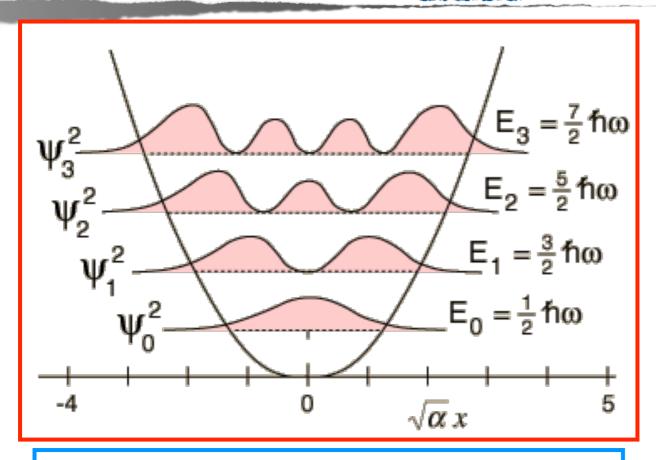
## **Note: Quantum SHO**

- Quantum Simple Harmonic Oscillator, qSHO
- Photons occupy an electromagnetic mode (referred as the modes in quantum optics, typically a plane wave)
  - **Hamiltonian**
  - Number operator
  - ☐ Energy Quantization (equally spacing in energy)
- •The energy in a mode is not continuous but discrete in quanta.
  - ☑ Vacuum state with zero-point energy
- There is a zero point energy inherent to each mode, which is equivalent with fluctuations of the electromagnetic field in vacuum, due to the uncertainty principle.
  - Schrodinger picture
  - ☐ Heisenberg picture
- •The observables are just represented by probabilities as usual in QM.

## Note: Quantum Mechanics

- Axioms
- ✓ State
- ☑ Density Matrix
- More on States
- □ Coherent States
- □ Squeezed States
- □ Uncertainty Relation → Minimum Uncertainty States
- Entropy
- Purity
- □ bi-particle States → Entanglement (Schmidt decomposition)
- □ Cat states

# **Quantum Simple Harmonic Oscillator (SHO)**



- Energy quantization
- Equally spacing in energy difference
- Zero-point energy  $\neq 0$

$$\psi(\xi) = H_n(\xi) \exp[-\xi^2/2], \qquad \epsilon = 2n+1, \qquad n = 0, 1, 2, 3 \dots$$

$$E = \frac{\hbar \omega}{2} \epsilon = \hbar \omega (n + \frac{1}{2}), \qquad n = 0, 1, 2, 3, \dots$$

$$\hat{H} = \frac{1}{2} \frac{\hat{p}^2}{m} + \frac{1}{2} k \, \hat{x}^2, \ \ [\hat{x}, \hat{p}] = i \hbar.$$

$$\hat{H}=\hbar\omega(\hat{a}^{\dagger}\hat{a}+\frac{1}{2}).\quad [\hat{a},\hat{a}^{\dagger}]=1,$$

$$\hat{N}|n\rangle = n|n\rangle,$$
 $\hat{a}|n\rangle = \sqrt{n}|n-1\rangle,$ 
 $\hat{a}^{\dagger}|n\rangle = \sqrt{n+1}|n+1\rangle,$ 
 $E_n = \hbar\omega(n+\frac{1}{2}).$ 



## **Poisson Distribution:**

$$P(n) = \frac{\bar{n}^n \exp(-\bar{n})}{n!},$$

$$\langle \hat{n} \rangle = \sum_{n} nP(n) = |\alpha|^2 \equiv \bar{n},$$
  
 $\langle \Delta \hat{n}^2 \rangle = \langle \hat{n}^2 \rangle - \langle \hat{n} \rangle^2 = |\alpha|^2 = \langle \hat{n} \rangle.$ 

mean = variance

### **Bose-Einstein Distribution:**

Boltzmann's law

$$P(n) \propto \exp[-E_n/k_BT],$$

$$P(n) = \frac{\exp[-E_n/k_B T]}{\sum_{n=0}^{\infty} \exp[-E_n/k_B T]},$$
  
=  $\exp[-E_n/k_B T] (1 - \exp[-\hbar\omega/k_B T]); \qquad E_n = n \hbar\omega$ 

$$\bar{n} = \sum_{n=0}^{\infty} n \, P(n) = \frac{1}{\exp[\hbar \omega/k_B T] - 1},$$
 • average photon number at temperature T

$$P(n) = \frac{1}{\bar{n}+1} (\frac{\bar{n}}{\bar{n}+1})^n,$$

## **Bose-Einstein Distribution:**

thermal state

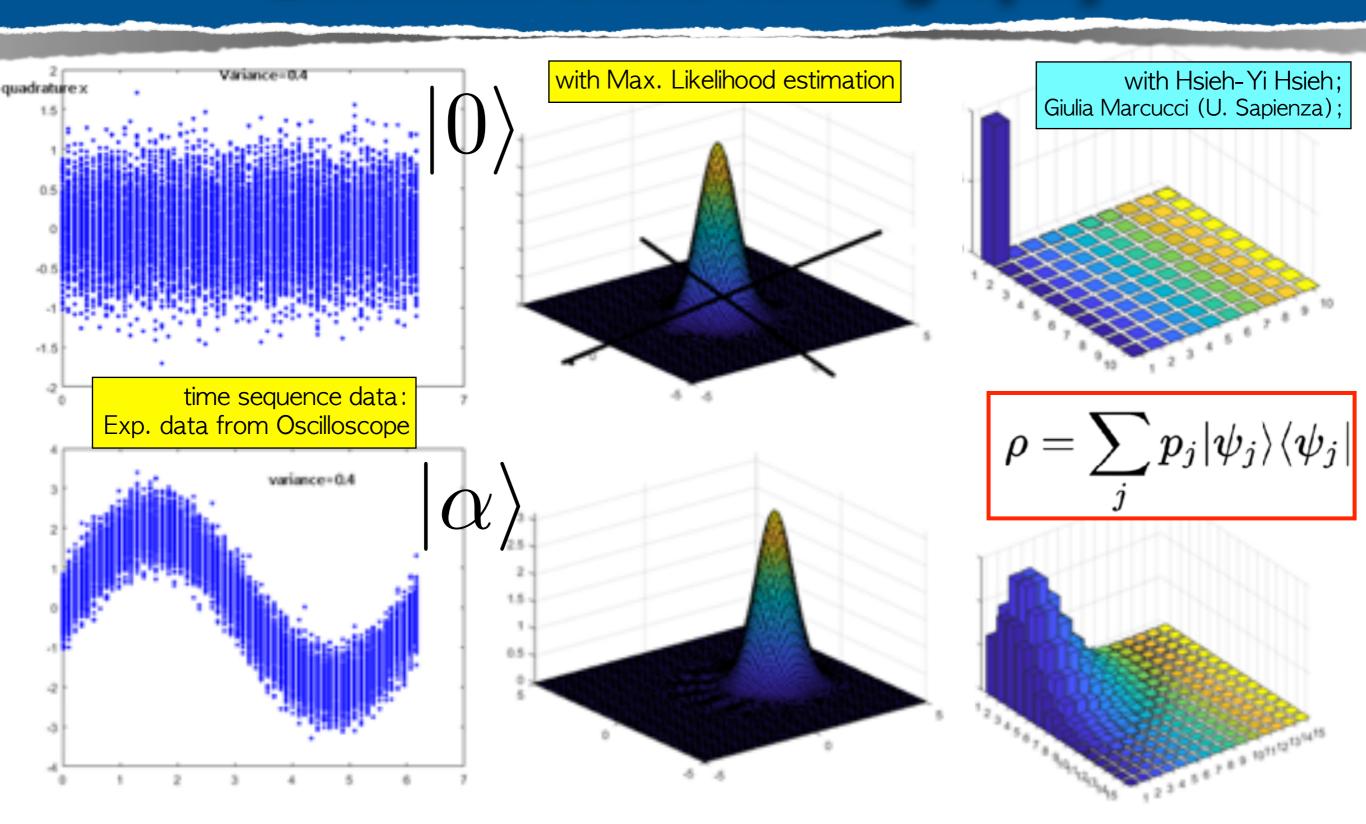
$$\rho_{th} = \sum_{n} = \frac{1}{\bar{n}+1} \left(\frac{\bar{n}}{\bar{n}+1}\right)^{n} |n\rangle\langle n|.$$

$$\Delta n^2 = \bar{n} + \bar{n}^2,$$

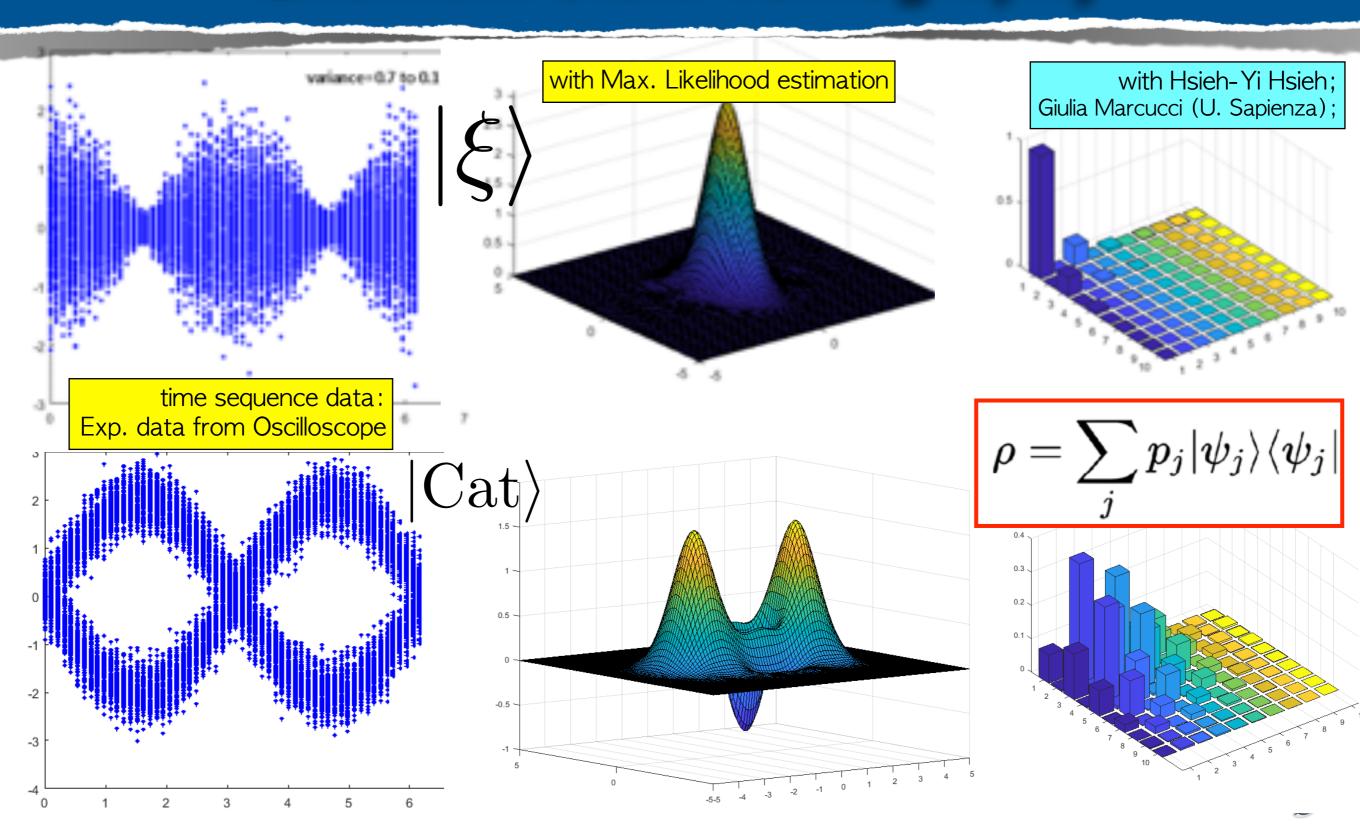
# Note: Coherent States (CS)

- □ Eigenstate of Annihilation operator
- Displacement Operator
- Properties of CS
- Representation of CS
- □ Expectation Value of E-fields
- Generation of CS
- More on States
- Minimum Uncertainty States
- □ Uncertainty Relation → Minimum Uncertainty States
- □ Squeezed States
- □ CS in Phase space
- ¬ Max. Mixed CS
- □ Generalized CS
- □ Spin Coherent States
- — Fermionic Coherent States

# **Quantum State Tomography**



# Quantum State Tomography



## **Poisson Distribution:**

$$P(n) = \frac{\bar{n}^n \exp(-\bar{n})}{n!},$$

$$\langle \hat{n} \rangle = \sum_{n} nP(n) = |\alpha|^2 \equiv \bar{n},$$
  
 $\langle \Delta \hat{n}^2 \rangle = \langle \hat{n}^2 \rangle - \langle \hat{n} \rangle^2 = |\alpha|^2 = \langle \hat{n} \rangle.$ 

• mean = variance

We introduce the eigenstate of annihilation operator, called the *coherent state*,

$$\hat{a}|\alpha\rangle = \alpha|\alpha\rangle,$$

# Eigenstate of $\hat{a}$ :

We introduce the eigenstate of annihilation operator, called the *coherent state*,

$$\hat{a}|\alpha\rangle = \alpha|\alpha\rangle,$$

$$|lpha
angle = e^{-rac{1}{2}|lpha|^2} \sum_{n=0}^{\infty} rac{lpha^n}{\sqrt{n!}} |n
angle.$$

• mean = variance

# Displacement Operator:

$$|lpha
angle=\hat{D}(lpha)|0
angle=e^{+lpha\hat{a}^{\dagger}-lpha^{*}\hat{a}}|0
angle,$$

- 1. The probability of finding n photons in  $|\alpha\rangle$  is given by a Poisson distribution.
- 2. The coherent state is a minimum-uncertainty states,
- 3. The set of all coherent states  $|\alpha\rangle$  is a complete set,

$$\int |\alpha\rangle\langle\alpha|d^2\alpha = \pi \sum_n |n\rangle\langle n|, \quad \text{or} \quad \frac{1}{\pi} \int |\alpha\rangle\langle\alpha|d^2\alpha = 1.$$
 (1)

4. Two coherent states corresponding to different eigenstates  $\alpha$  and  $\beta$  are not orthogonal,

$$\langle \alpha | \beta \rangle = \exp(-\frac{1}{2}|\alpha|^2 + \alpha^*\beta - \frac{1}{2}|\beta|^2) = \exp(-\frac{1}{2}|\alpha - \beta|^2). \tag{2}$$

5. Coherent states are approximately orthogonal only in the limit of large separation of the two eigenvalues,  $|\alpha - \beta| \to \infty$ . Therefore, any coherent state can be expanded using other coherent state,

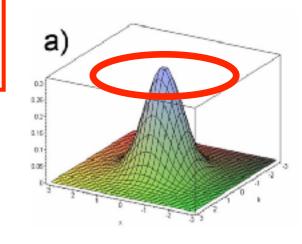
$$|\alpha\rangle = \frac{1}{\pi} \int d^2\beta |\beta\rangle \langle\beta|\alpha\rangle = \frac{1}{\pi} \int d^2\beta e^{-\frac{1}{2}|\beta-\alpha|^2} |\beta\rangle.$$
 (3)

This means that a coherent state forms an overcomplete set.

6. The simultaneous measurement of  $\hat{a}_1$  and  $\hat{a}_2$ , represented by the projection operator  $|\alpha\rangle\langle\alpha|$ , is not an exact measurement but instead an approximate measurement with a finite measurement error.

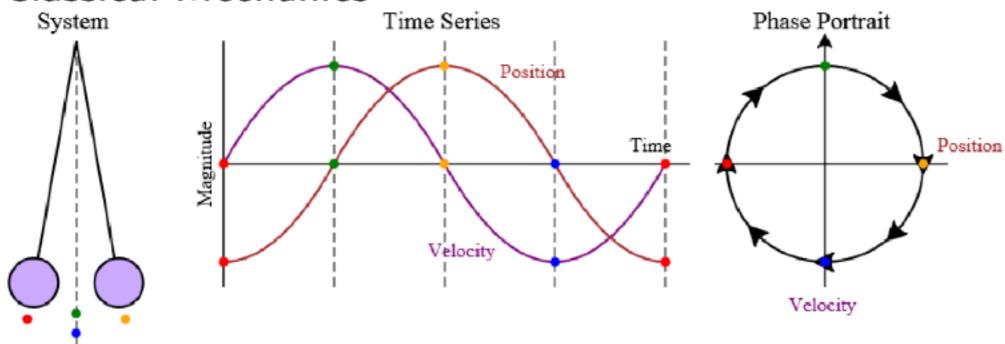
# Representation:

$$\langle q | \alpha \rangle = (\frac{\omega}{\pi \hbar})^{1/4} \exp[-\frac{\omega}{2\hbar} (q - \langle q \rangle)^2 + i \frac{\langle p \rangle}{\hbar} q + i\theta],$$



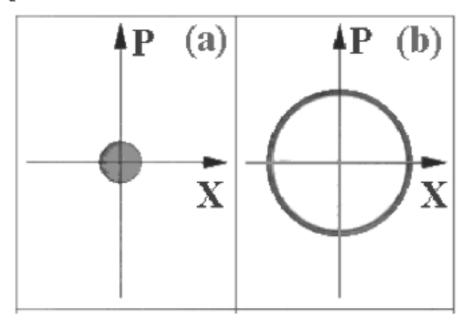
## Phase space

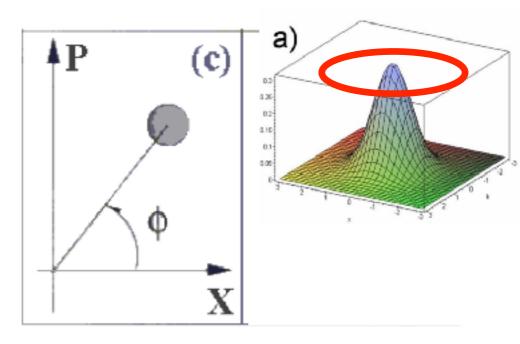
#### Classical Mechanics

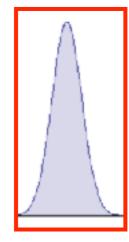


wave-nature

#### Quantum Mechanics

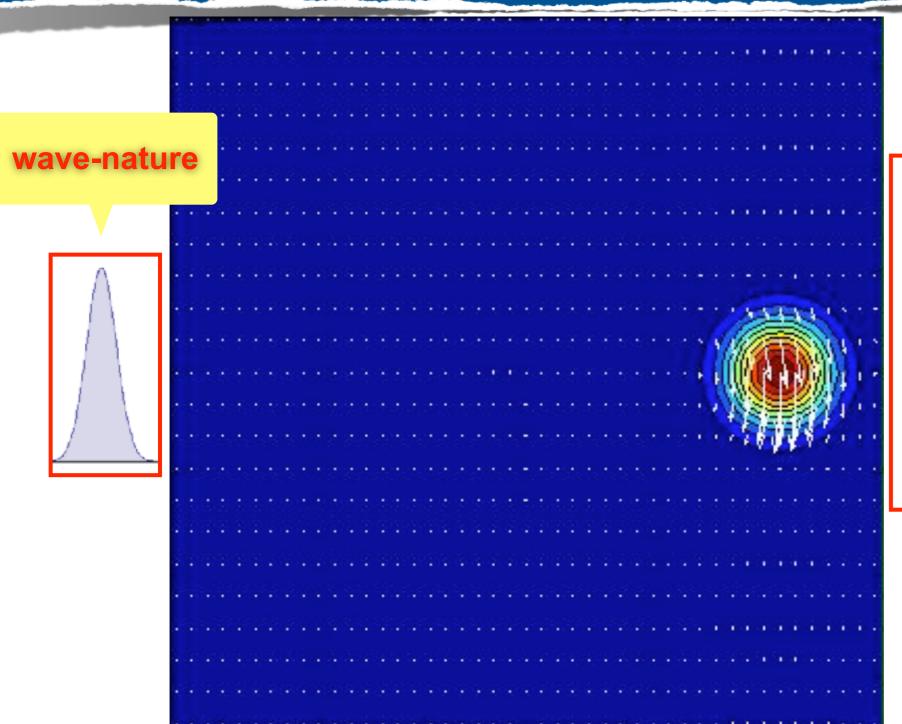


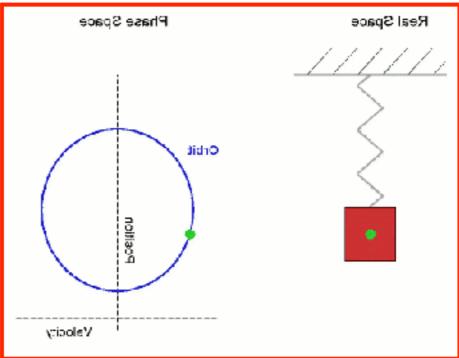




from Wiki

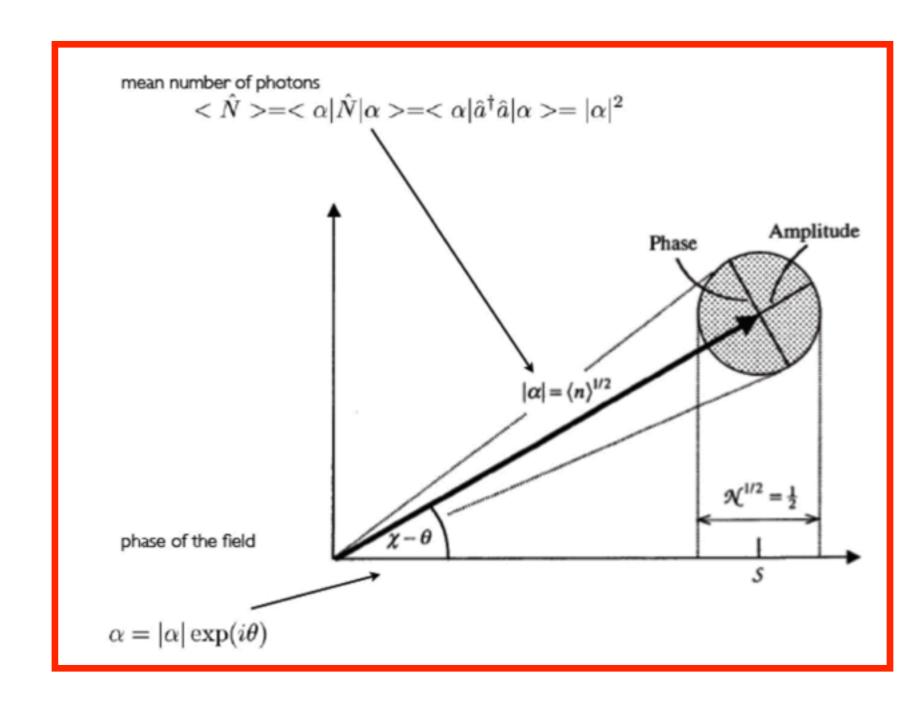
### **Coherent states**





with Popo Yang

# **Expectation value of E-fields:**



# **Generation of CS:**