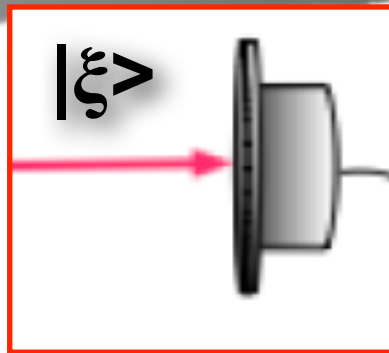


● phase-squeezed field



Non-classical states

Squeezed Vacuum State: $|\xi\rangle$



time-sequence

histogram

Gaussian wave-package

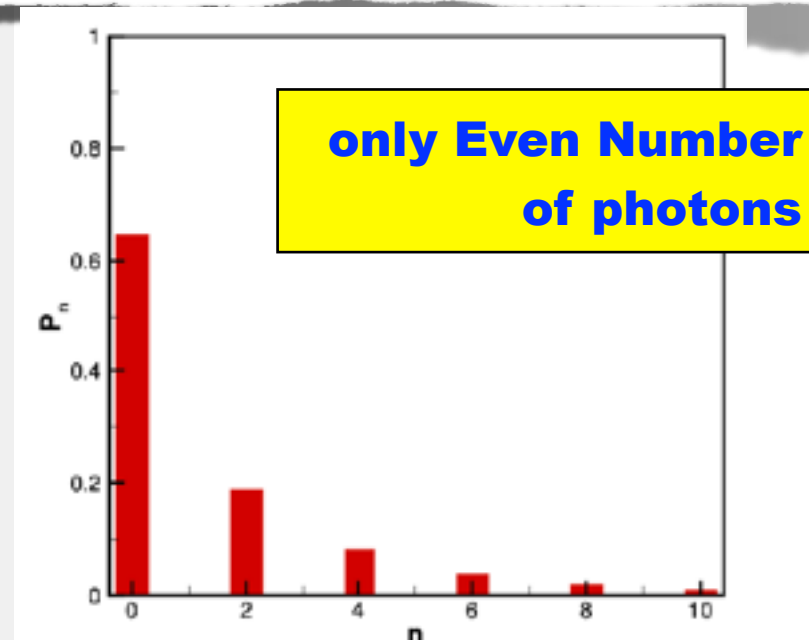
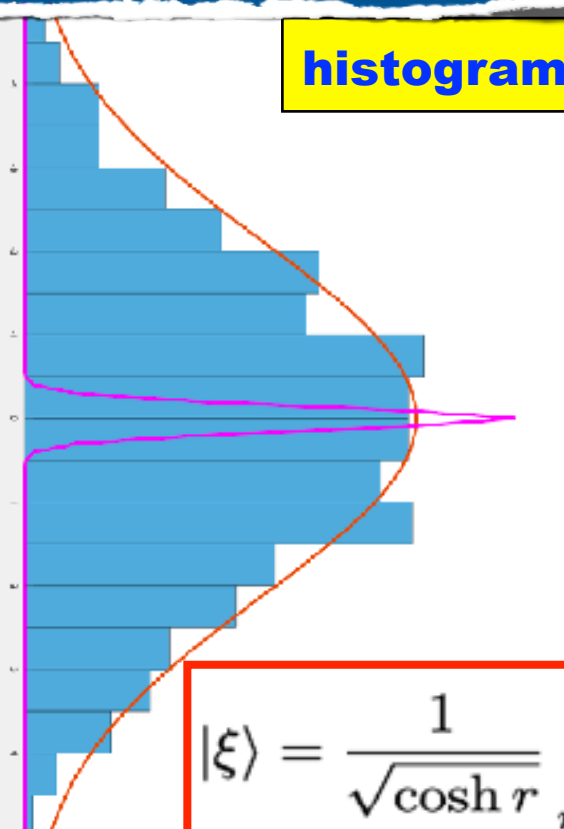
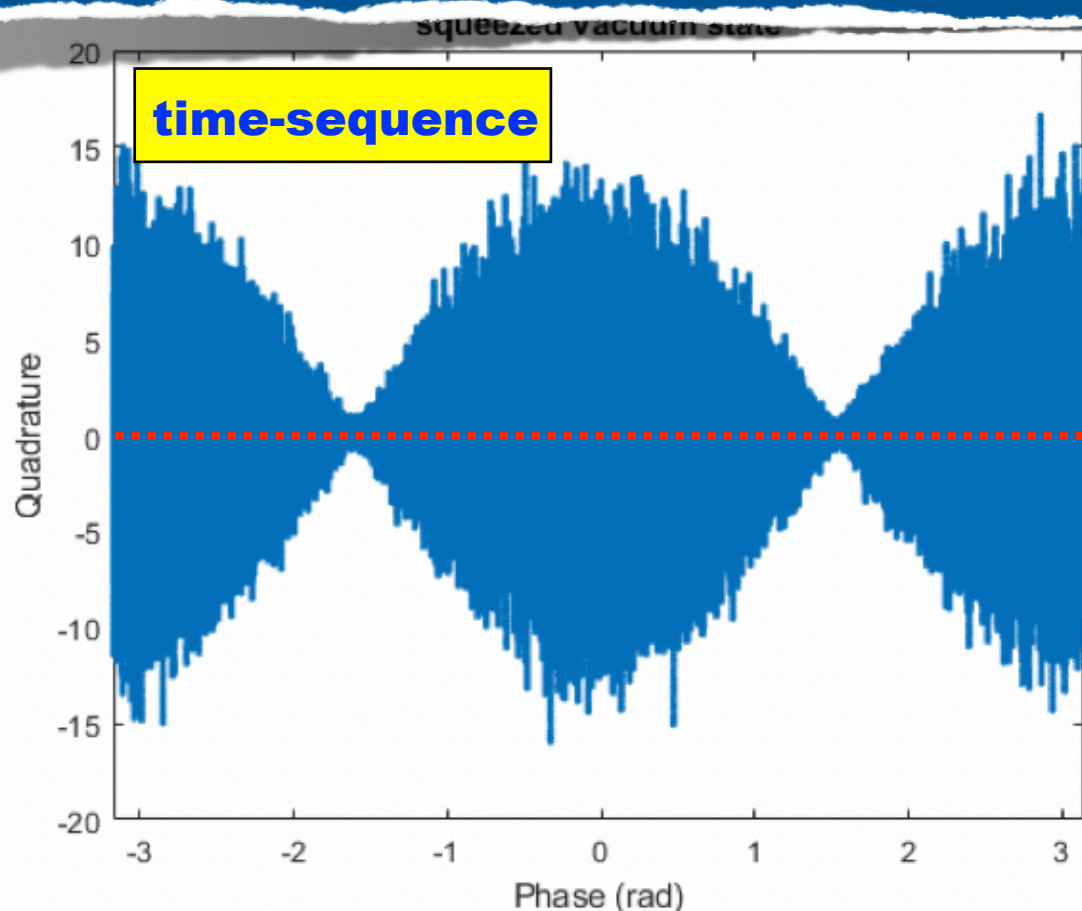
c
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-20 -3 -2 -1 0 1 2 3

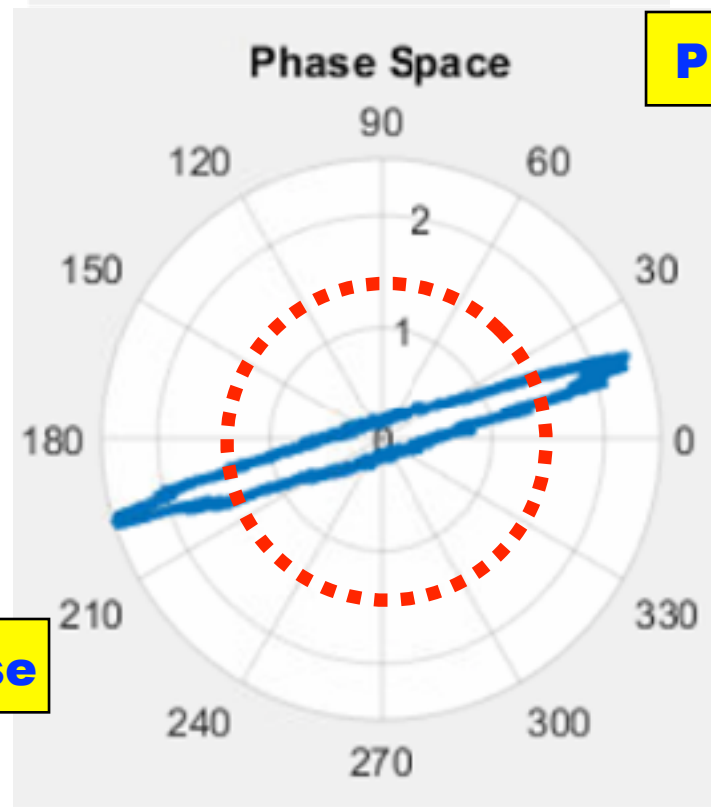
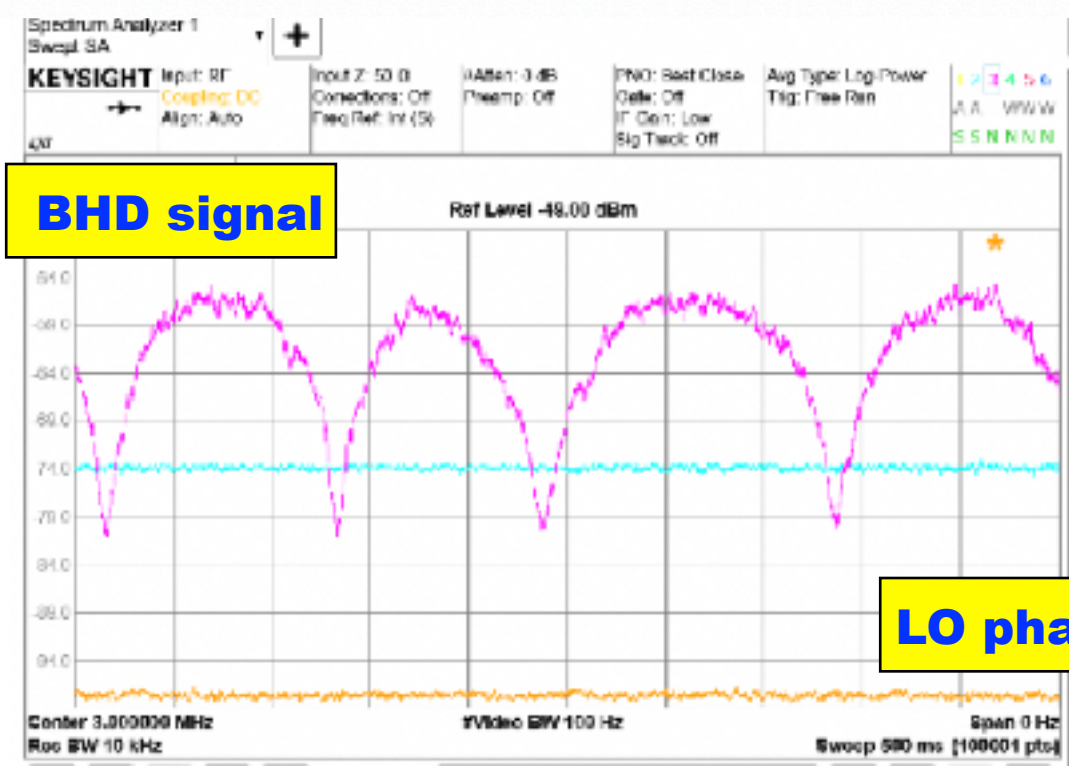
time

with Zero Mean

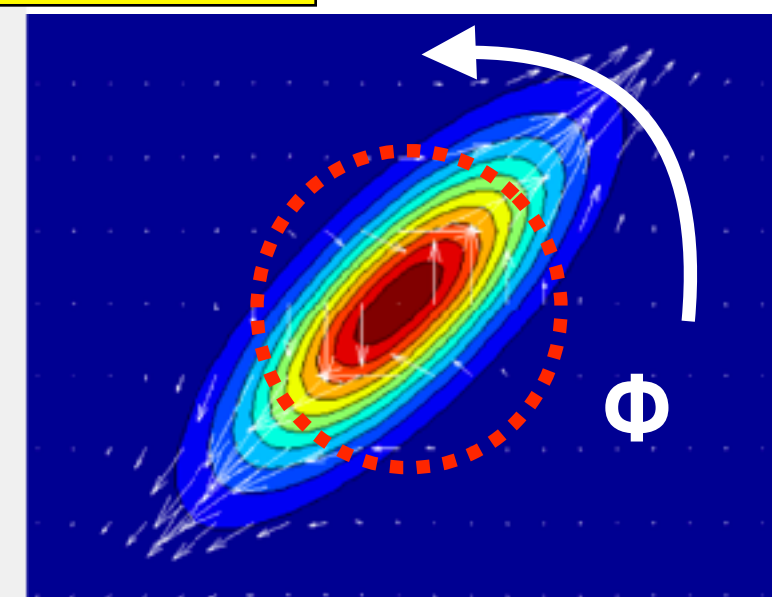
Squeezed States



$$|\xi\rangle = \frac{1}{\sqrt{\cosh r}} \sum_{m=0}^{\infty} (-1)^m \frac{\sqrt{(2m)!}}{2^m m!} e^{im\theta} \tanh^m r |2m\rangle$$



Phase Sapce



Squeezed States and SHO:

Optical Parametric Oscillator, OPO

For the degenerate parametric process, *i.e.*, two-photon process, its Hamiltonian is

$$\hat{H} = i\hbar(g\hat{a}^{\dagger 2} - g^*\hat{a}^2),$$

Quantum Correlation

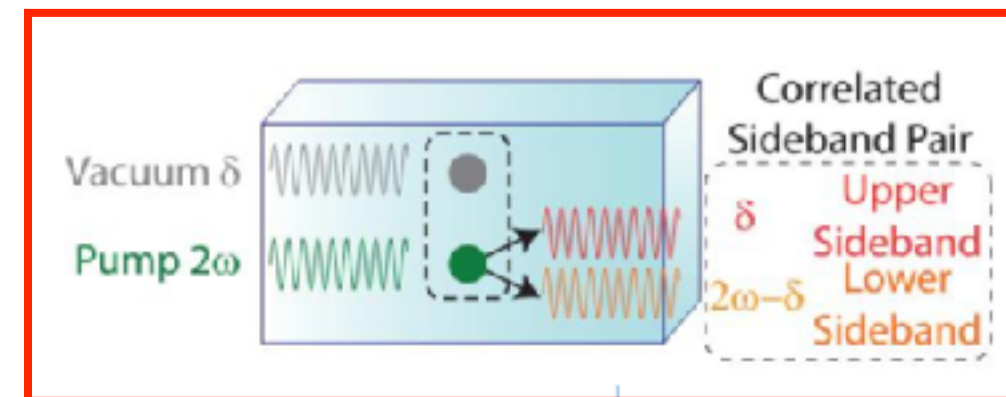
where g is a coupling constant. The state of fields generated by this Hamiltonian is

$$|\Psi(t)\rangle = \exp[(g\hat{a}^{\dagger 2} - g^*\hat{a}^2)t]|0\rangle.$$

Then, one can define the unitary squeeze operator

$$\hat{S}(\xi) = \exp\left[\frac{1}{2}\xi^*\hat{a}^2 - \frac{1}{2}\xi\hat{a}^{\dagger 2}\right],$$

where $\xi = r\exp(i\theta)$ is an arbitrary complex number.



Squeezed Operator

$$\hat{S}^\dagger(\xi) = \hat{S}^{-1}(\xi) = \hat{S}(-\xi),$$

Squeezed Operator

- If $|\Psi\rangle$ is the vacuum state $|0\rangle$, then $|\Psi_s\rangle$ state is the *squeezed vacuum*,

$$|\xi\rangle = \hat{S}(\xi)|0\rangle.$$

- The variances for squeezed vacuum are

$$\Delta\hat{a}_1^2 = \frac{1}{4}[\cosh^2 r + \sinh^2 r - 2 \sinh r \cosh r \cos \theta],$$

$$\Delta\hat{a}_2^2 = \frac{1}{4}[\cosh^2 r + \sinh^2 r + 2 \sinh r \cosh r \cos \theta],$$

- For $\theta = 0$, we have

$$\Delta\hat{a}_1^2 = \frac{1}{4}e^{-2r}, \quad \text{and} \quad \Delta\hat{a}_2^2 = \frac{1}{4}e^{+2r},$$

and squeezing exists in the $\hat{a}_1$ quadrature.

- For $\theta = \pi$, the squeezing will appear in the $\hat{a}_2$ quadrature.

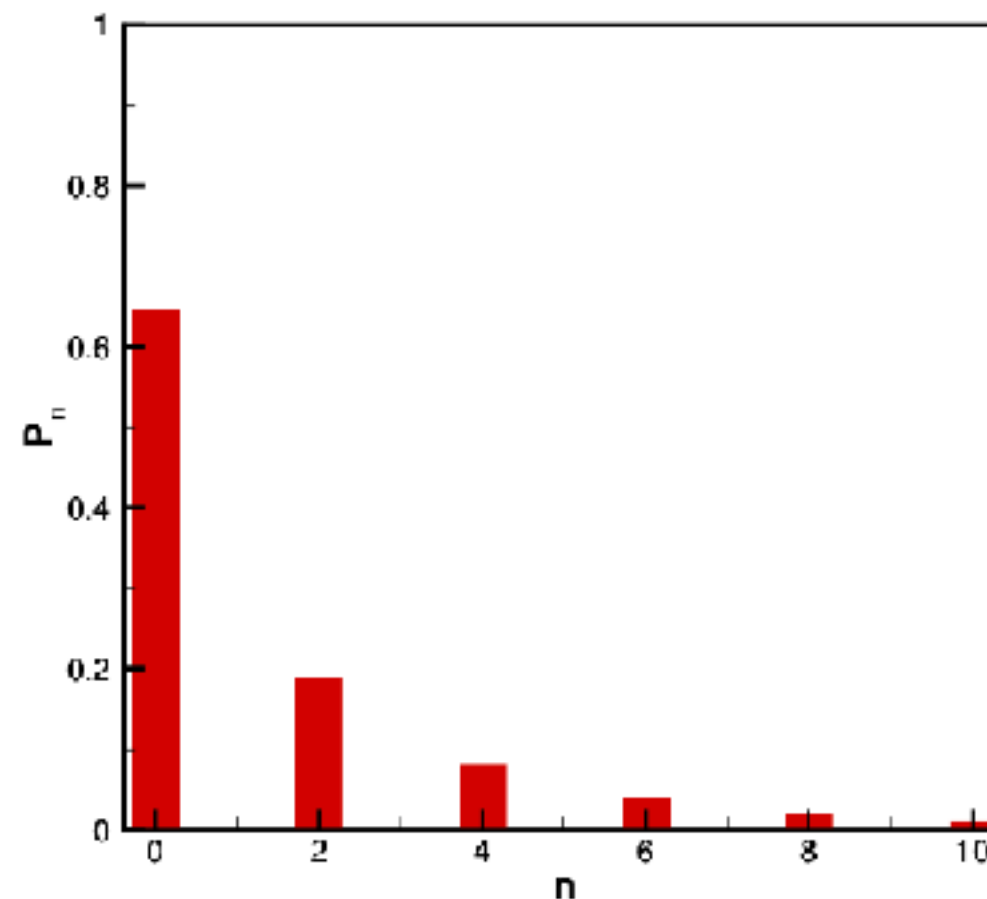


Squeezed State as a Eigen-state of UR

$$(\mu \hat{a} + \nu \hat{a}^\dagger) |\xi\rangle = 0,$$

Squeezed states in the Number state basis

$$|\xi\rangle = \frac{1}{\sqrt{\cosh r}} \sum_{m=0}^{\infty} (-1)^m \frac{\sqrt{(2m)!}}{2^m m!} e^{im\theta} \tanh^m r |2m\rangle$$



Sub-Poisson distribution

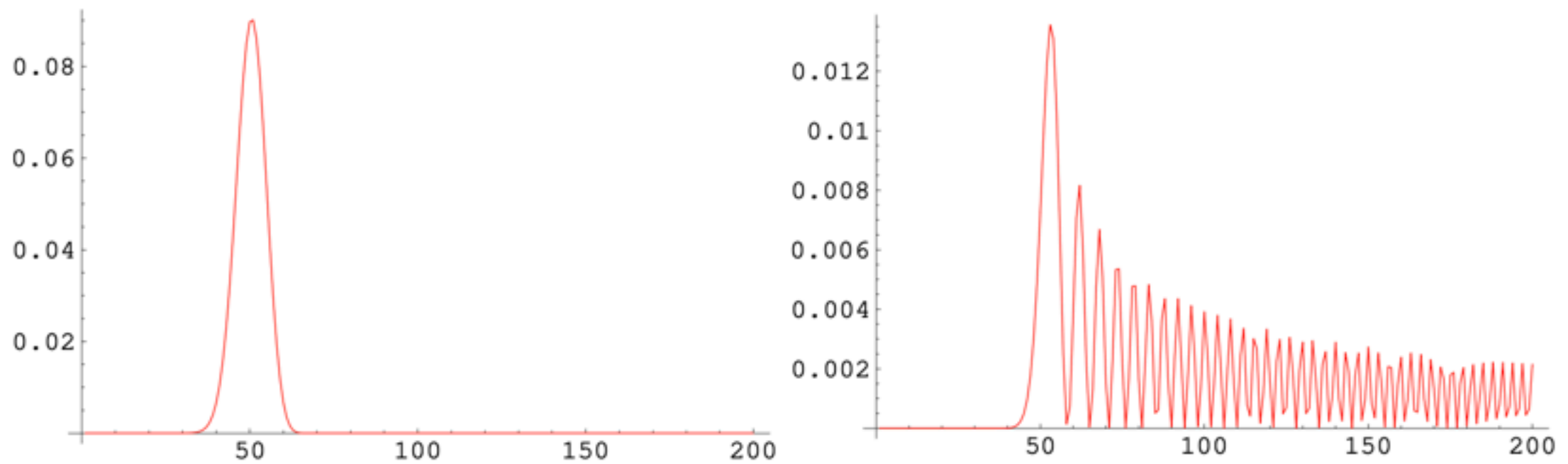
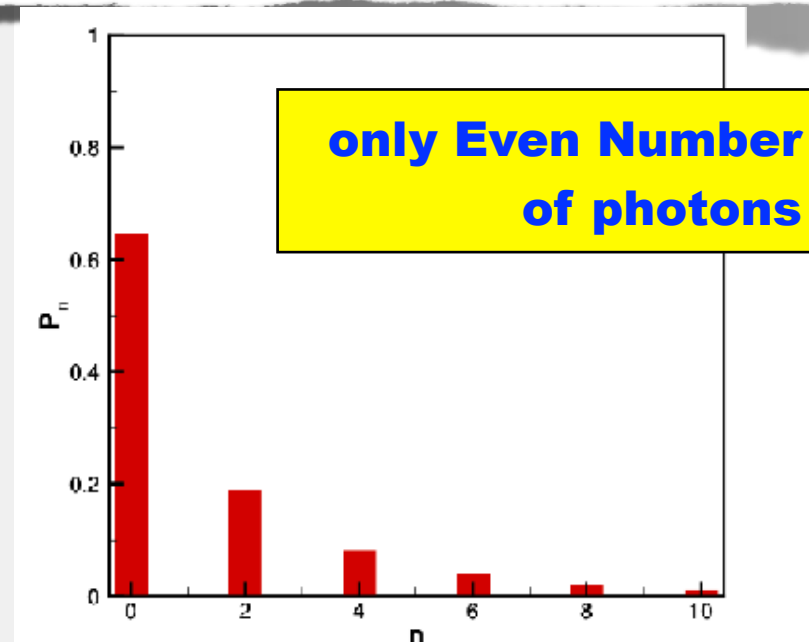
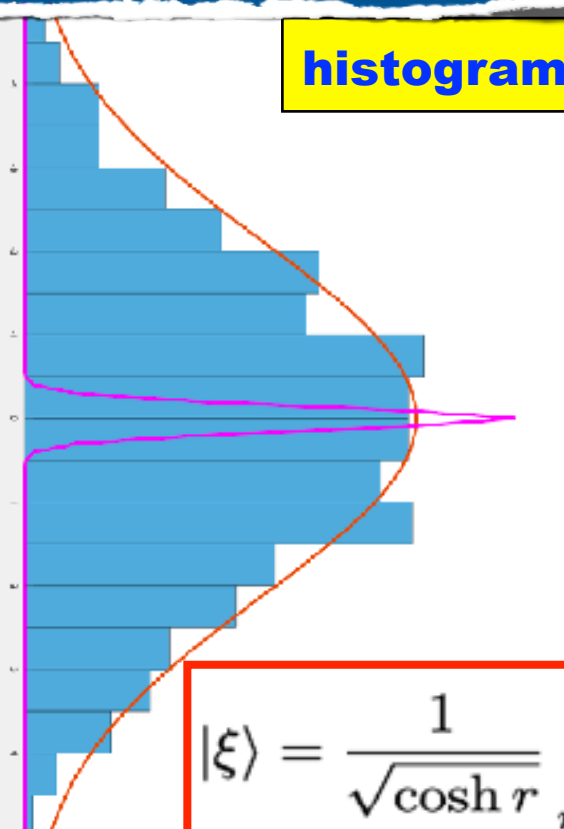
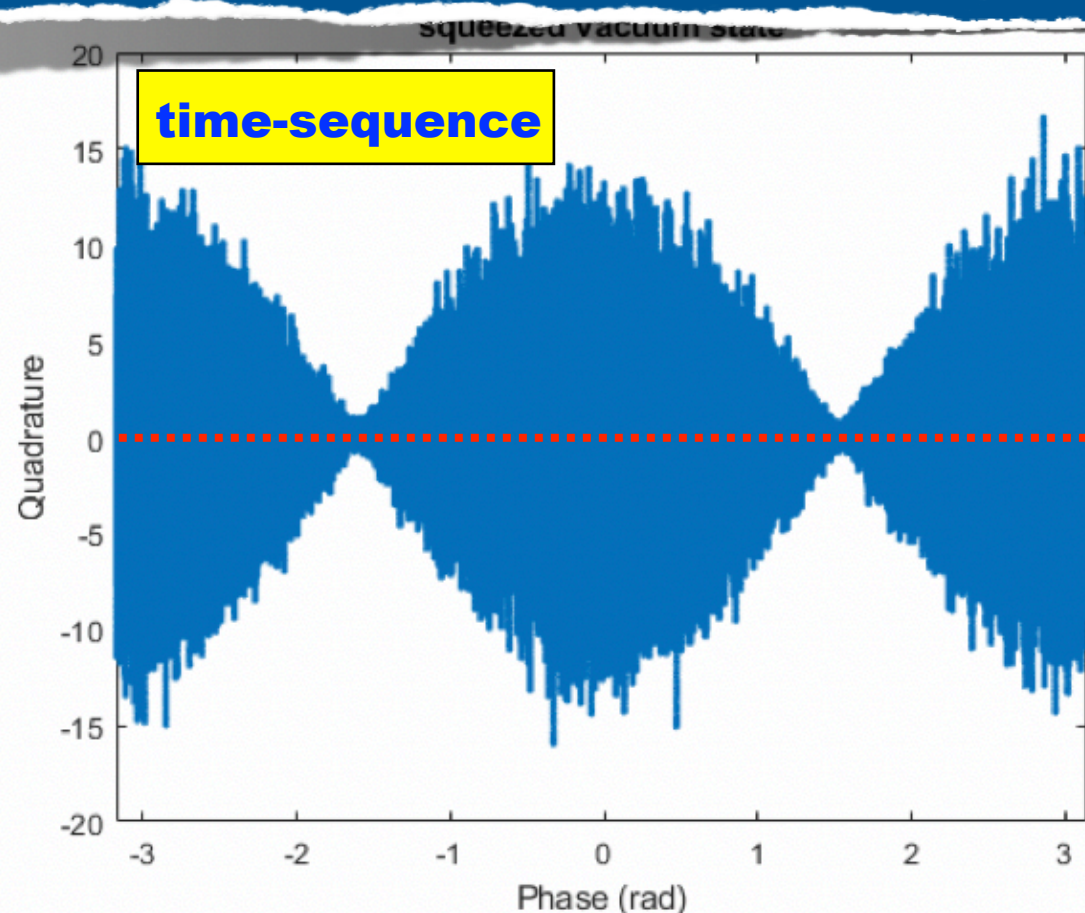


FIG. 2: Photon number distribution for squeezed coherent states. Left: $|\alpha|^2 = 50, \theta = 0, r = 0.5$; Right: $|\alpha|^2 = 50, \theta = 0, r = 4.0$.

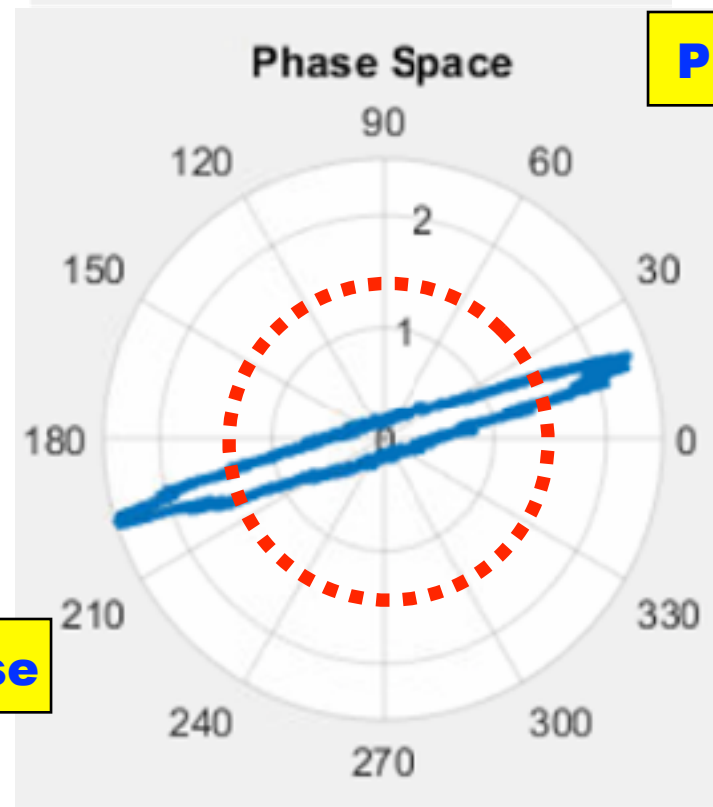
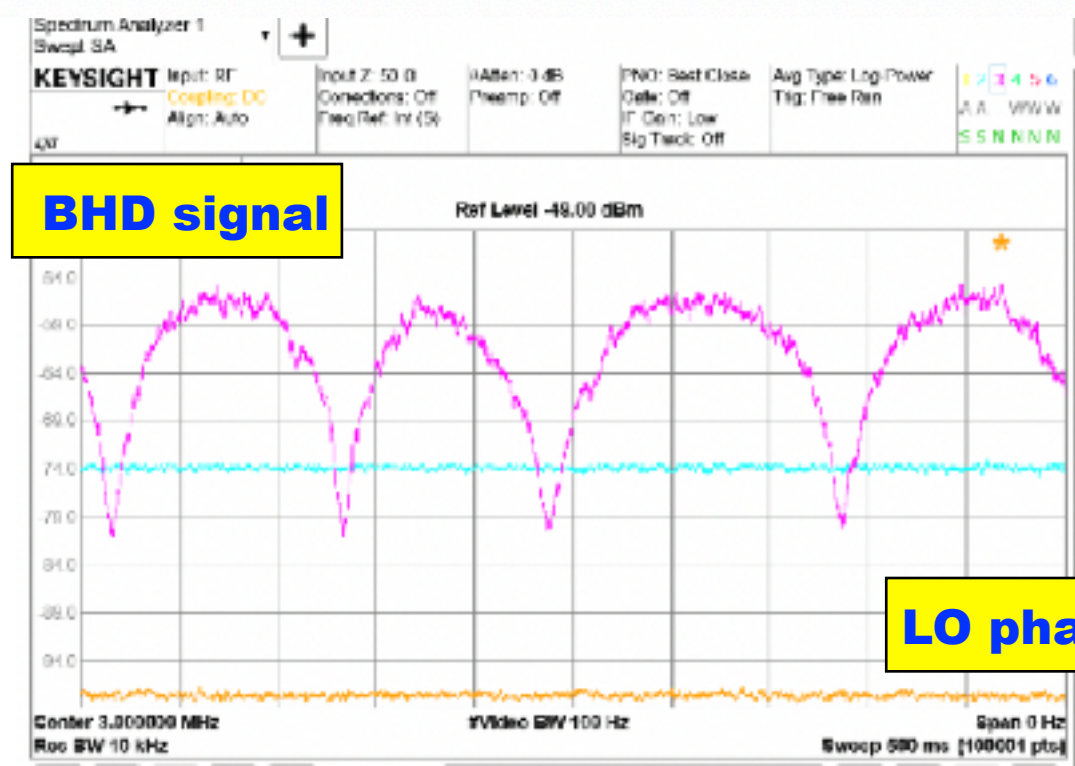
Squeezed State as a Eigen-state of UR

$$(\mu \hat{a} + \nu \hat{a}^\dagger) |\alpha, \xi\rangle = (\alpha \cosh r + \alpha^* \sinh r) |\alpha, \xi\rangle \equiv \gamma |\alpha, \xi\rangle,$$

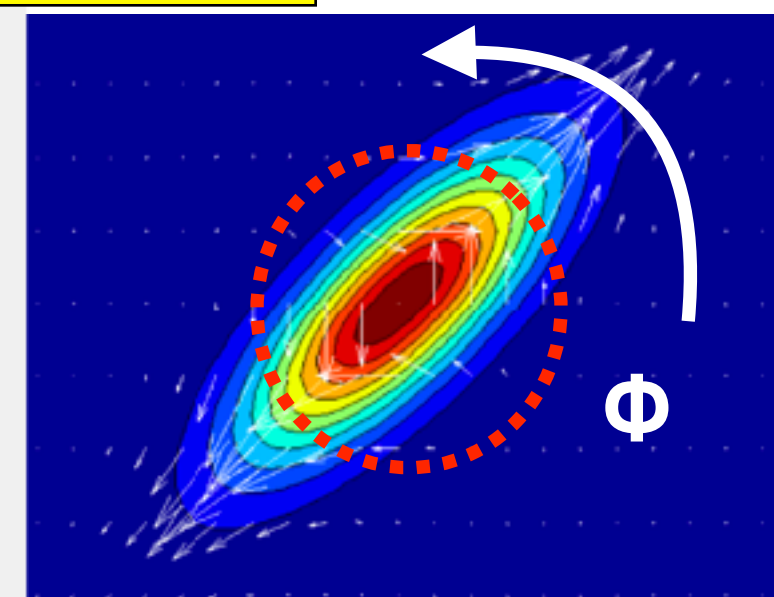
Squeezed States



$$|\xi\rangle = \frac{1}{\sqrt{\cosh r}} \sum_{m=0}^{\infty} (-1)^m \frac{\sqrt{(2m)!}}{2^m m!} e^{im\theta} \tanh^m r |2m\rangle$$

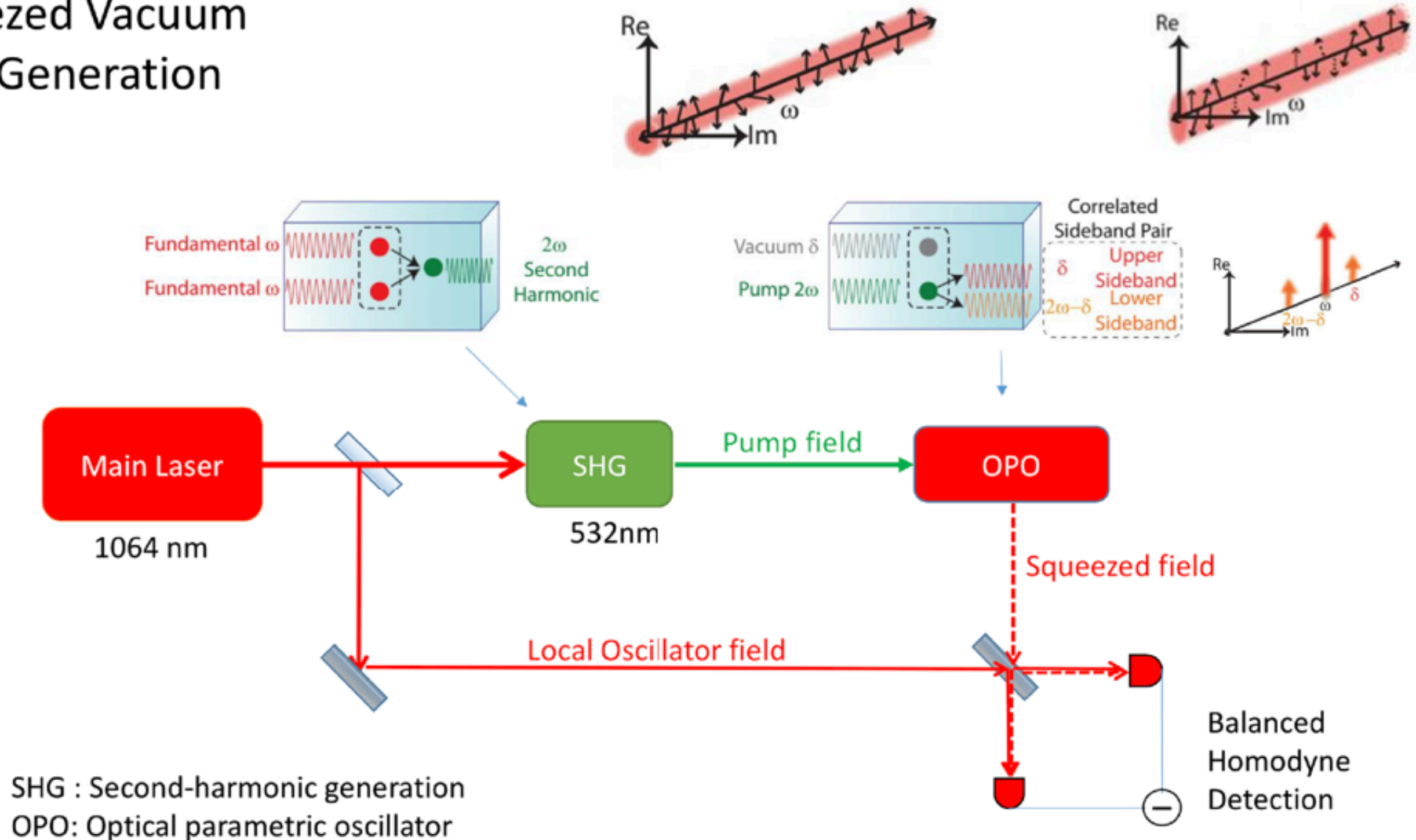


Phase Sapce



Optical Parametric Oscillator, OPO

Squeezed Vacuum State Generation



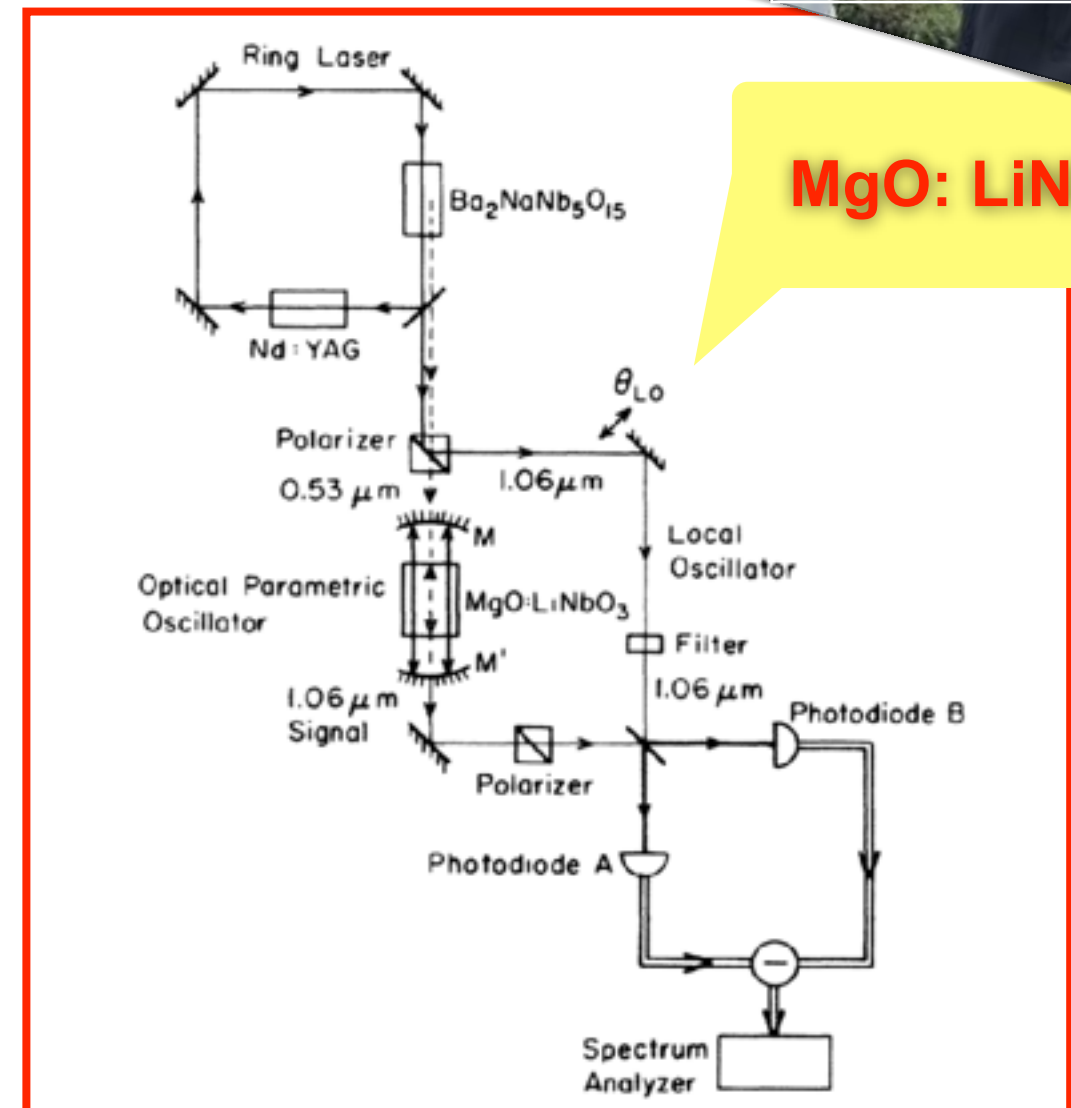
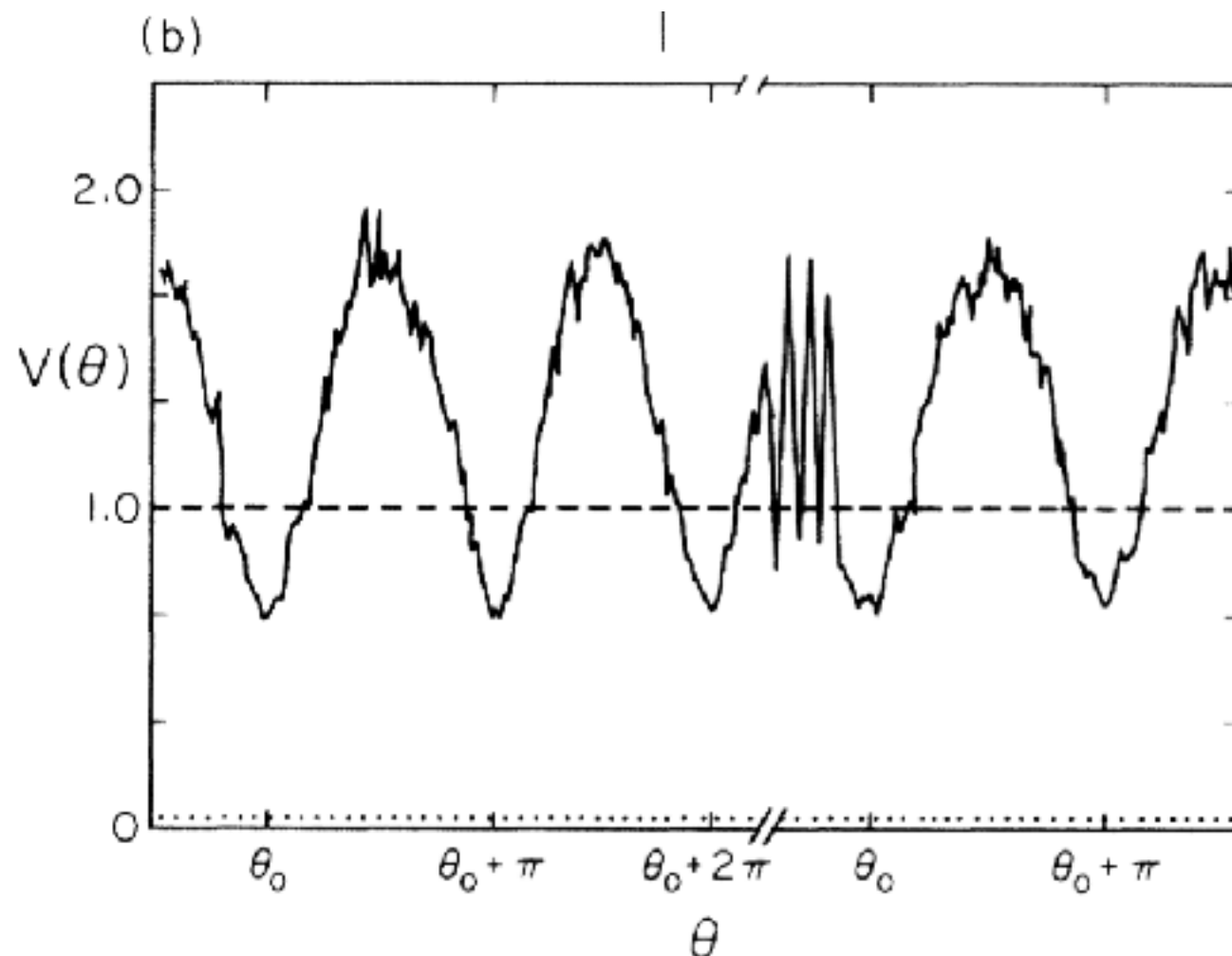
Generation of Squeezed States by Parametric

Ling-An Wu, H. J. Kimble, J. L. Hall,^(a)*Department of Physics, University of Texas at Austin*

(Received 11 September 1986)

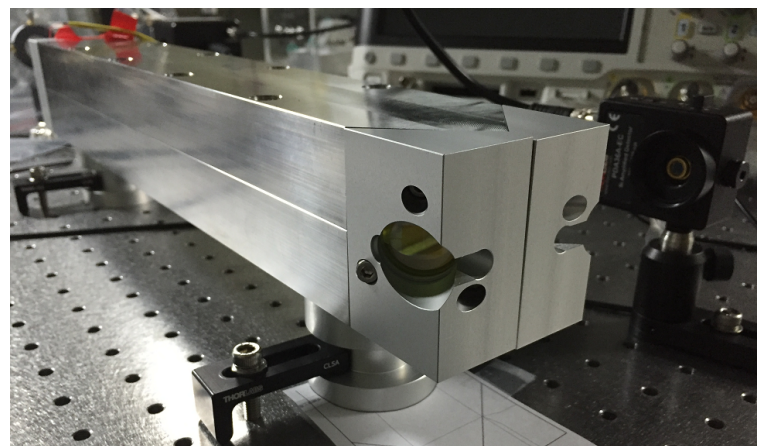
Squeezed states of the electromagnetic field are generated by in an optical cavity. Noise reductions greater than 50% relative in a balanced homodyne detector. A quantitative comparison with theory squeezing results from a field that in the absence of linear attenuation would be squ then tenfold.

PACS numbers: 42.50.Dv, 03.65-w, 42.65.Ky

MgO: LiNbO₃

Oct. 12th, 2019

1ST SQUEEZER @ TAIWAN



1064 nm, 2W,
BW < 1 kHz

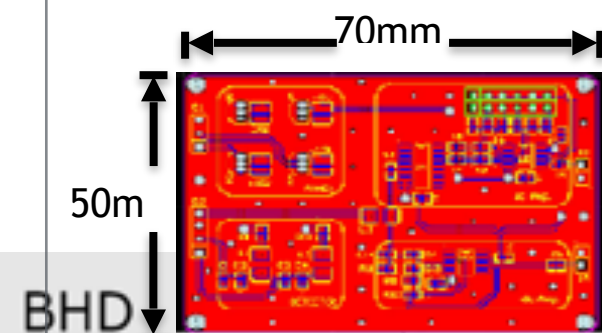
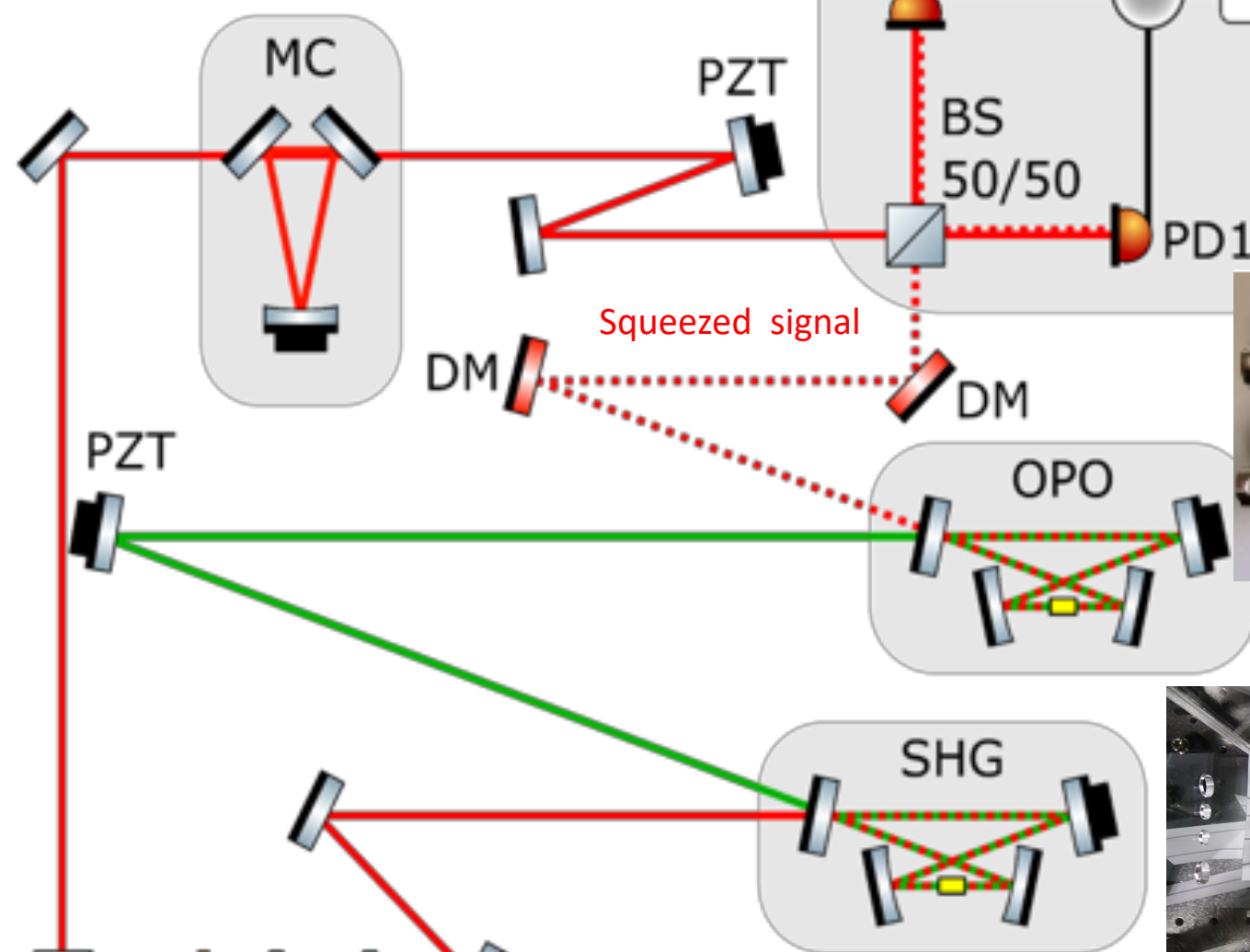
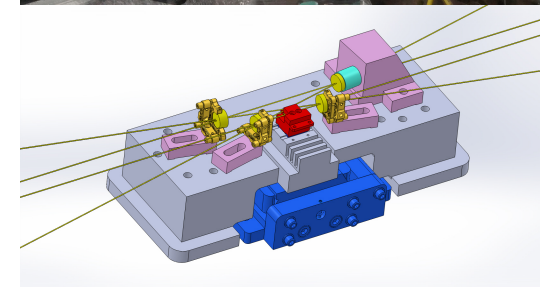
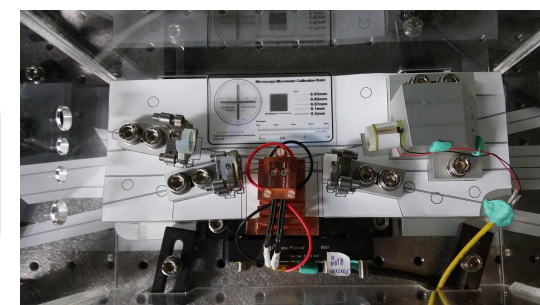
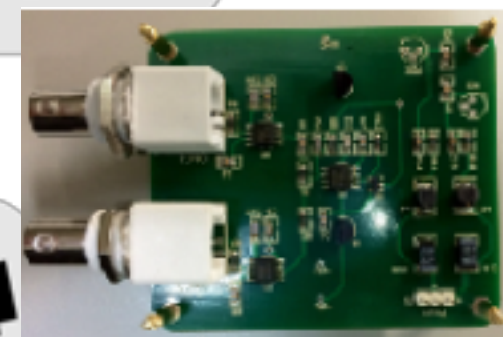


Fig.



Balanced Homodyne Detector, BHD

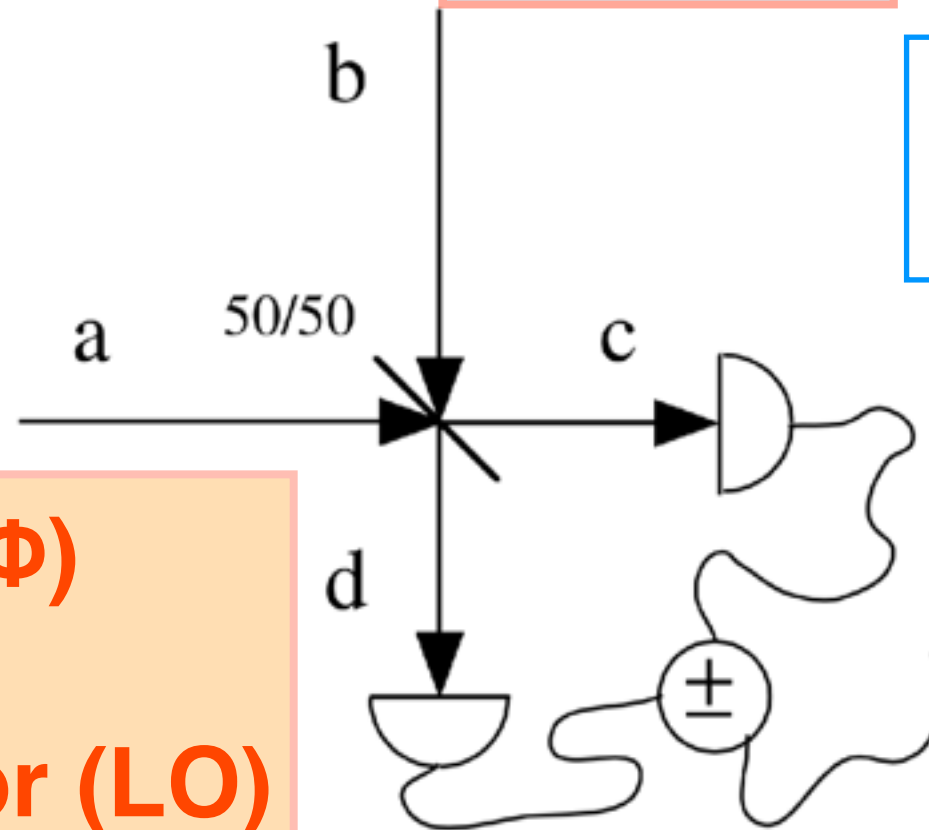
$$\begin{aligned}\tilde{a} &= \alpha + \delta\tilde{a}(\omega) \\ \tilde{b} &= \beta + \delta\tilde{b}(\omega)\end{aligned}$$

$$|0\rangle \text{ or } |\xi\rangle$$

$$\begin{aligned}\tilde{c} &= \sqrt{1-\xi}\tilde{a} - \sqrt{\xi}\tilde{b} \\ \tilde{d} &= \sqrt{\xi}\tilde{a} + \sqrt{1-\xi}\tilde{b}\end{aligned}$$

$$|\alpha\rangle * \exp(i\Phi)$$

Local Oscillator (LO)



$$\tilde{n}^- = 0 + \alpha \delta\tilde{X}_1^b$$

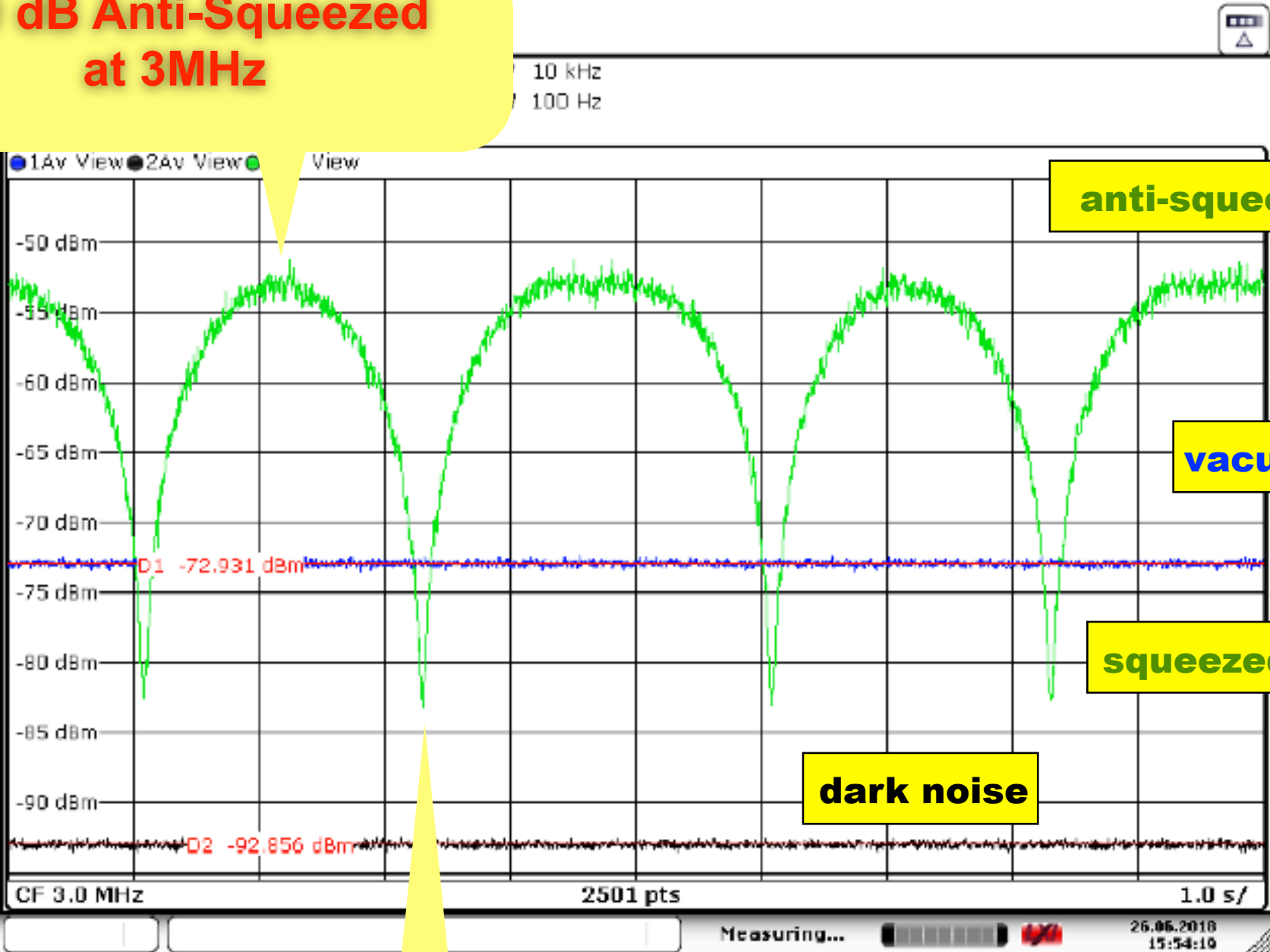
$$\begin{aligned}\tilde{n}^- &= \tilde{c}^\dagger \tilde{c} - \tilde{d}^\dagger \tilde{d} \\ &= 2\alpha\beta + \alpha(\delta\tilde{b} + \delta\tilde{b}^\dagger) + \beta(\delta\tilde{a} + \delta\tilde{a}^\dagger) \\ &= 2\alpha\beta + \alpha\delta\tilde{X}_1^b + \beta\delta\tilde{X}_1^a\end{aligned}$$



OPO 532nm incident power: 96mW
MC output(LO beam)= 14.5mW

specially ordered InGaAs Photodiodes
Laser Components GmbH $\phi=500\mu\text{m}$ $\text{QE}\geq 99\%$

+20 dB Anti-Squeezed
at 3MHz



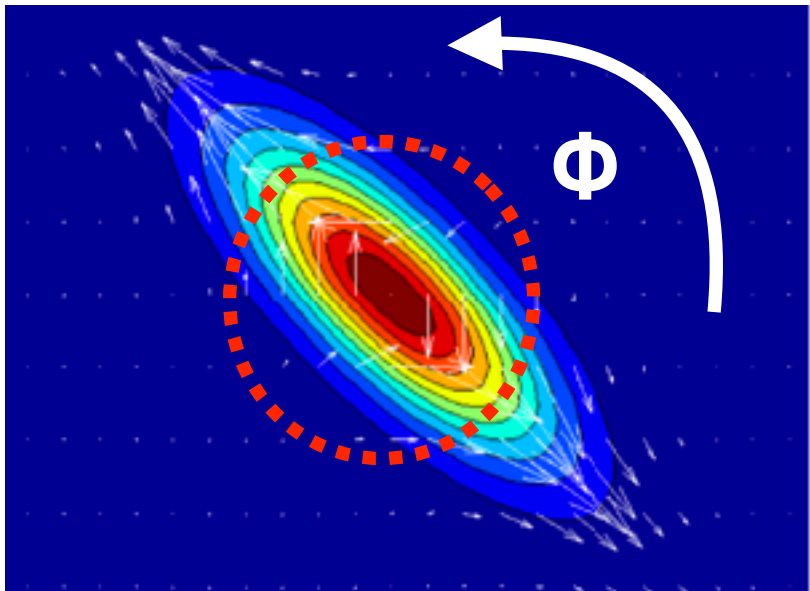
Clearance= 20dB

Zero Span mode at 3MHz
RBW=10kHz
VBW=100Hz

Blue Line: 14.5mW Vacuum noise
Black Line: Dark noise

-10 dB Squeezed vacuum
at 3MHz

Squeezing angle, ϕ



Date: June 26th, 2018

by Chien-Ming Wu