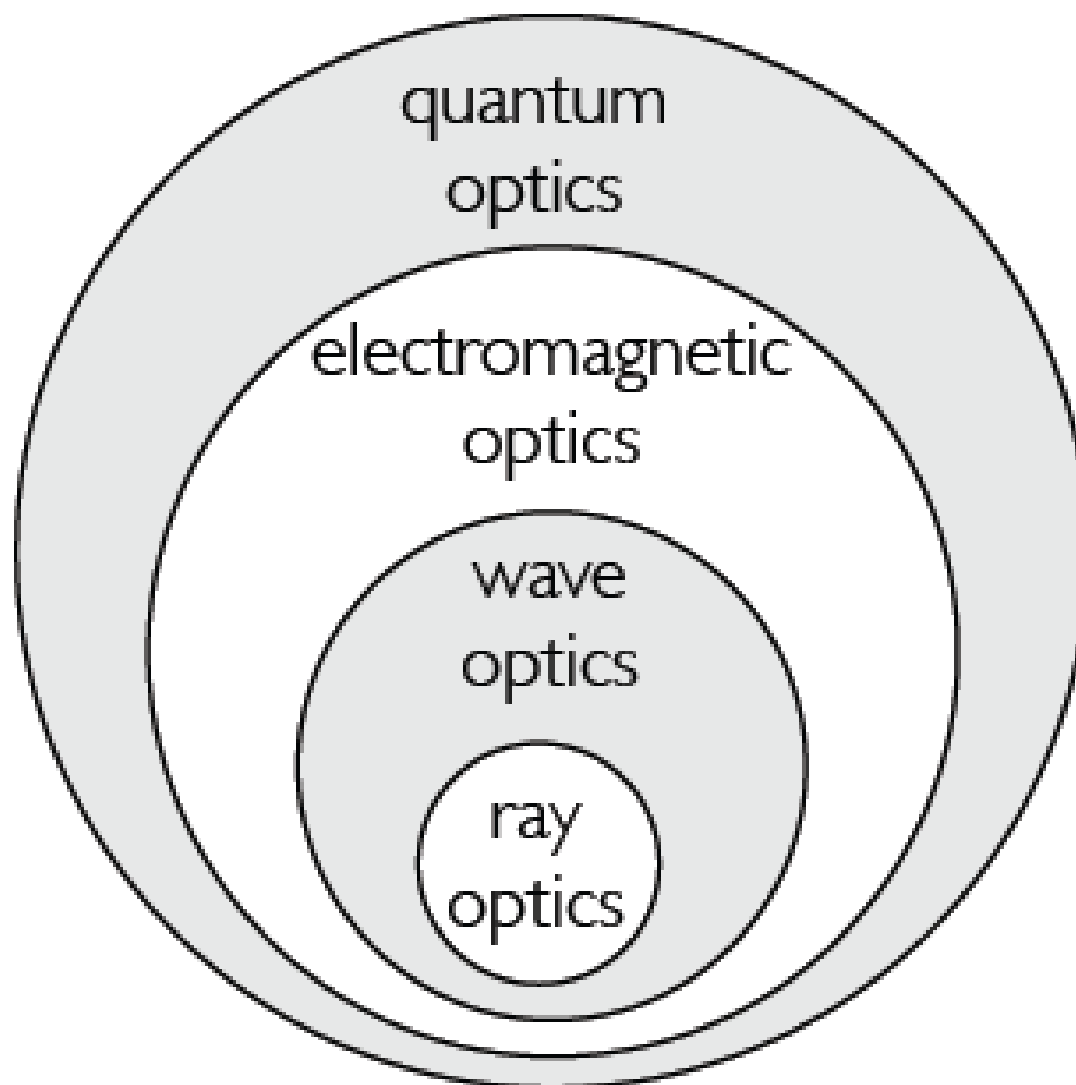


General Optics



Wave Optics

- Ray optics is wave optics for infinitely small wavelength
- Wave optics:
 - Plane waves
 - Spherical waves
 - Interference
- Diffraction
- Gaussian beams

Postulates of ray optics

- light travels in form of waves
- optical medium is characterized by a quantity n

$$c = \frac{c_0}{n},$$

- optical wave satisfies wave equation of type

$$\nabla^2 u - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} = 0,$$

- wave equation is linear - superposition
if u_1 and u_2 is solution then also $a_1 u_1 + a_2 u_2$,
- optical intensity, $I(r, t) = 2\langle u^2(r, t) \rangle$,
i.e. averaging over times longer than 1 optical cycle,
- optical power, $P(t) = \int_A I(r, t) dA$,

Wave equations

- ➔ For a *source-free* medium, $\rho = \mathbf{J} = 0$,

$$\begin{aligned}\nabla \times (\nabla \times E) &= -\mu\epsilon \frac{\partial^2}{\partial t^2} E, \\ \Rightarrow \nabla(\nabla \cdot E) - \nabla^2 E &= -\mu\epsilon \frac{\partial^2}{\partial t^2} E.\end{aligned}$$

- ➔ When $\nabla \cdot E = 0$, one has *wave equation*,

$$\nabla^2 E = \mu\epsilon \frac{\partial^2}{\partial t^2} E$$

- ➔ which has following expression of the solutions, in 1D,

$$\begin{aligned}E &= \hat{x}[f_+(z - vt) + f_-(z + vt)], \\ H &= \sqrt{\frac{\epsilon}{\mu}} \hat{y}[f_+(z - vt) - f_-(z + vt)],\end{aligned}$$

with

$$\mu\epsilon = \mu_0\epsilon_0(1 + \chi) = \frac{n^2}{c^2},$$

Plane waves

→ 1D wave equation,

$$\frac{\partial^2}{\partial z^2} E = \mu\epsilon \frac{\partial^2}{\partial t^2} E,$$

which has the solutions of

$$E = \hat{x} [f_+(z - vt) + f_-(z + vt)], \text{ with}$$

$$v^2 = \frac{1}{\mu\epsilon} = \frac{n^2}{c_0^2},$$

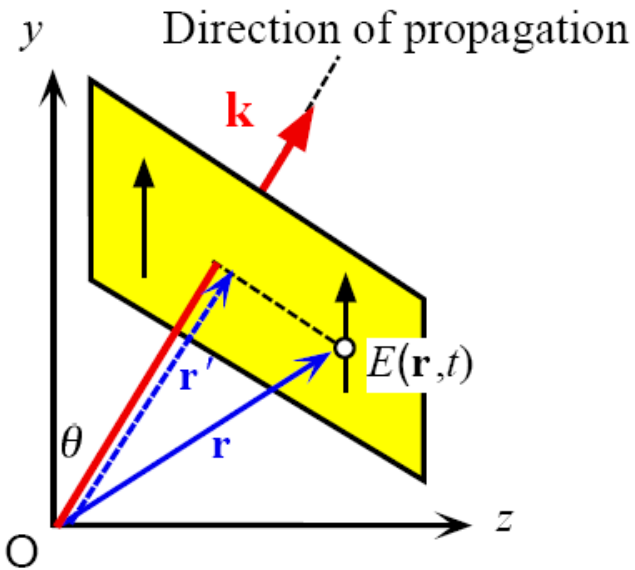
→ plane wave solutions:

$$E_+ = E_0 \cos(kz - \omega t),$$

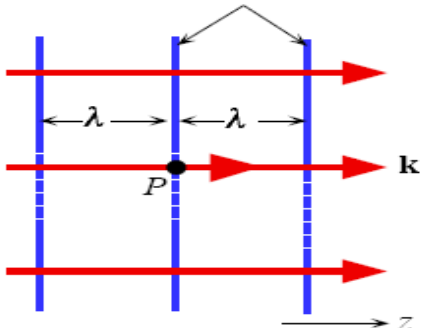
where $\frac{\omega}{k} = \frac{c_0}{n}$.

Travelling waves

A travelling plane EM wave along a direction k .

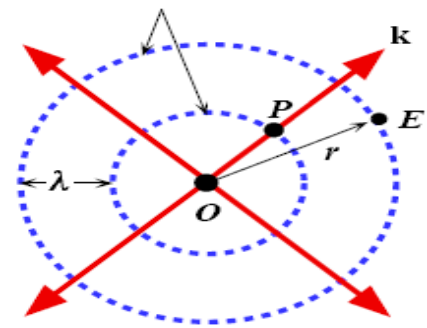


Wave fronts (constant phase surfaces)



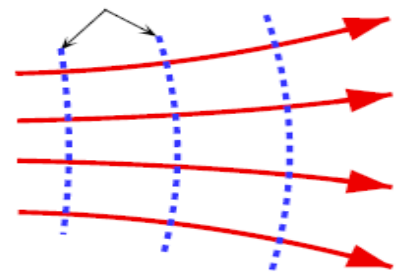
A perfect plane wave (a)

Wave fronts



A perfect spherical wave (b)

Wave fronts



A divergent beam (c)

solution of wave equation

- monochromatic wave is solution of wave equation

$$u(r, t) = a(r) \cos[\omega t + \phi(r)],$$

where $a(r)$ is the *amplitude*, $\omega = 2\pi\mu$ is the *frequency*, and $\phi(r)$ is the *phase*.

- complex representation:

$$u(r, t) = \mathbf{Re}\{U(r, t)\} = \frac{1}{2}[U(r, t) + U^*(r, t)],$$

where

$$U(r, t) = a(r)\exp[i\phi(r)]\exp(i\omega t) = U(r)\exp(i\omega t),$$

has to satisfy

$$\nabla^2 U - \frac{1}{c^2} \frac{\partial^2}{\partial t^2} U = 0,$$

complex representation and Helmholtz Equation

- separation time/space

$$U(r, t) = a(r)\exp[i\phi(r)]\exp(i\omega t) = U(r)\exp(i\omega t),$$

where

$$U(r) \text{ is complex,}$$

with the amplitude $|U(r)|$ and the phase $\arg\{U(r)\}$,

- Helmholtz wave equation

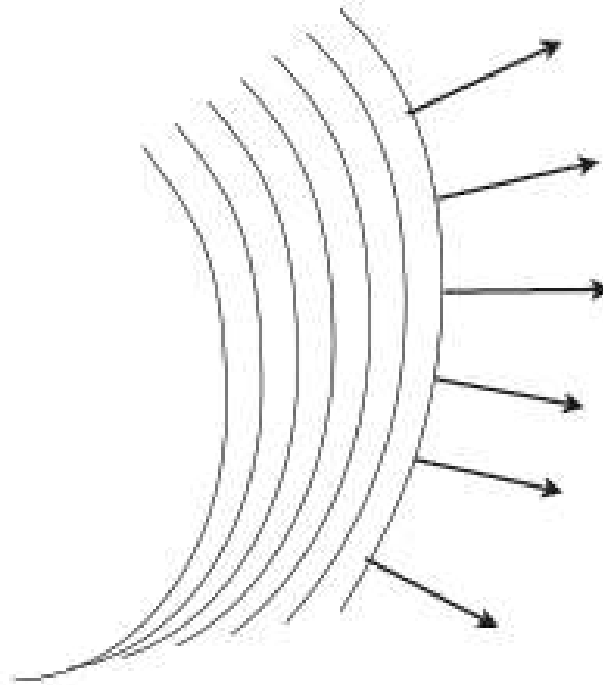
$$(\nabla^2 + k^2)U(r) = 0,$$

where

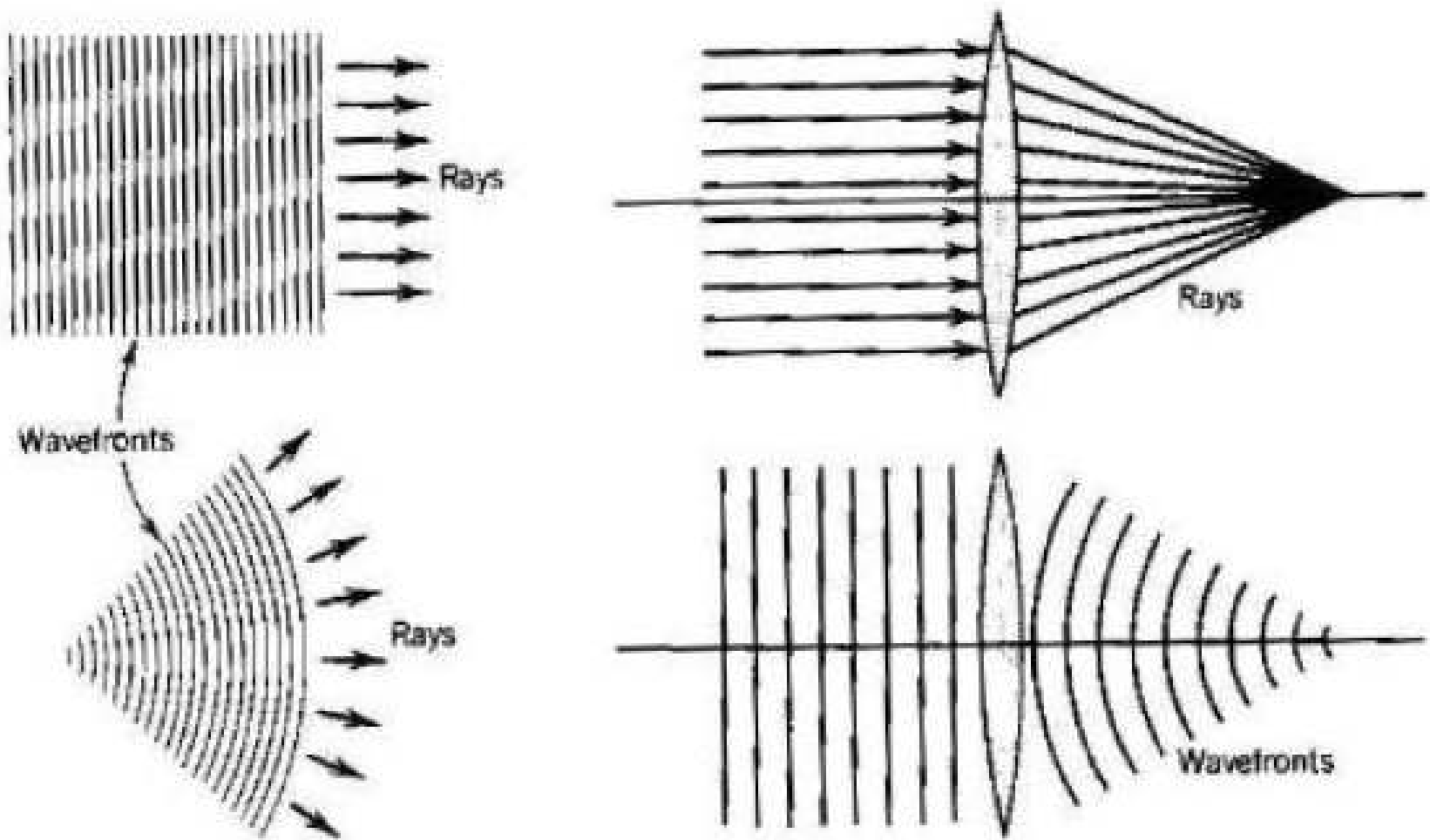
$$k = \frac{\omega}{c}, \quad \text{wavenumber,}$$

relation to ray optics

- intensity: $I(r) = |U(r)|^2$,
- wavefront: $\phi(r) = \text{constant}$,
- rays are normals to the wavefronts change in the curvature of wavefronts bends rays



Link between wave optics and ray optics



plane waves

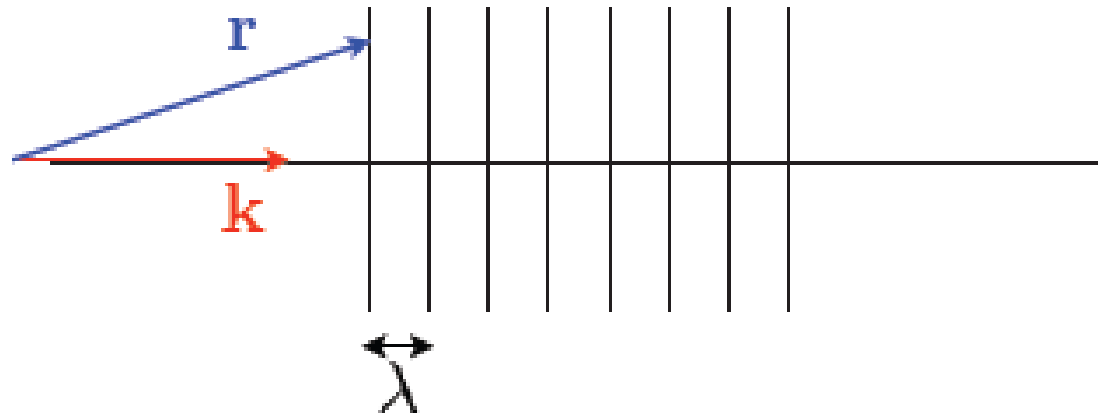
- plane wave


$$U(\mathbf{r}) = A \exp(-i\mathbf{k} \cdot \mathbf{r}),$$

where \mathbf{k} is the wavevector, defines propagation direction,

$$\mathbf{k} \cdot \mathbf{r} = 2\pi n, \quad n \text{ is an integer,}$$

- distance between neighboring wavefronts, $\lambda = \frac{2\pi}{k} = \frac{c}{\mu}$,
- in a medium with refractive index n , $\lambda = \frac{c}{\mu} = \frac{c_0}{n\mu} = \frac{\lambda_0}{n}$, and $k = nk_0$,
- intensity, $I(\mathbf{r}) = |A|^2$





As a monochromatic wave propagates through media of different refractive indices its **frequency remains the same** but its **velocity, wavelength and wavenumber are altered**.

spherical waves

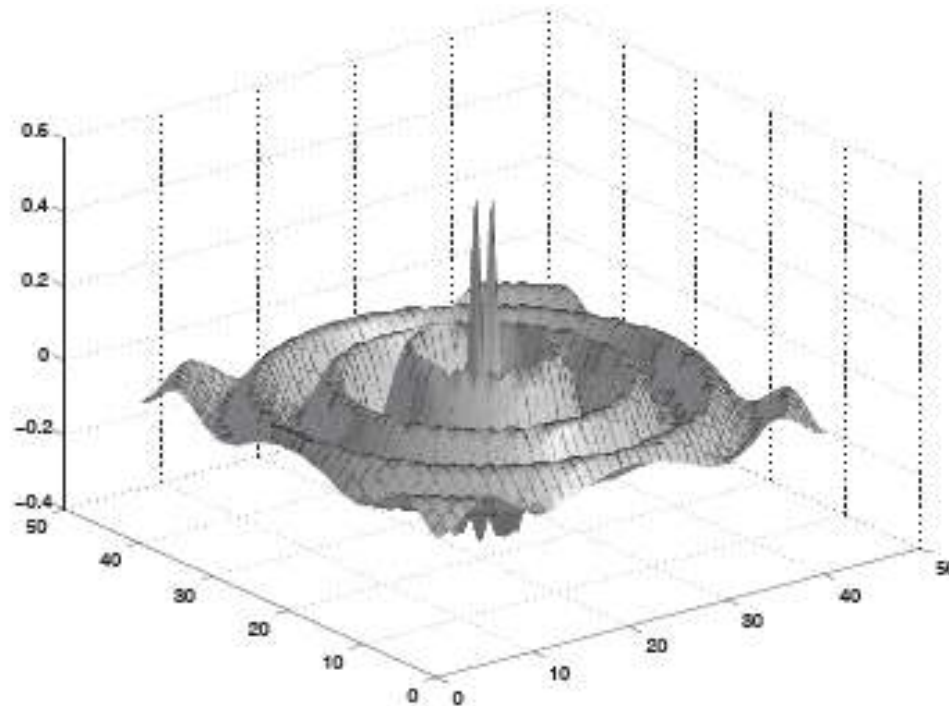
→ spherical wave:

$$U(r) = \frac{A}{|r - r_0|} \exp(-ik|r - r_0|),$$

where $k|r - r_0| = \text{constant}$, wavefronts resemble sphere surfaces,

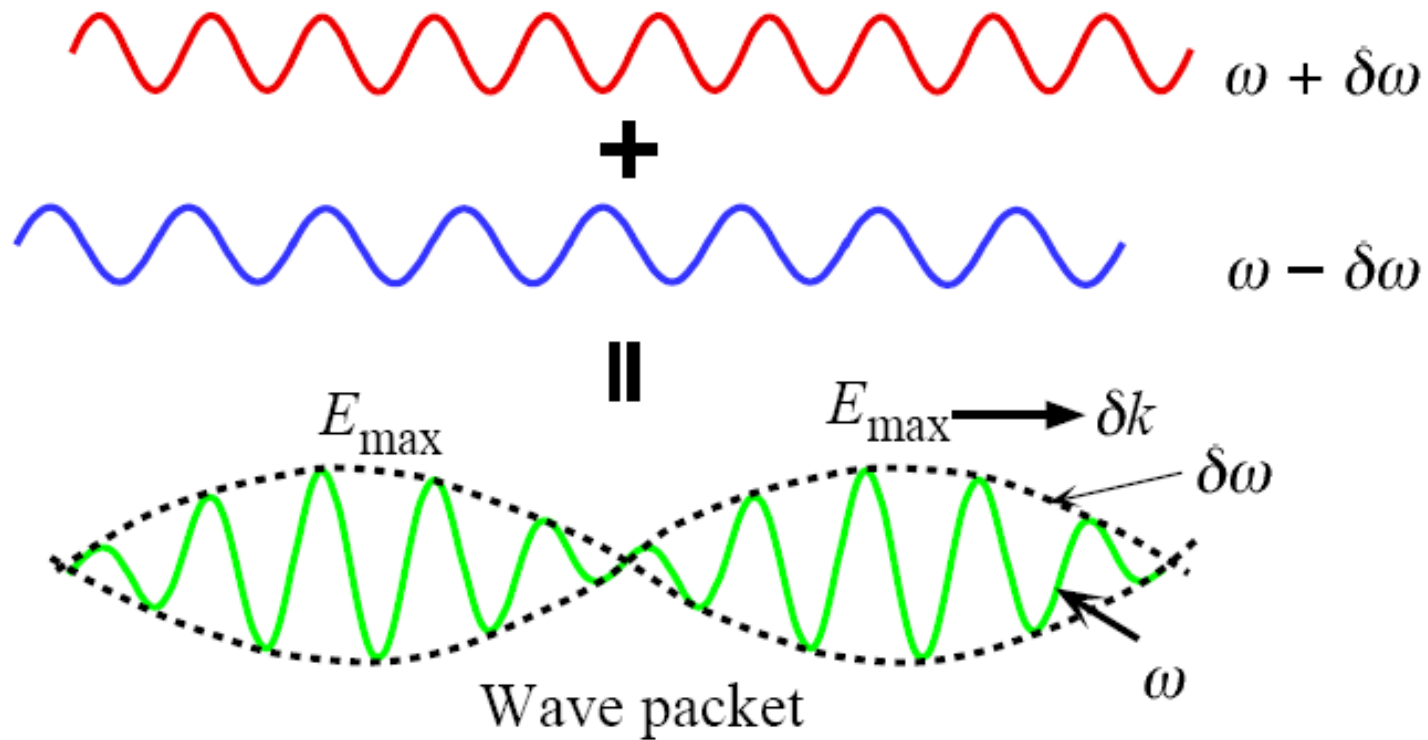
→ intensity:

$$I(r) = \frac{|A|^2}{r^2},$$



Phase velocity and Group velocity

- phase velocity: $v_p = c = \frac{c_0}{n} = \frac{\omega}{k}$,
- group velocity: $v_G = \frac{d\omega}{dk}$,



Interference (spatial)

- superposition of two monochromatic waves of the same frequency,

$$U(r) = U_1(r) + U_2(r),$$

- intensity:

$$I = |U(r)|^2 = |U_1 + U_2|^2 = |U_1|^2 + |U_2|^2 + U_1^*U_2 + U_1U_2^*,$$

- define:

$$U_1 = I_1^{1/2} \exp(i\phi_1), \quad U_2 = I_2^{1/2} \exp(i\phi_2),$$

then

$$I = I_1 + I_2 + 2(I_1 I_2)^{1/2} \cos \phi,$$

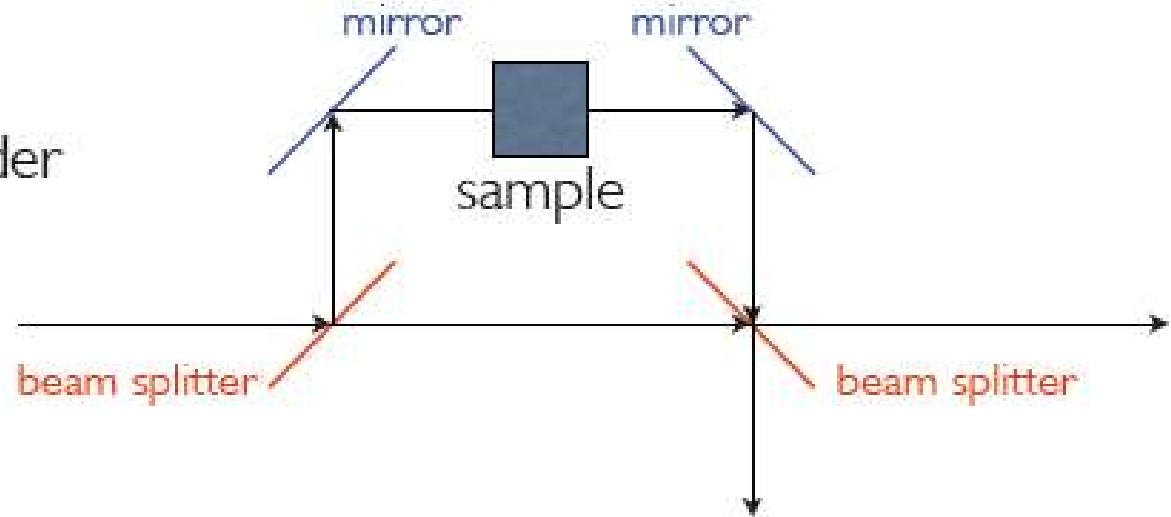
where

$$\phi = \phi_2 - \phi_1,$$

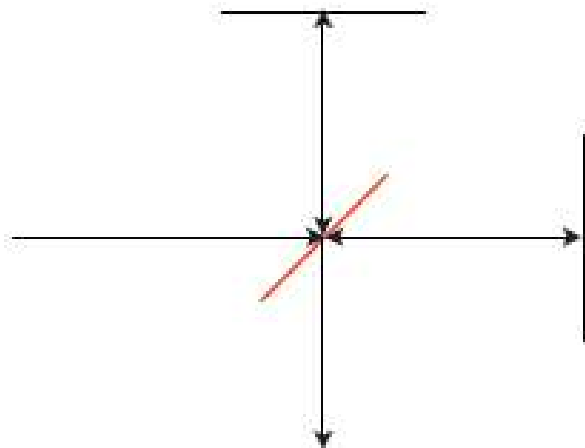
the phase can be measured by interference.

Interferometers

Mach-Zehnder

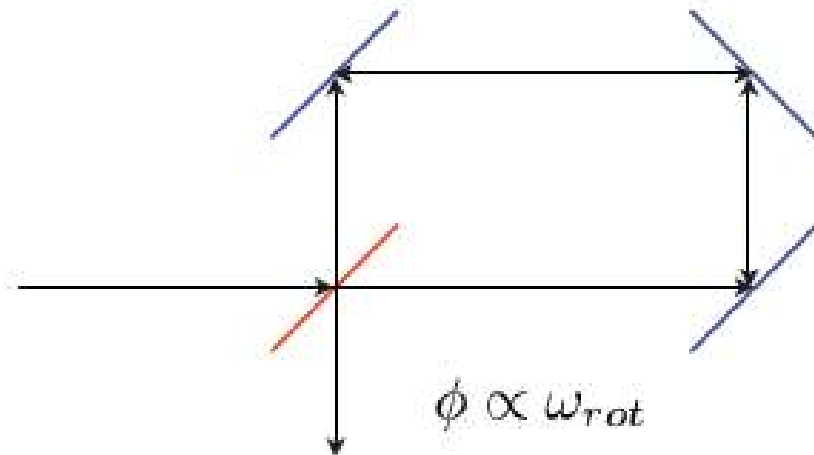


Michelson




Sagnac

laser gyroscope



laser gyro in F16

Honeywell



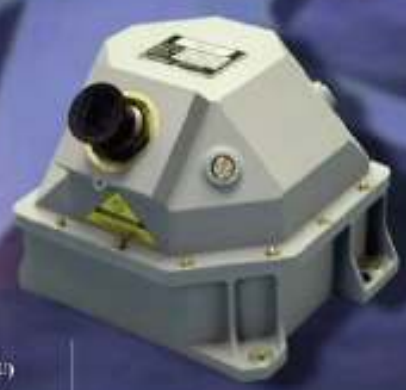

HG9900 Navigation-Grade Inertial Measurement Unit (IMU)

New
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Honeywell has fielded more than 100,000 Ring Laser Gyros (RLGs) since 1990.



System Features

- Inertial Measurement Unit (IMU)
 - ◆ Honeywell GG1320AN Digital Laser Gyros
 - ◆ Honeywell QA2800 Accelerometers
 - ◆ Honeywell Smart Inertial Electronics
 - ◆ Proven 0.5 nm/s performance
- IMU Processor
 - ◆ TI DSP TMS320VC33 (60 Mips)
 - ◆ 128 Kbytes SRAM, 512 Kbytes Flash EEPROM
- Standard I/O
 - ◆ SDLC RS-422
 - ◆ 300 Hz filtered angular rate and linear acceleration (other frequencies available)
 - ◆ 300 Hz compensated AVs and DVs (other frequencies available)
- Power Input
 - ◆ 5, ±15 Vdc input

Interference (temporal)

- superposition of two monochromatic waves of different frequency,

$$U_1(\mathbf{r}, t) = A_1 \exp(-i\mathbf{k}_1 \cdot \mathbf{r}) \exp(i\omega_1 t),$$

$$U_2(\mathbf{r}, t) = A_2 \exp(-i\mathbf{k}_2 \cdot \mathbf{r}) \exp(i\omega_2 t),$$

- at fixed \mathbf{r} ,

$$U(t) = I_1^{1/2} \exp(i\omega_1 t) + I_2^{1/2} \exp(i\omega_2 t),$$

- the intensity,

$$I(t) = I_1 + I_2 + 2(I_1 I_2)^{1/2} \cos[(\omega_2 - \omega_1)t],$$

- light beating at the frequency,

$$\mu = \frac{\omega_2 - \omega_1}{2\pi},$$

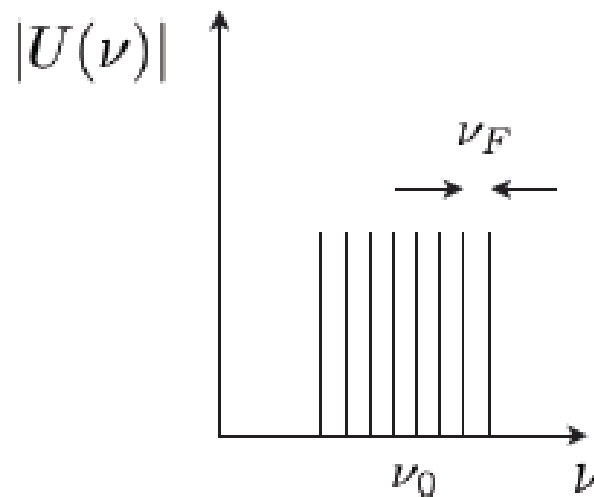
multiple-wave interference (temporal with M waves)

- equal amplitudes and equal phase differences, such as Fabry-Perot filter, Bragg filter,
- the total scalar field is thus the summation

$$U(t) = I_0^{1/2} \sum_{q=-M}^M \exp(i\omega_q t),$$

where

$$\omega_q = 2\pi\mu_q = 2\pi(\mu_0 + q\mu_F),$$



multiple-wave interference (temporal with M waves)

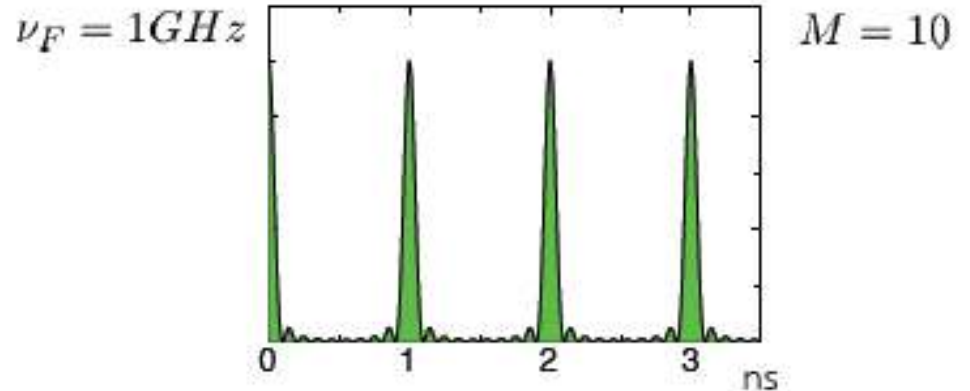
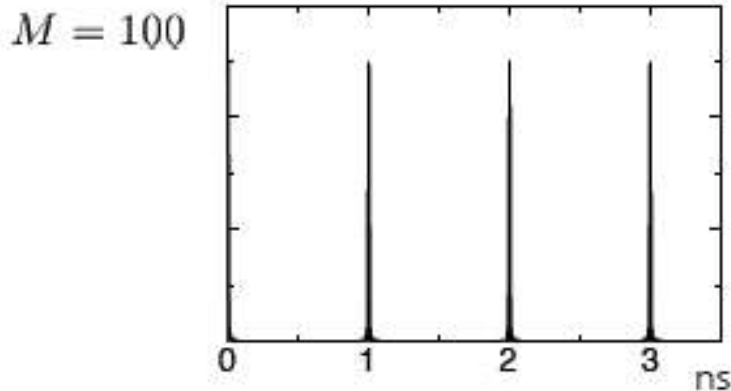
→ multiple-wave interference,

$$U(t) = I_0^{1/2} \sum_{q=-M}^M \exp(i\omega_q t),$$

→ intensity:

$$I(t) = |U(t)|^2 = I_0 \frac{\sin^2(M\pi\mu_F t)}{\sin^2 \pi\mu_F t},$$

→ acts as a high-finesse/high-Q filter,



Summary of Wave Optics

- light propagates in form of waves
- wave equation in its simplest form is linear, which gives rise to superposition and separation of time and space dependence (interference, diffraction)
- waves are characterized by wavelength and frequency
- propagation through media is characterized by refractive index n , which describes the change in phase velocity
- media with refractive index n alter velocity, wavelength and wavenumber but not frequency
- lenses alter the curvature of wavefronts

paraxial wave approximation

- paraxial wave = wavefronts normals are paraxial rays

$$U(r) = A(r)\exp(-ikz),$$

- $A(r)$ slowly varying with at a distance of λ ,
- paraxial Helmholtz equation

$$\begin{aligned}(\nabla^2 + k^2)U(r) &= 0, \\ \rightarrow \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} - 2ik\frac{\partial}{\partial z}\right)A(r) &= 0,\end{aligned}$$

- solution of the paraxial Helmholtz equation is the *Gaussian* beams,