

IPT 5260

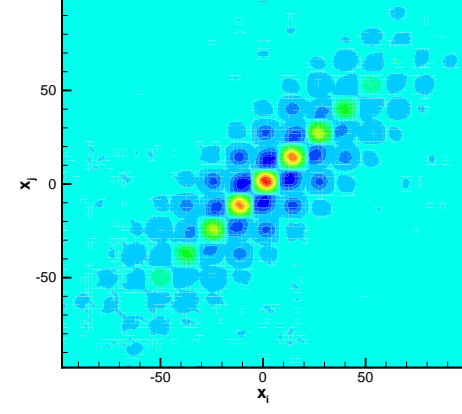
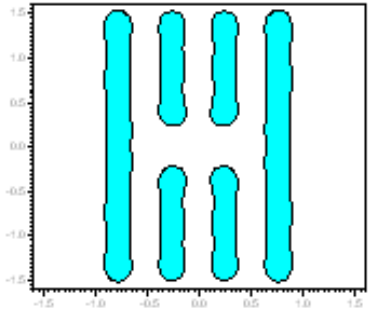
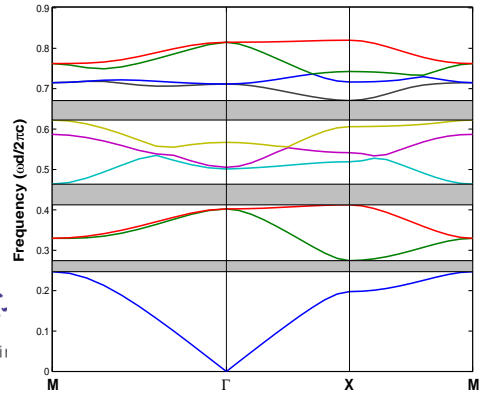
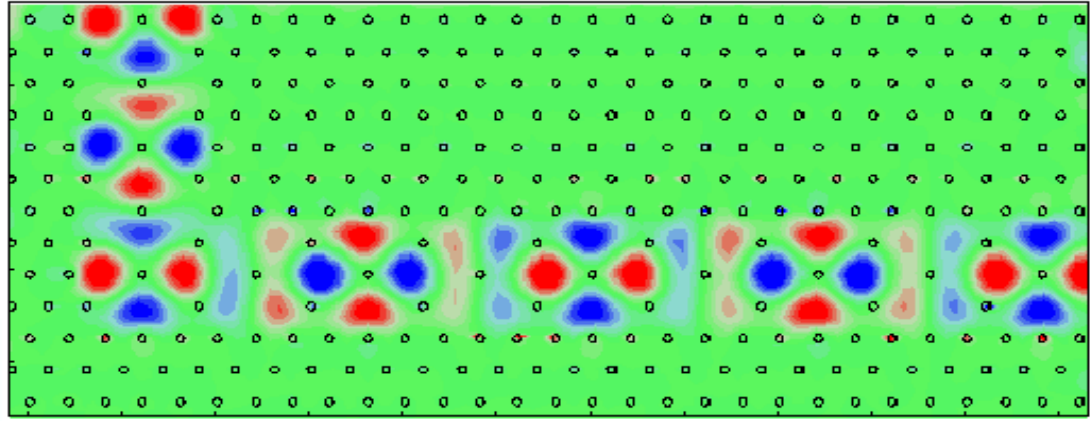
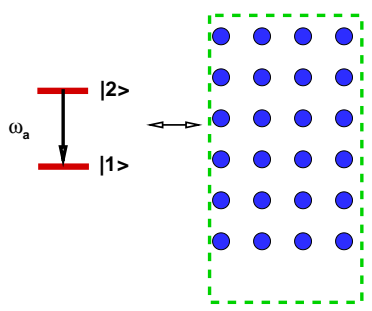
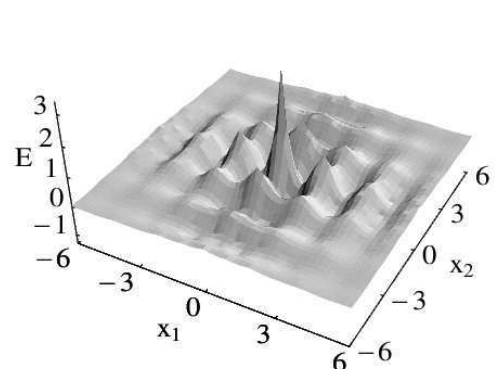
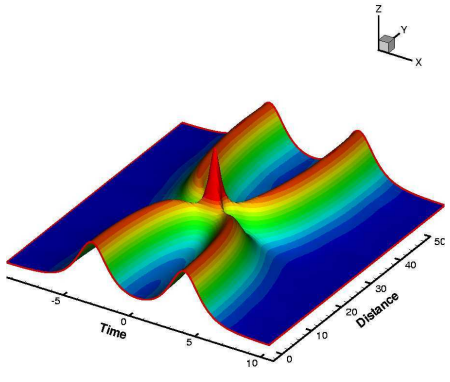
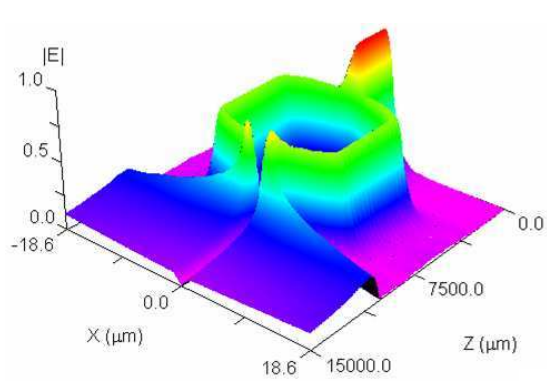
Computational Methods for Optoelectronics

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Computational Optoelectronics



IPT 52600

Time: M2M3M4 (09:30-12:30 AM, Monday)

Course Description:

- Fundamental numerical simulation techniques for solving Optoelectronics related problems.
- Taking this course you are asked to program by yourself.
- Although this course is given primarily for the first year graduate students, those who are undergraduates or senior graduates are encouraged to take this course.
- **Background:** No required but you must learn how to program in C/C++, Fortran, Matlab, Mathematica, or Mapple (at least one of these programming languages).
- **Remark:** This course is more focused on numerical methods than *numerical analysis*.

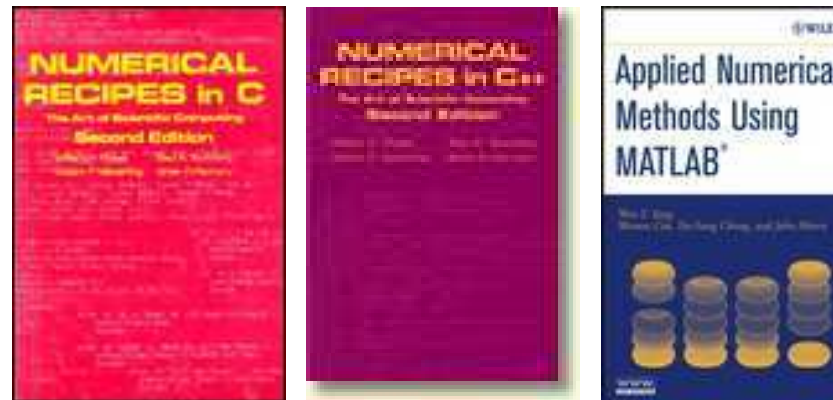
Teaching Method: in-class lectures with examples and projects studies.

Evaluation:

1. Two Homeworks, 70%;
2. One Project, 30%.

Reference Books

- W. H. Press et al., "Numerical Recipes (in C, C++, or Fortran)," Cambridge University Press (1992).
- W. Y. Yang et al., "Applied Numerical Methods Using MATLAB," Wiley (2005).



<http://www.nr.com/>

ftp://ftp.wiley.com/public/sci_tech_med/applied_numerical/

Related courses

- Data Structures
- Algorithm
- Numerical Analysis
- Numerical Partial Differential Equations
- Scientific Computation
- Computational Physics
- Computational Fluid Dynamics, CFD
- First Principle calculation, ab-initio
- Finite Element Method, FEM
- Monte Carlo/Molecular Dynamics calculation, MC/MD
- Computer-Aided Design, CAD

1, Linear Algebraic Equations

$$\mathbf{A} \cdot \mathbf{x} = \mathbf{b},$$

where

$$\mathbf{A} = \begin{bmatrix} a_{11} & a_{12} & \cdots & a_{1N} \\ a_{21} & a_{22} & \cdots & a_{2N} \\ \cdots & \cdots & \cdots & \cdots \\ a_{M1} & a_{M2} & \cdots & a_{MN} \end{bmatrix}, \mathbf{x} = \begin{bmatrix} x_1 \\ x_2 \\ \vdots \\ x_N \end{bmatrix}, \text{ and } \mathbf{b} = \begin{bmatrix} b_1 \\ b_2 \\ \vdots \\ b_N \end{bmatrix}.$$

- ➔ Gauss-Jordan elimination
- ➔ LU decomposition
- ➔ Jacobi iteration
- ➔ Gauss-Seidel iteration

Multiple scattering method for PBG

Helmholtz equation:

$$\nabla^2 E + \tilde{k}^2(M)E = 0$$

with

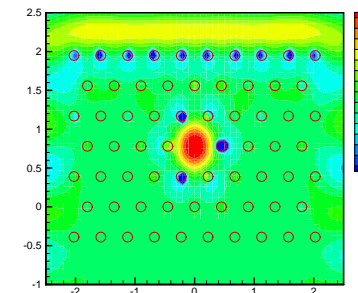
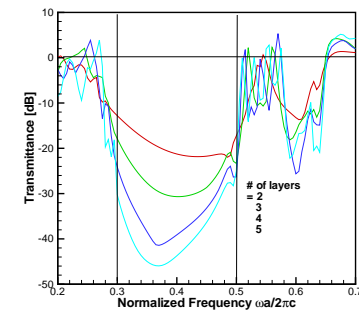
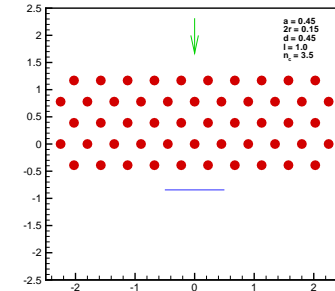
$$\tilde{k}^2(M) = k^2 \tilde{\epsilon} = \begin{cases} k^2 \epsilon_j & \text{if } M \in C_j (j = 1, 2, \dots, N) \\ k^2 & \text{if } M \notin C_j (j = 1, 2, \dots, N) \end{cases}$$

write down the total field in decomposition,

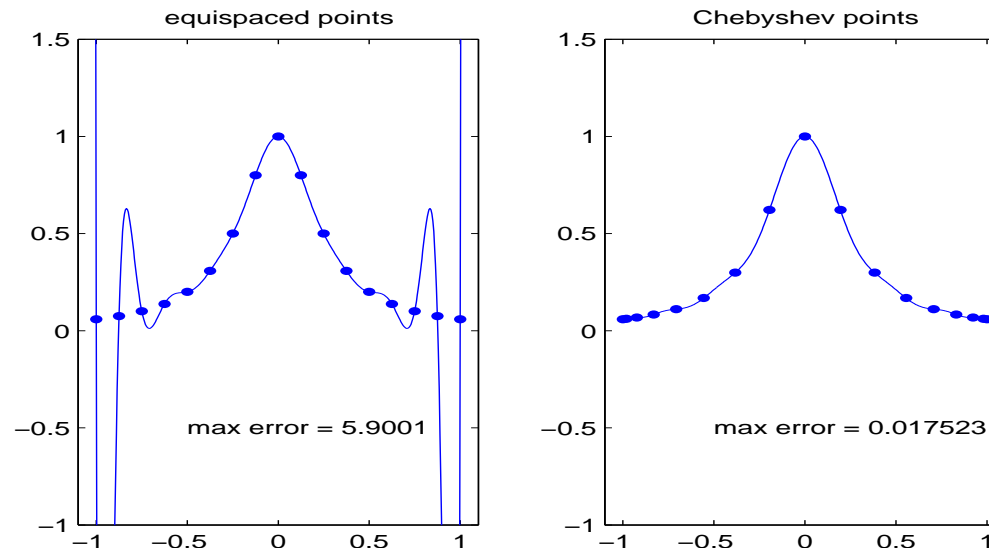
$$E(P) = \sum_{m=-\infty}^{\infty} a_{l,m} J_m[kr_l(P)] e^{im\theta_l(P)} + \sum_{m=-\infty}^{\infty} b_{l,m} H_m^{(1)}[kr_l(P)] e^{im\theta_l(P)},$$

then all we need is to solve

$$\hat{\mathbf{b}}_l - \sum_{j \neq l} \mathbf{S}_l \mathbf{T}_{l,j} \hat{\mathbf{b}}_j = \mathbf{S}_l \mathbf{Q}_l.$$



2, Interpolation, Curve Fitting, and Integration



- ➔ Polynomial interpolation and extrapolation
- ➔ Rational function interpolation
- ➔ Chebyshev approximation
- ➔ Padé approximation
- ➔ Fast Fourier Transform and Discrete Fourier Transform
- ➔ Trapezoidal and Simpson method for integration

Finite difference approximation

Second-order FD approximation for $u'(x_j)$

$$u'(x_j) \approx \frac{u_{j+1} - u_{j-1}}{2h}$$

in the matrix-vector form (with periodic boundary)

$$\begin{pmatrix} u'_1 \\ \vdots \\ u'_j \\ \vdots \\ u'_N \end{pmatrix} = h^{-1} \begin{pmatrix} 0 & \frac{1}{2} & & & -\frac{1}{2} \\ -\frac{1}{2} & 0 & & & \\ & & \ddots & & \\ & & & -\frac{1}{2} & 0 & \frac{1}{2} \\ & & & & \ddots & \\ & & & & & 0 & \frac{1}{2} \\ \frac{1}{2} & & & & & -\frac{1}{2} & 0 \end{pmatrix} \begin{pmatrix} u_1 \\ \vdots \\ u_j \\ \vdots \\ u_N \end{pmatrix}$$

3, Ordinary Differential Equations

$$\frac{d^2y}{dx^2} + q(x)\frac{dy}{dx} = r(x),$$

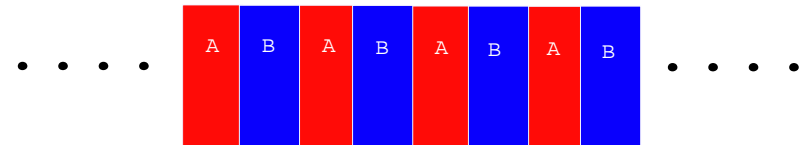
- ➔ Euler's method
- ➔ Runge-Kutta method
- ➔ Adaptive stepsize control
- ➔ Predictor-corrector method
- ➔ Boundary value problems
- ➔ Relaxation method

ODE with B. C.: 1D Bragg reflector

coupled-mode equation:

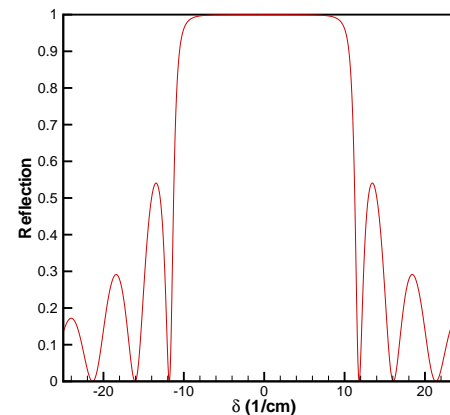
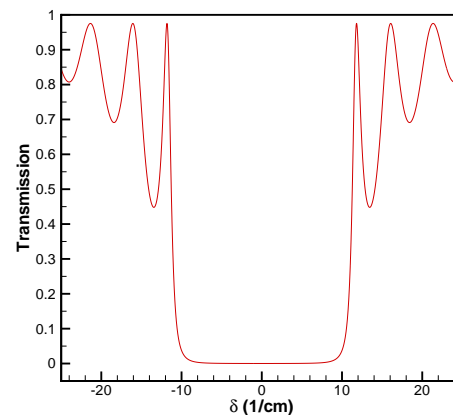
$$\frac{dE_+(z)}{dz} = i\delta E_+(z) + i\kappa E_-(z)$$
$$\frac{dE_-(z)}{dz} = i\delta E_-(z) + i\kappa E_+(z)$$

with the Boundary Condition:



$$E_+(z = 0) = 1$$

$$E_-(z = L) = 0$$

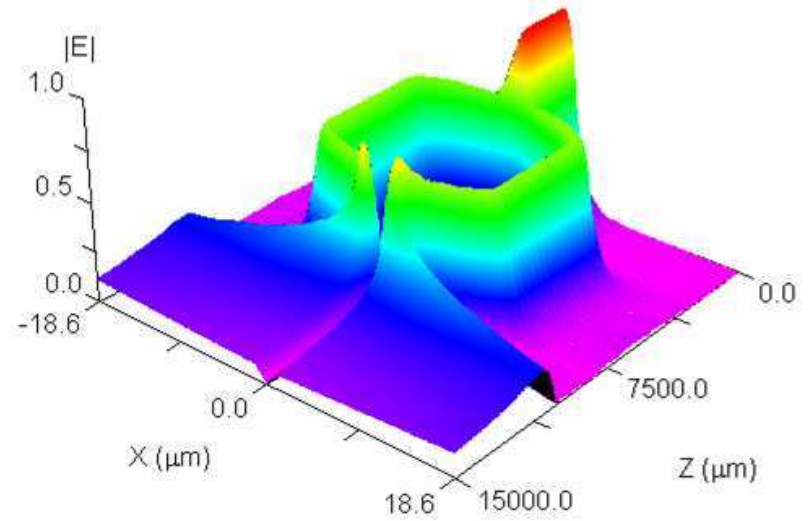
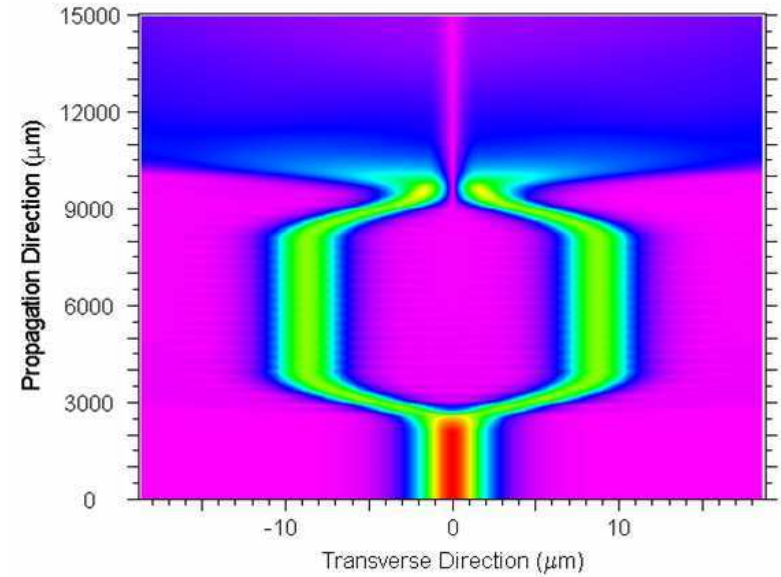
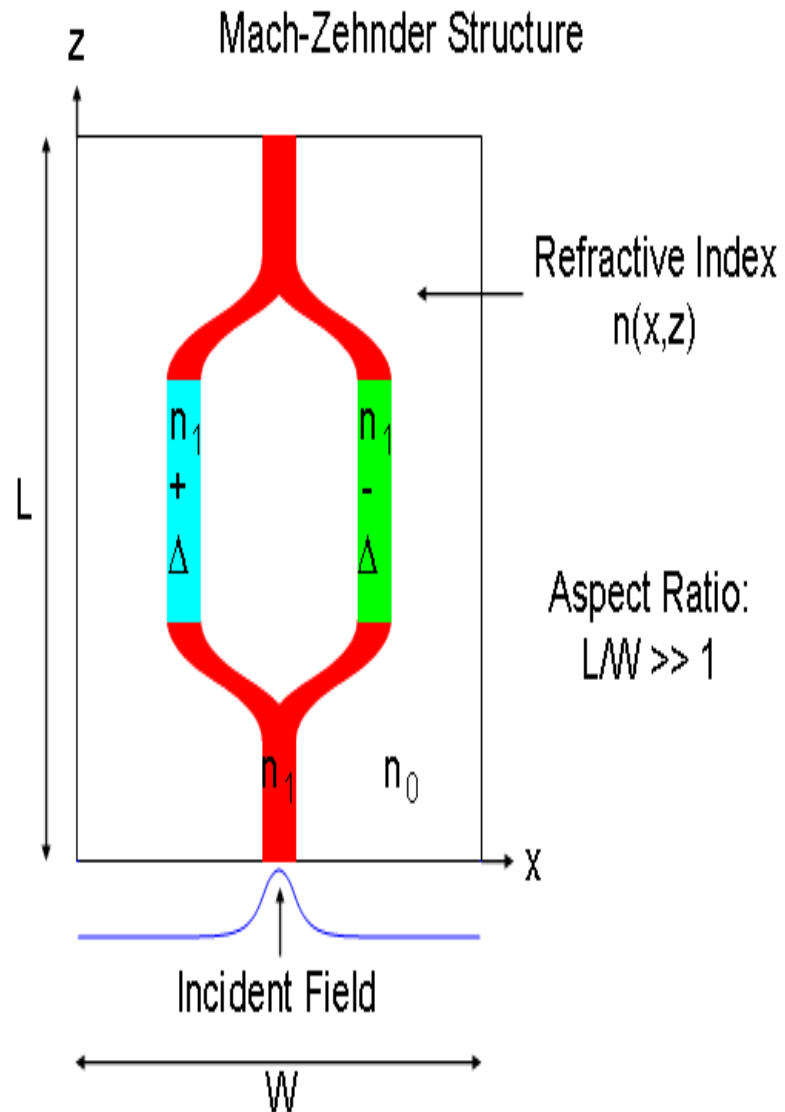


4, Partial Differential Equations

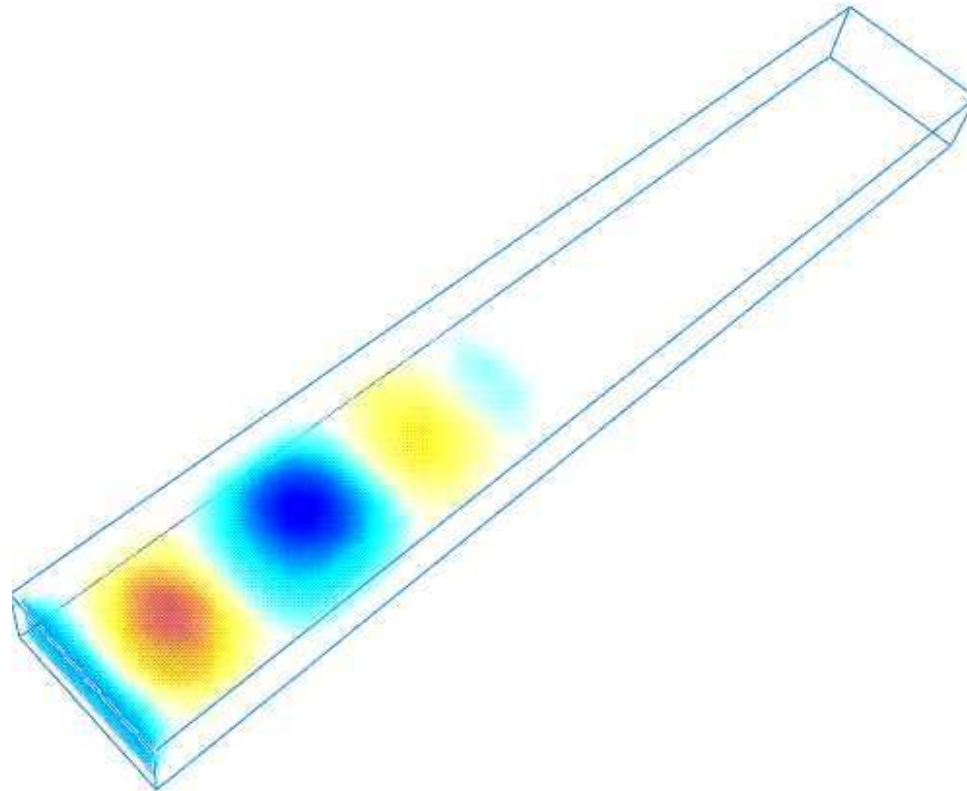
$$A(x, y) \frac{\partial^2 u}{\partial x^2} + B(x, y) \frac{\partial^2 u}{\partial x \partial y} + C(x, y) \frac{\partial^2 u}{\partial y^2} = f(x, y, u, \frac{\partial u}{\partial x}, \frac{\partial u}{\partial y}),$$

- Euler method
- Crank-Nicholson method
- Beam Propagation Method
- Slit-Step Fast Fourier Transform
- Finite Difference Time Domain method
- Absorption Boundary Condition
- Periodic Boundary Condition
- Perfect Matching Layers

Mach-Zehnder structure



Metallic Waveguide



5, Nonlinear Equations and Nonlinear PDE

Nonlinear equation:

$$f(x) = x^3 + x^2 = 5,$$

- ➔ Iterative method
- ➔ Newton-Raphson method
- ➔ Secant Method

Nonlinear Schrödinger equation:

$$i \frac{\partial U}{\partial t} + \frac{\partial^2 U}{\partial x^2} + |U|^2 U = 0$$

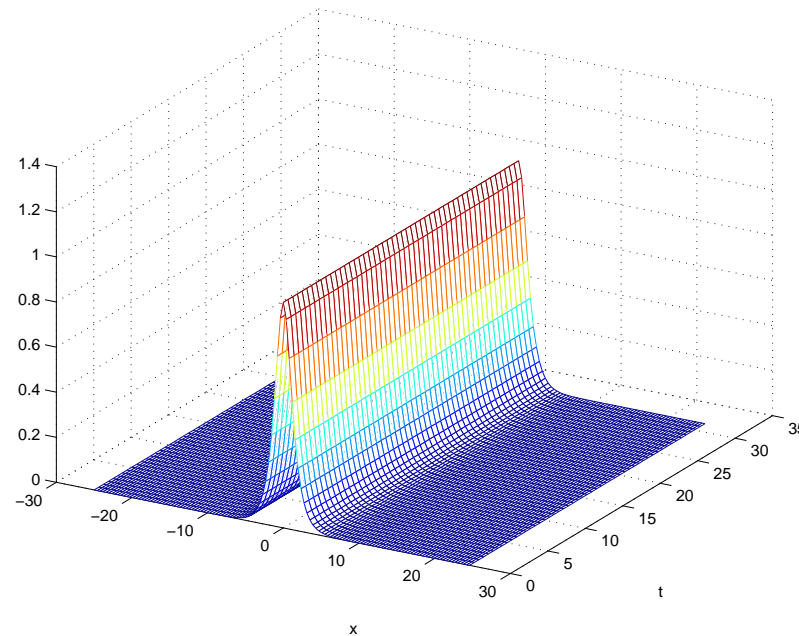
- ➔ Crank-Nicholson method
- ➔ Split-Step Fast Fourier Transform
- ➔ Pseudospectral method

Propagation of solitons

Nonlinear Schrödinger equation:

$$i\frac{\partial U}{\partial t} + \frac{\partial^2 U}{\partial x^2} + |U|^2 U = 0$$

supports soliton solutions, $U(t = 0, x) = \text{sech}(x)$.



simulated by Fourier spectral + 4th-order explicit Runge-Kutta methods,

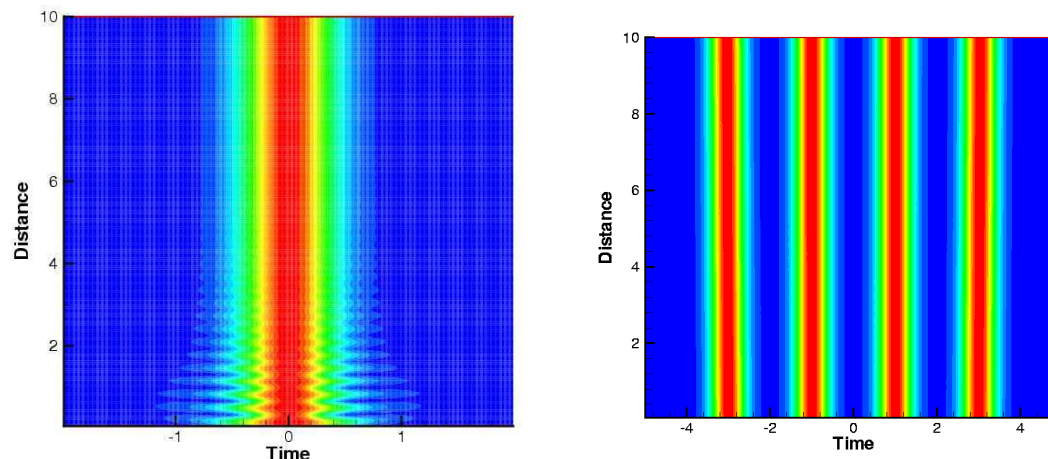
Bound solitons in CGLE

Complex Ginzburg-Landau Equation:

$$iU_z + \frac{D}{2}U_{tt} + |U|^2U = i\delta U + i\epsilon|U|^2U + i\beta U_{tt} \\ + i\mu|U|^4U - \nu|U|^4U,$$

seek for bound-state solutions by **propagation** method.

$$U(z, t) = \sum^N U_0(z, t + \rho_j) e^{i\theta_j}$$

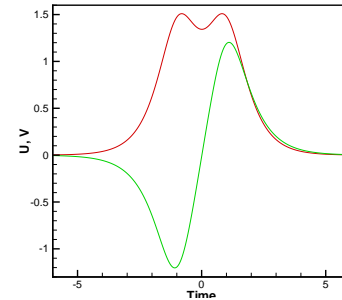


Vector bound solitons

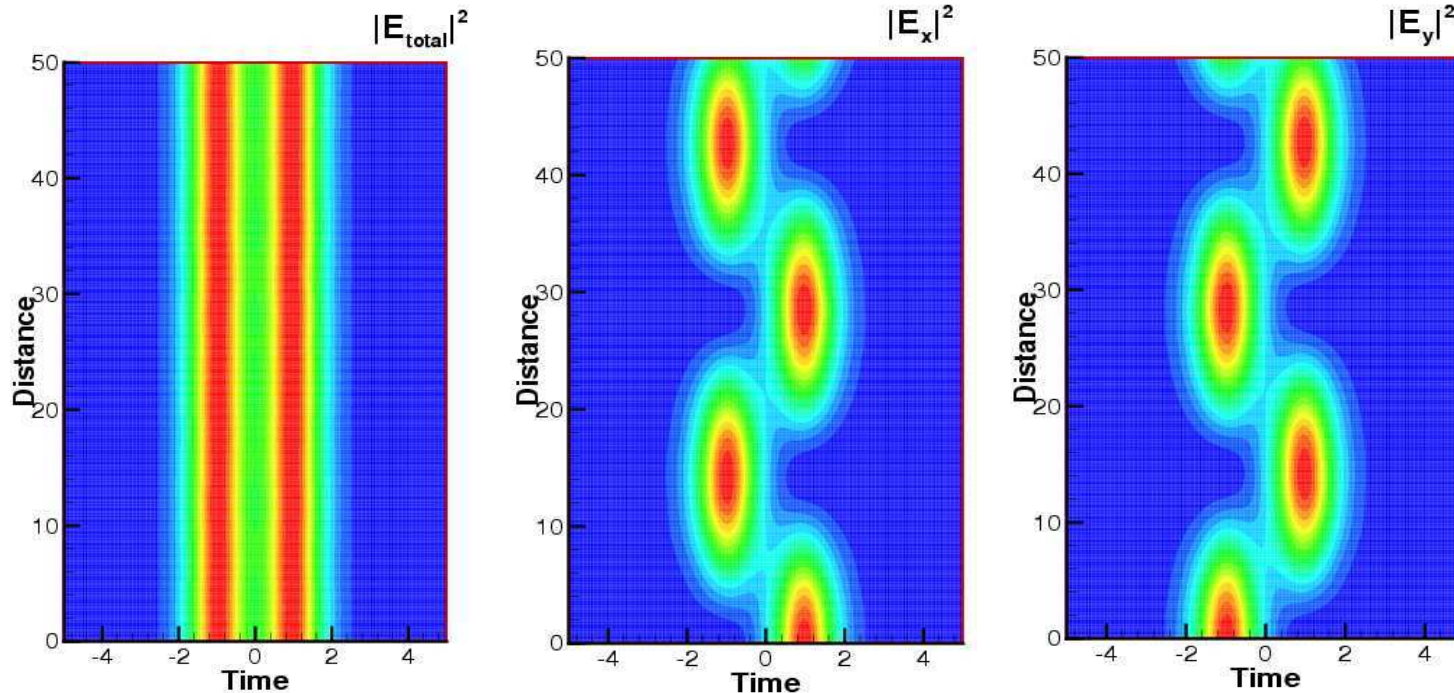
Coupled Nonlinear Schrödinger Equations:

$$i\frac{\partial U}{\partial z} + \frac{1}{2}\frac{\partial^2 U}{\partial t^2} + A|U|^2U + B|V|^2U = 0$$

$$i\frac{\partial V}{\partial z} + \frac{1}{2}\frac{\partial^2 V}{\partial t^2} + A|V|^2V + B|U|^2V = 0$$



where $A = 1/3$, $B = 2/3$; and U, V are circular polarization fields.

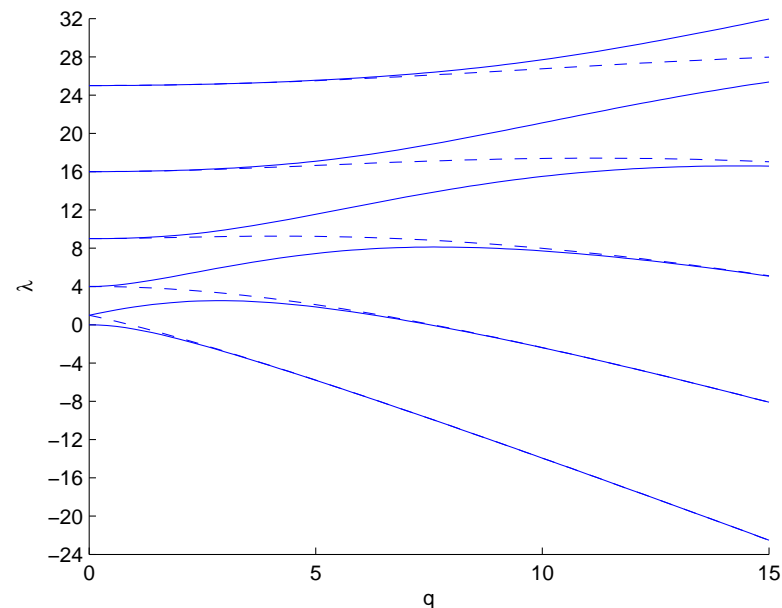


6, Eigenvalues and Eigenvectors

$$\mathbf{A} \cdot \mathbf{x} = \lambda \mathbf{x},$$

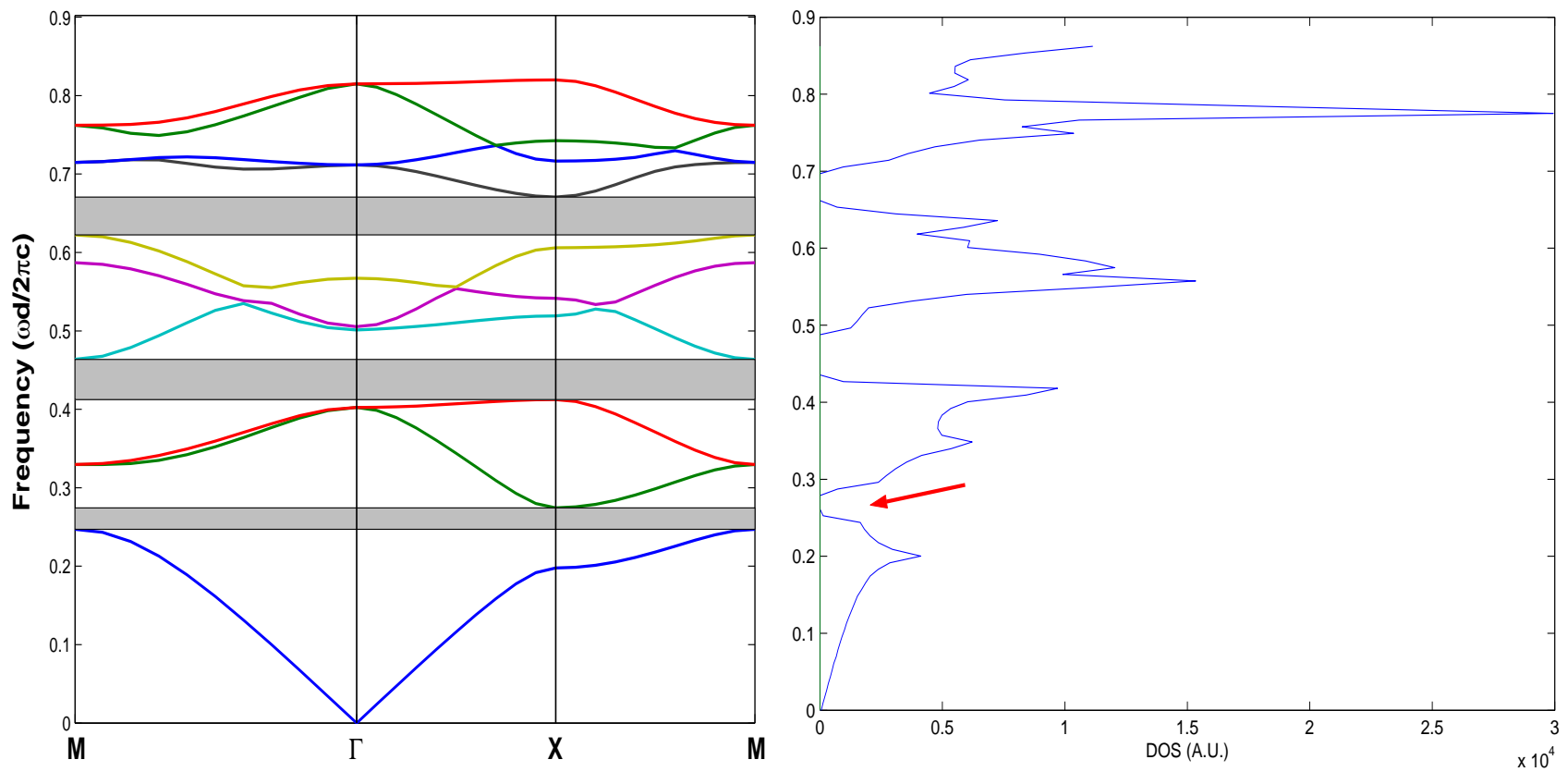
- ➔ Similarity transformation and Diagonalization
- ➔ Jacobi method
- ➔ The QR algorithm

Mathieu equation: $-u_{xx} + 2q \cos(2x)u = \lambda x,$



Band diagram and Density of States

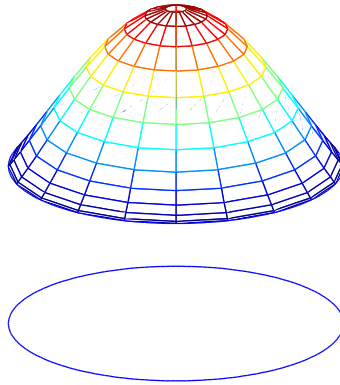
$$\frac{1}{\epsilon(\mathbf{r})} \nabla \times \{ \nabla \times \mathbf{E}(\mathbf{r}) \} = \frac{\omega^2}{c^2} \mathbf{E}(\mathbf{r}),$$



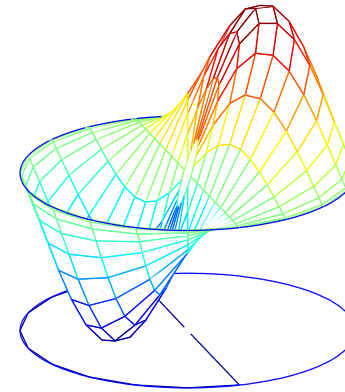
Laplacian equation in a disk

Eigenmodes of Laplacian equations, $[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}]u(x, y) = -\lambda f(x, y)$.

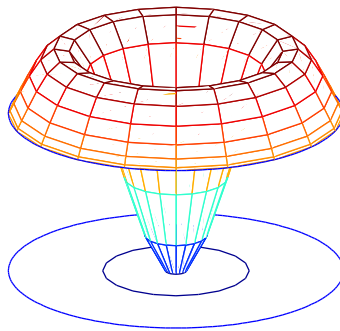
Mode 1
 $\lambda = 1.0000000000$



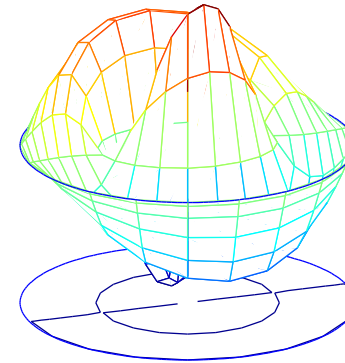
Mode 3
 $\lambda = 1.5933405057$



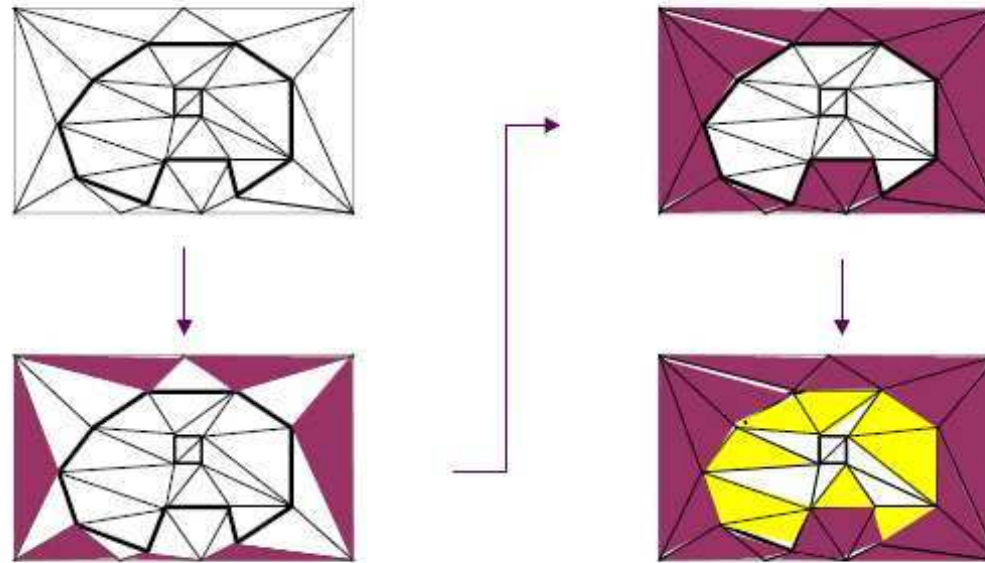
Mode 6
 $\lambda = 2.2954172674$



Mode 10
 $\lambda = 2.9172954551$

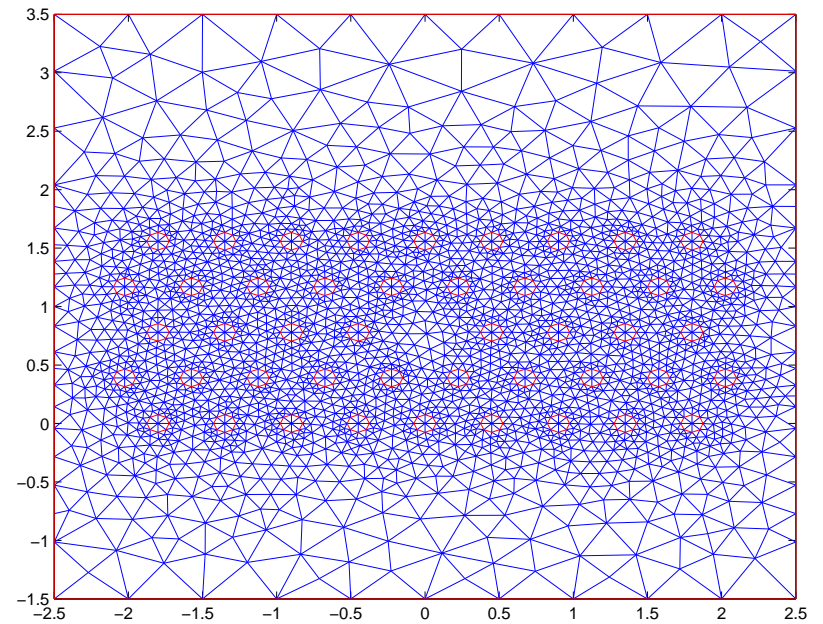
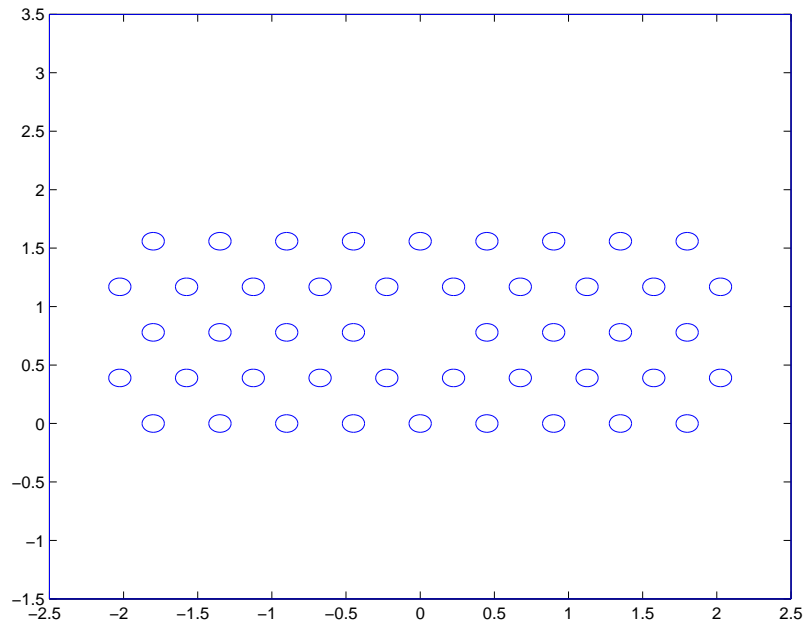


7, Finite Element Method

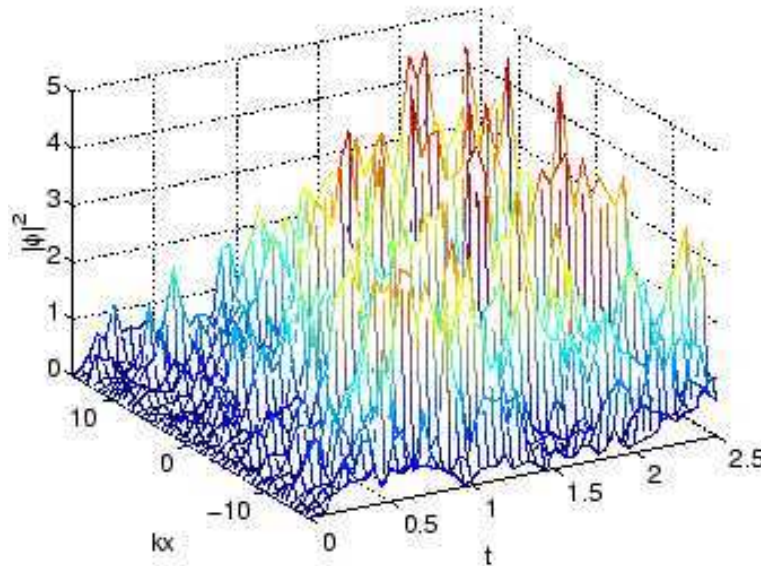


- Elements
- Mesh generation
- Element assembly
- Boundary condition

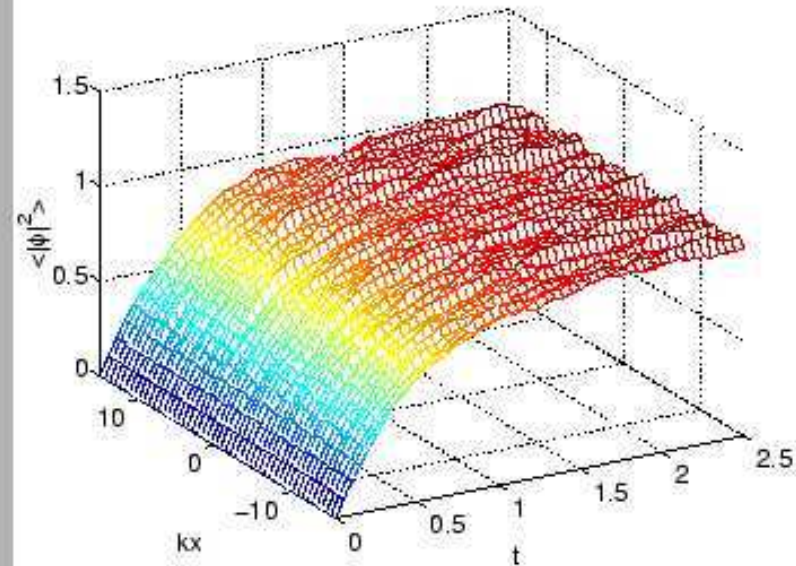
Photonic Crystals



8, Monte Carlo Method



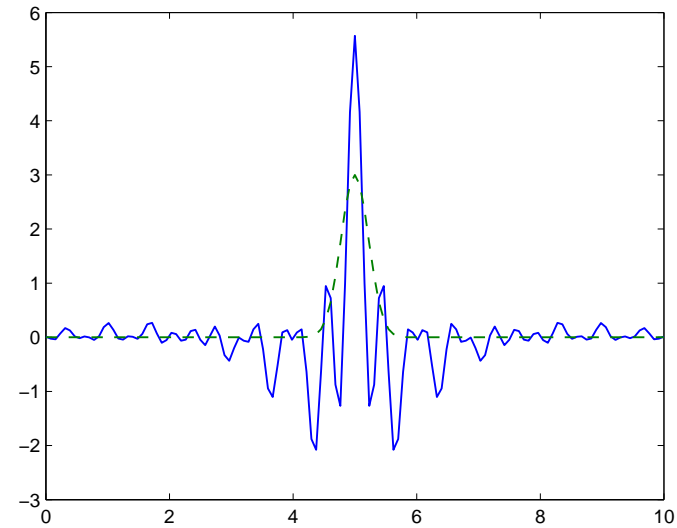
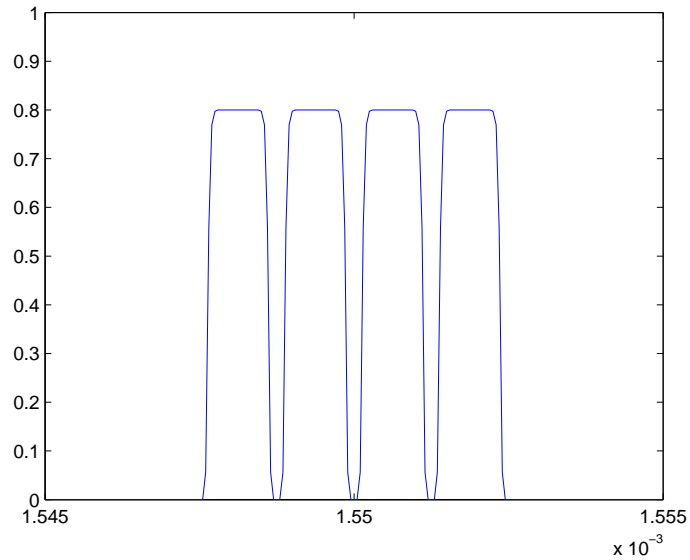
a) Single path



b) 1024 path mean

- Random numbers with uniform deviates
- Transformation method
- Rejection method
- Random bits
- Monte Carlo methods

9, Optimization



- ➔ Simulated annealing
- ➔ Genetic algorithm
- ➔ Penalty function
- ➔ Optimal control method
- ➔ Matlab built-in routines

10, Case studies

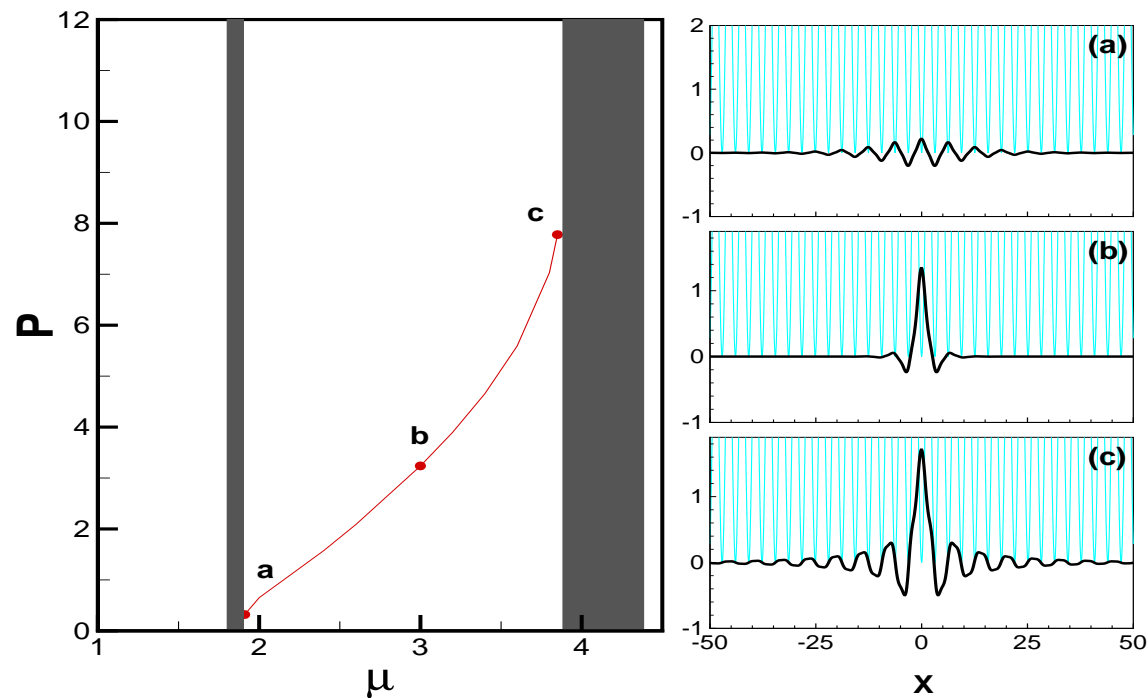
- 2D Gross-Pitaevskii equation
- Moment method
- ... your project
- ... your project
- Numerical Libraries
 - Matlab
 - IMSL
- Programming in Unix
- GNU Make and Concurrent Version System (CVS)
- Parallel programming

Eigenfunction of Nonlinear PDF

Gap solitons in optical lattices,

$$i\hbar \frac{\partial}{\partial t} \Phi_0(t, x) = -\frac{1}{2} \frac{\partial^2}{\partial x^2} \Phi_0(t, x) + V(x) \Phi_0(t, x) + g_{1D} |\phi_0(t, x)|^2 \phi_0(t, x)$$

which has gap soliton solutions.



Syllabus

1. Linear Algebraic Equations (Sep. 17th)
2. Interpolation, Curve Fitting, and Integration (Oct. 1st)
3. Ordinary Differential Equations (Oct. 8th, 15th)
Homework # 1: two-weeks to finish, (deadline: Oct. 22th)
4. Partial Differential Equations (Oct. 22th, 29th)
5. Nonlinear Equations and Nonlinear PDE (Nov. 5th, 12th)
Homework # 2: two-weeks to finish (deadline: Nov. 19th),
6. Eigenvalues and Eigenvectors (Nov. 19th)
7. Finite Element Method (Nov. 26th)
8. Monte Carlo Method (Dec. 10th)
9. Optimization (Dec. 17th)
Project: one-month to finish (deadline: Dec. 31th),
10. Case studies (Dec. 24th, 31th)

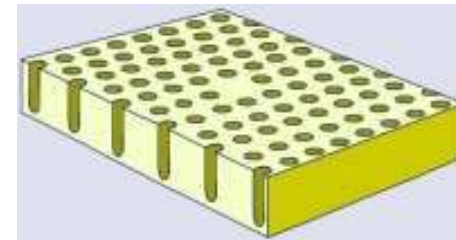
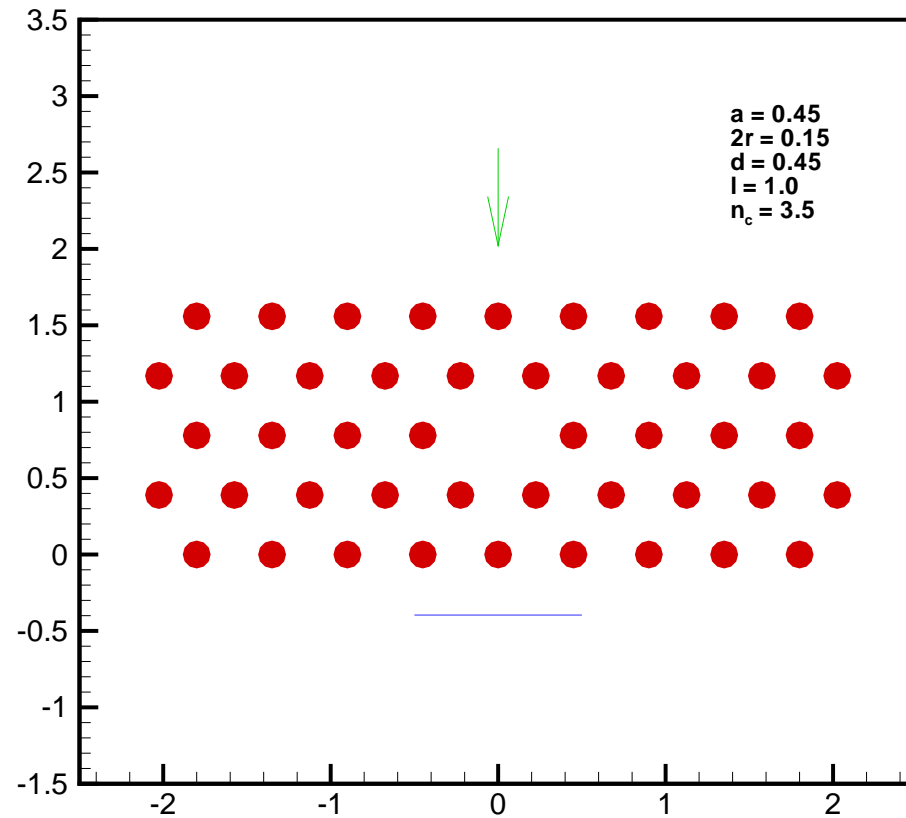
Evaluation

1. Two Homeworks (assigned), 70%
 - ➔ HW1, Ordinary Differential Equations
 - ➔ HW2, Nonlinear Partial Differential Equations
2. One Project (chosen one that is related to your research), 30%
 - ➔ Two/Three-dimensional PDE
 - ➔ Nonlinear/Coupled ODE/PDE
 - ➔ Finite Element Method
 - ➔ Monte Carlo Method
 - ➔ Optimization
 - ➔ ...
 - ➔ other suggestions

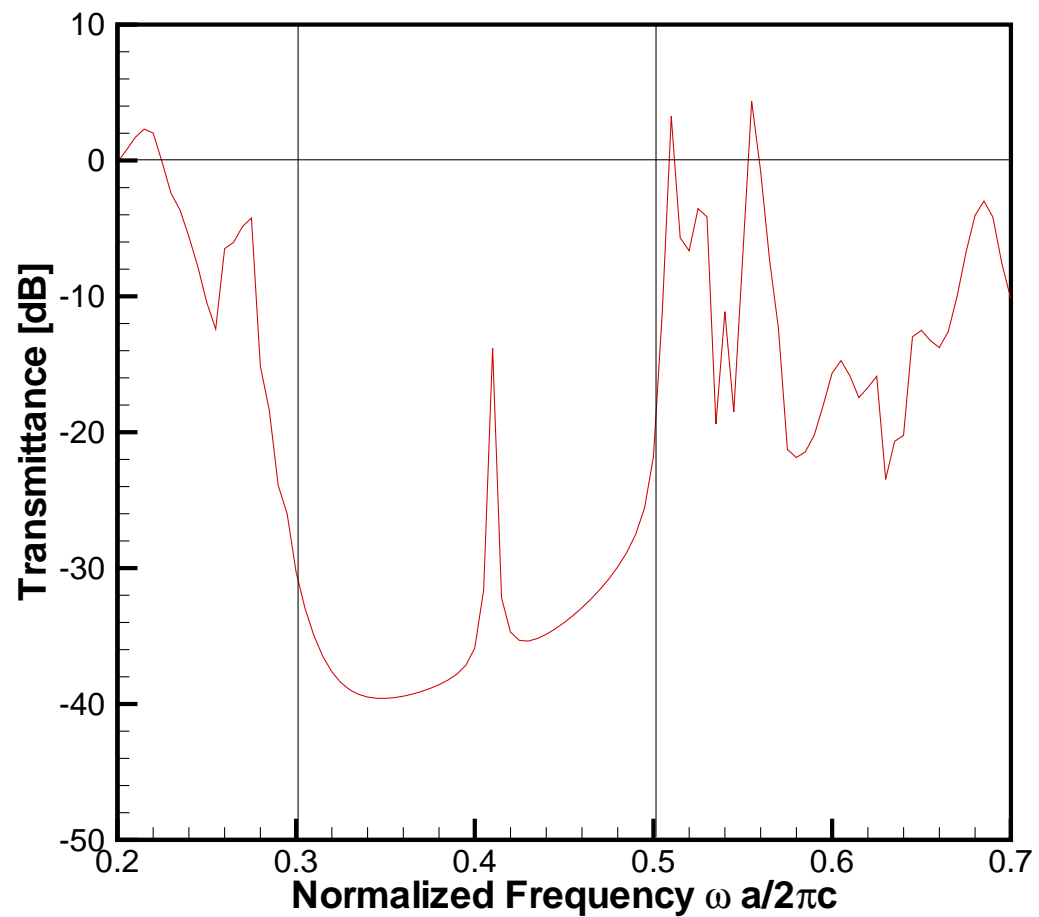
Course Projects

- 2D/3D Finite-Difference Time-Domain method
- Band-spectrum for 1D nonlinear Schrödinger equation
- Band-spectrum for 2D/3D Maxwell equation
- Coupled nonlinear PDE in 2+1 dimensions
- Soliton solutions for nonlinear PDE
- Finite Element method
- Optimization problem
- Monte-Carlo simulation

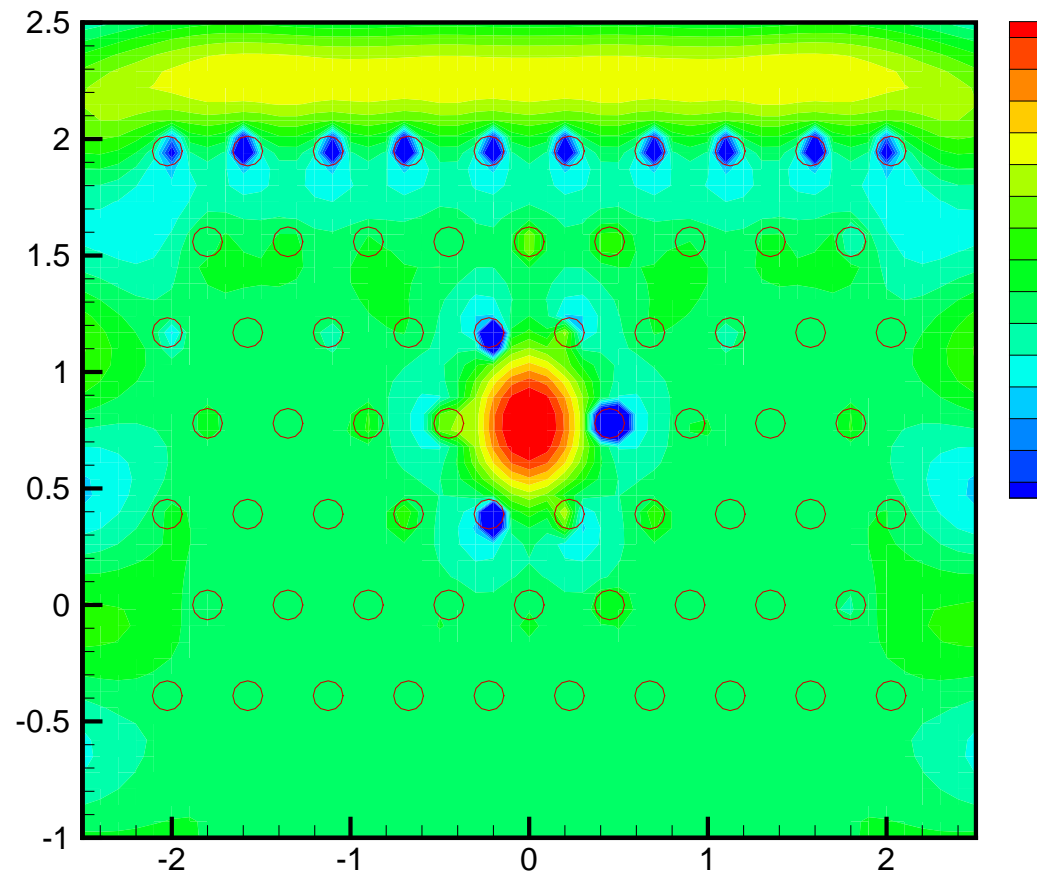
One-defect in photonic crystal



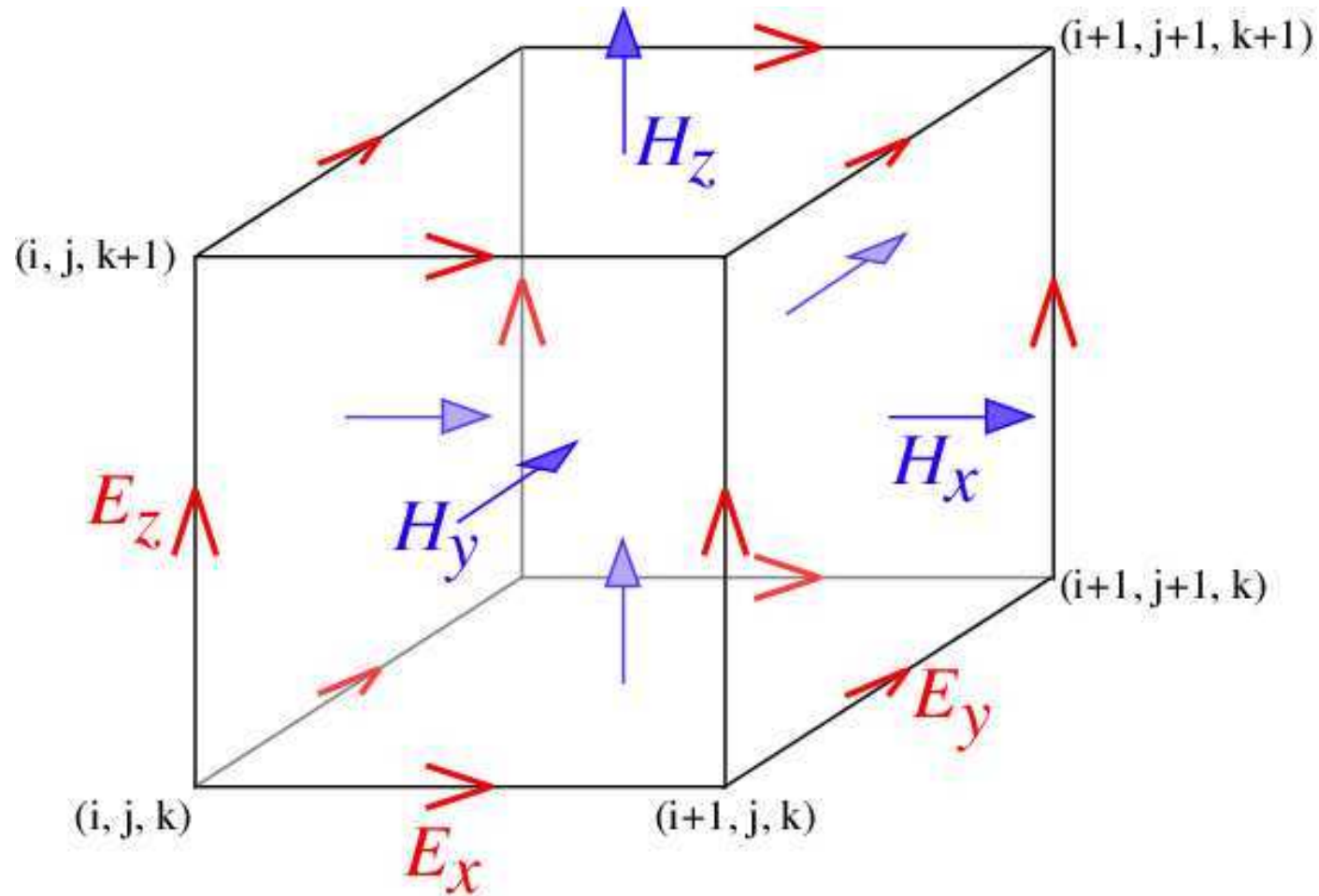
Transmittance: one-defect



Localized field: one-defect



Finite-Difference Time-Domain, FDTD: Yee's algorithm



Course Project, Spring 2006

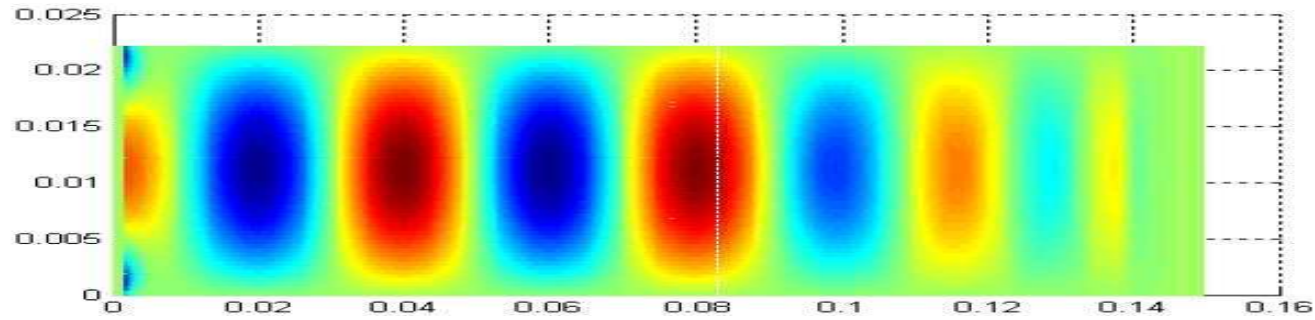


Figure 1: H_x 在 xz 平面

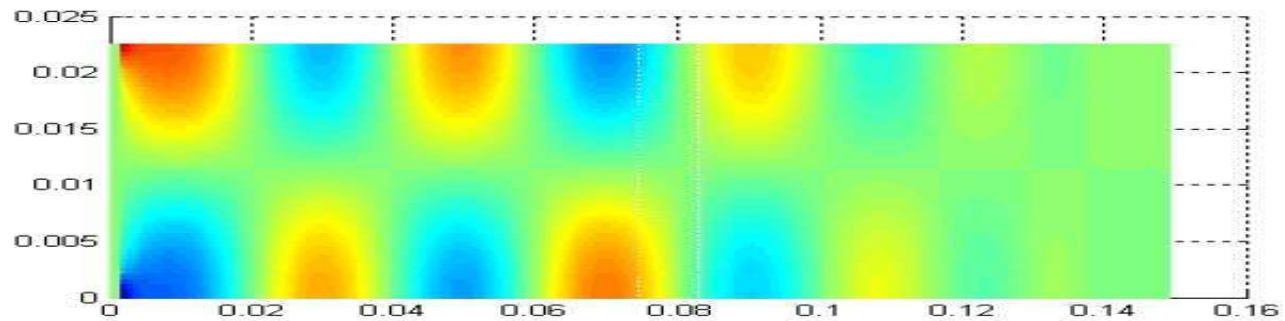


Figure 2: H_z 在 xz 平面

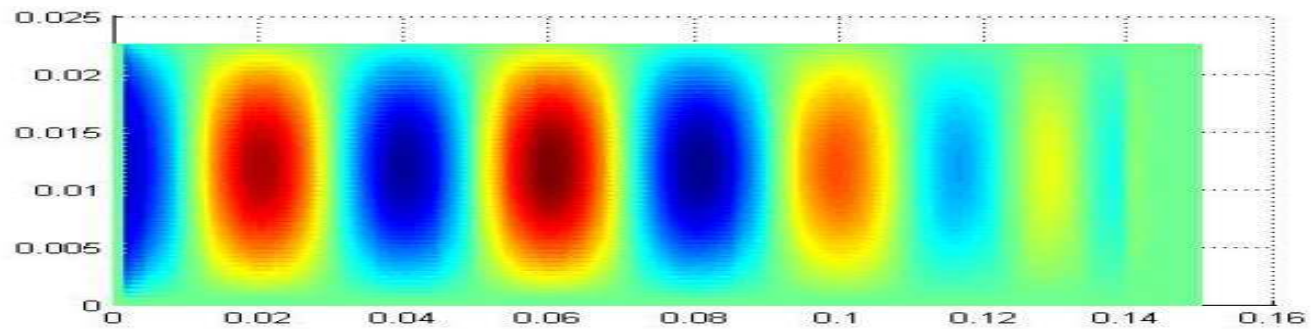
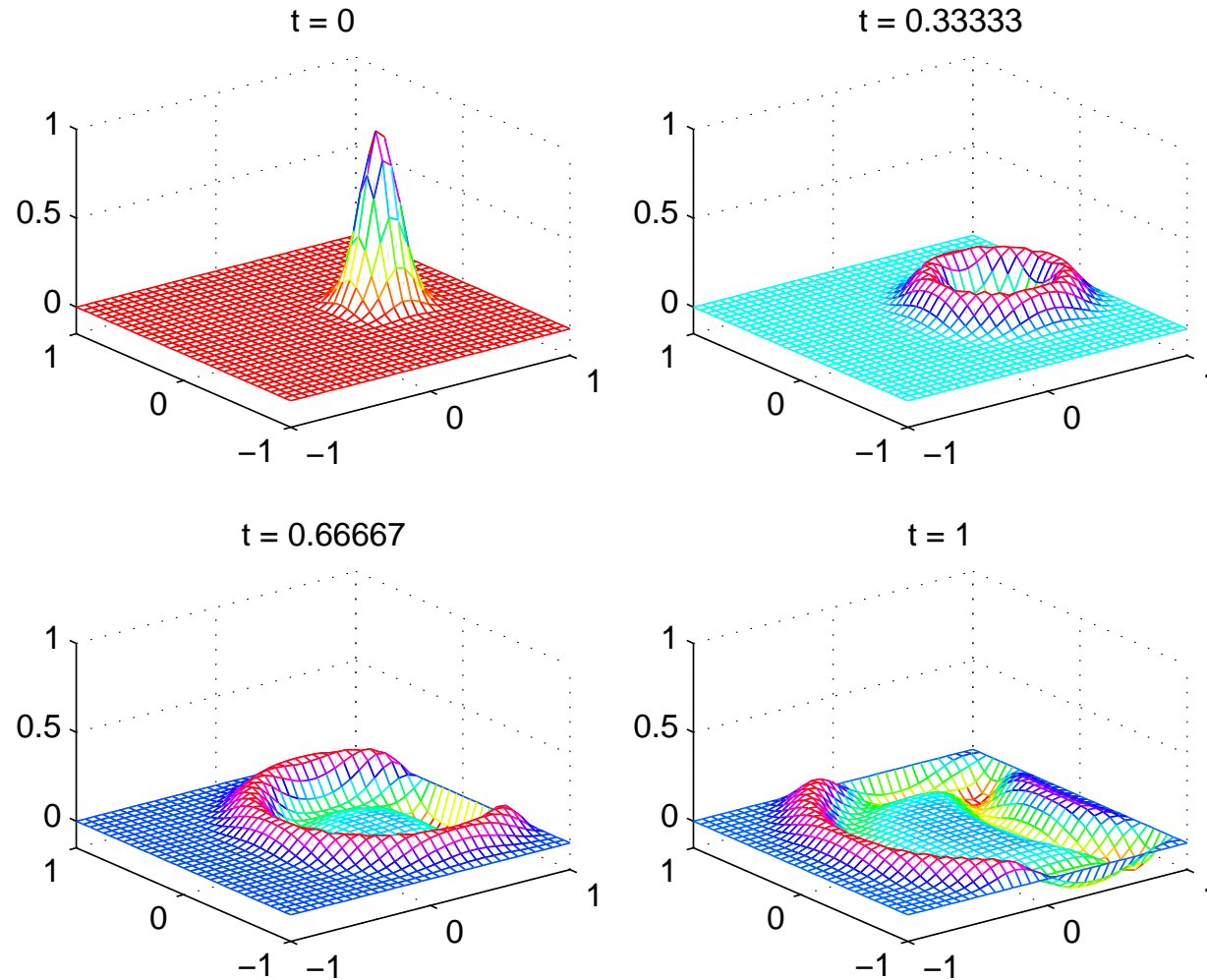


Figure 3: E_y 在 xz 平面

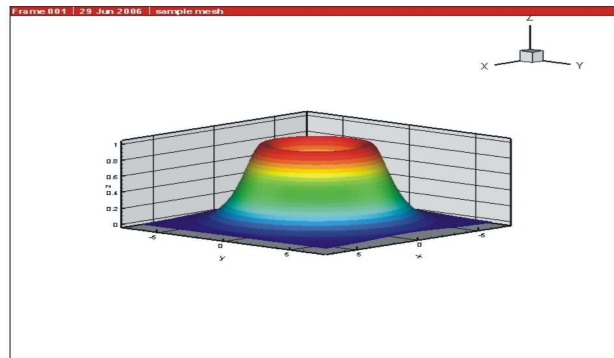
FFT method for wave equation

$$u_{tt} = u_{xx} + u_{yy}, \quad -1 < x, y < 1, \quad t > 0, \quad u = 0 \quad \text{on the boundary}$$

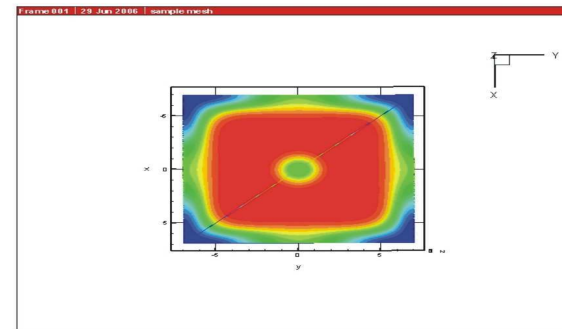


Course Project, Spring 2006

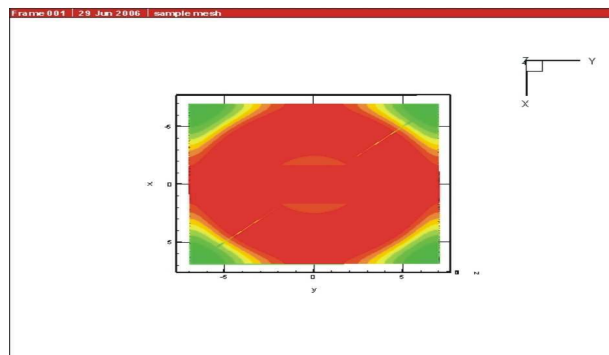
2-D sine-Gordon solitons,



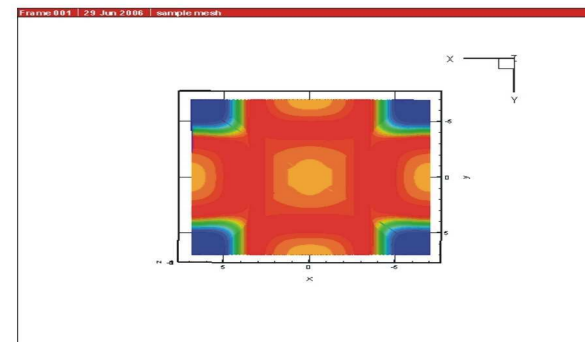
$T=0$



$T=5.6$

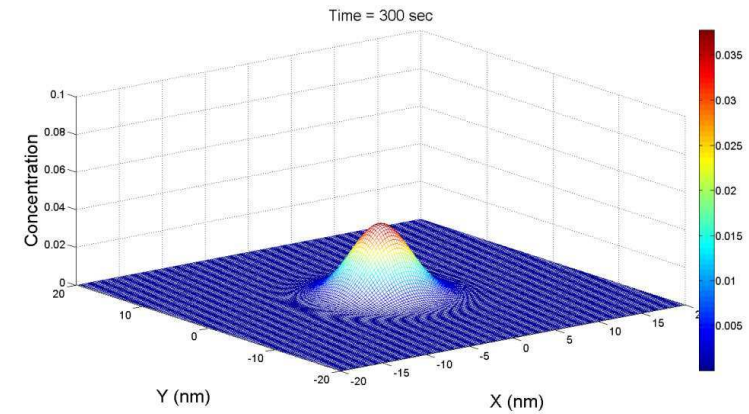
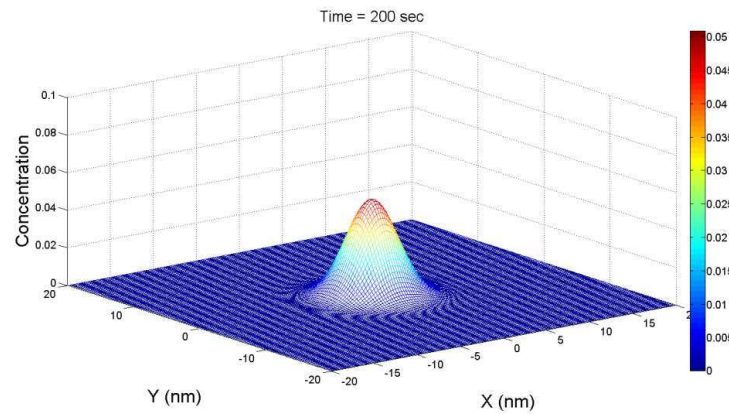
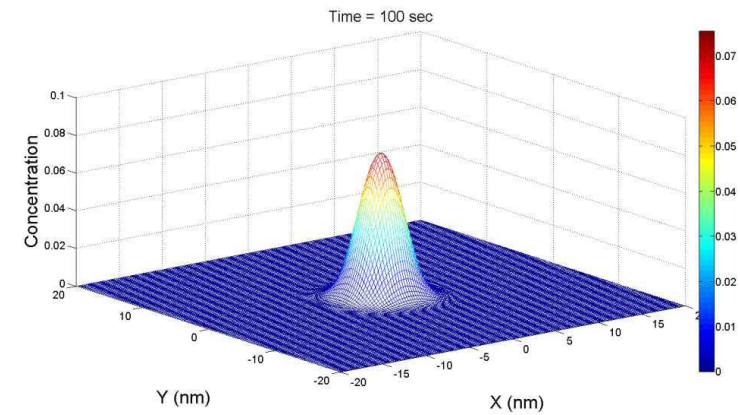
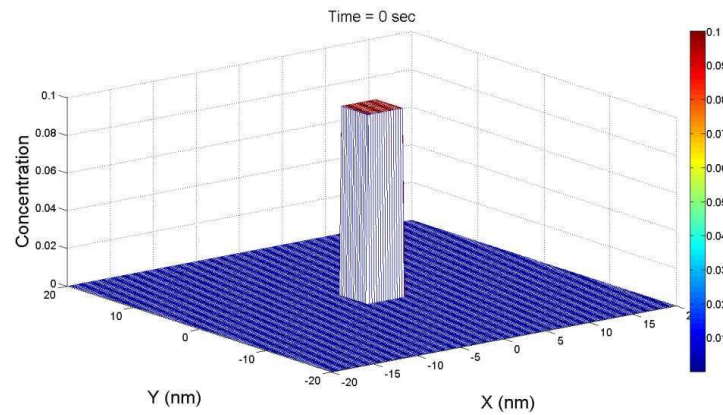


$T=2.8$



$T=8.1$

Diffusion of dopant, Spring 2007



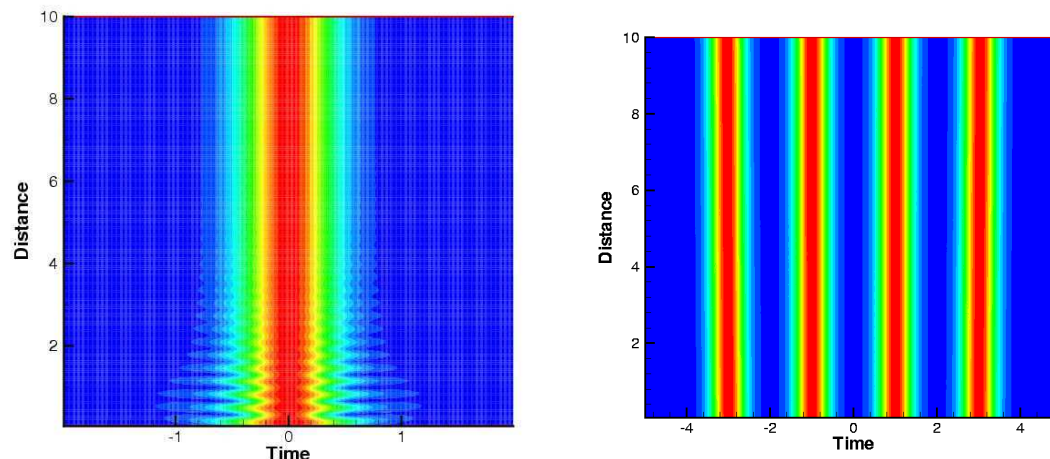
Bound solitons in CGLE

Complex Ginzburg-Landau Equation:

$$iU_z + \frac{D}{2}U_{tt} + |U|^2U = i\delta U + i\epsilon|U|^2U + i\beta U_{tt} + i\mu|U|^4U - \nu|U|^4U,$$

seek for bound-state solutions by **propagation** method.

$$U(z, t) = \sum^N U_0(z, t + \rho_j) e^{i\theta_j}$$

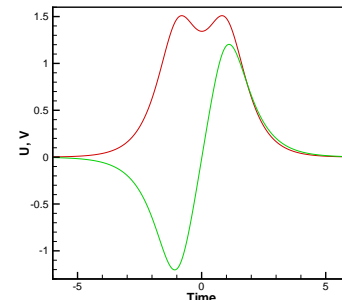


Vector bound solitons

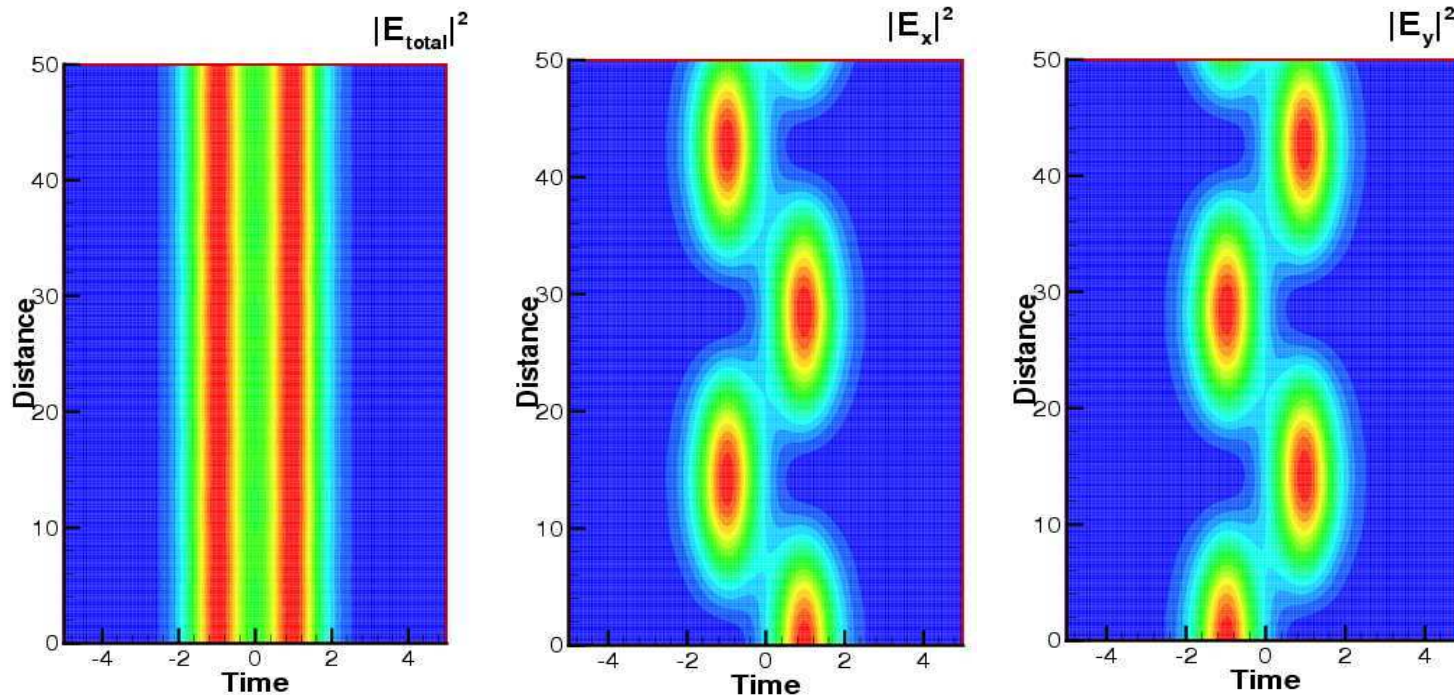
Coupled Nonlinear Schrödinger Equations:

$$i \frac{\partial U}{\partial z} + \frac{1}{2} \frac{\partial^2 U}{\partial t^2} + A|U|^2 U + B|V|^2 U = 0$$

$$i \frac{\partial V}{\partial z} + \frac{1}{2} \frac{\partial^2 V}{\partial t^2} + A|V|^2 V + B|U|^2 V = 0$$



where $A = 1/3$, $B = 2/3$; and U, V are circular polarization fields.

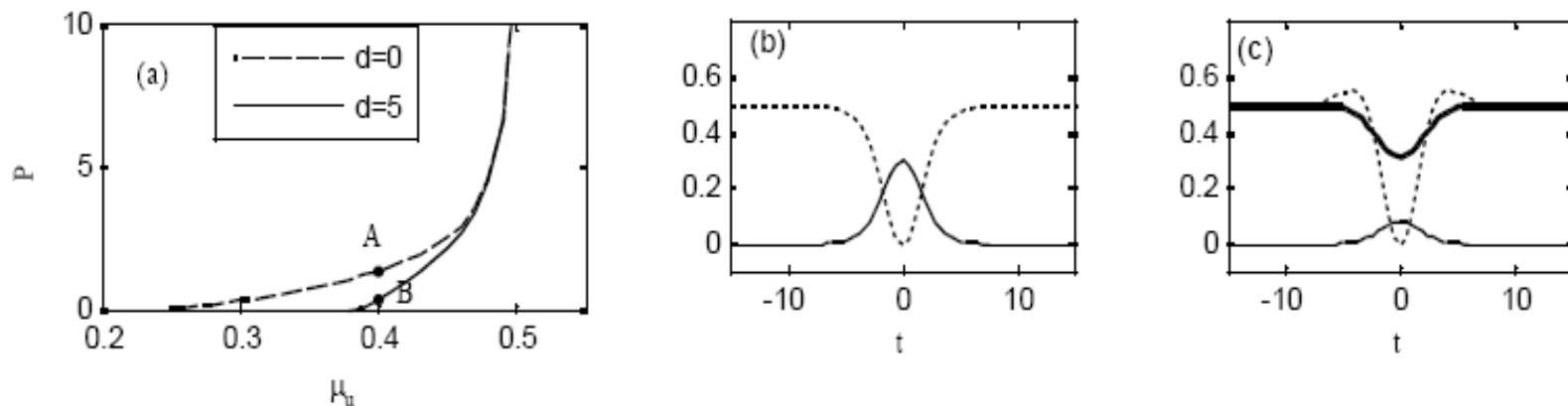


Nonlocal vector dark-bright soliton pairs, Spring 2006

- Nonlocal response of materials can stabilize solitons due to the diffusion of the Kerr nonlinearity.
- In this situation, solitons also need to increase their formation power to compensate the diffusion effect in nonlocal materials.
- We show that in normal dispersive media with Kerr-type nonlocal nonlinearity, one can stabilize dark-bright vector soliton pairs as well as **reduce the forming threshold power for guided bright solitons**.

$$i \frac{\partial \Psi}{\partial z} - \frac{1}{2} \frac{\partial^2}{\partial t^2} \Psi + n(t, z) \Psi = 0, \quad \Psi = U, V,$$

$$n(t, z) = \int_{-\infty}^{\infty} dt' R(t - t') (|U|^2 + |V|^2), \quad \text{and} \quad R(t) = \frac{1}{2\sqrt{d}} \exp\left(-\frac{|t|}{\sqrt{d}}\right),$$

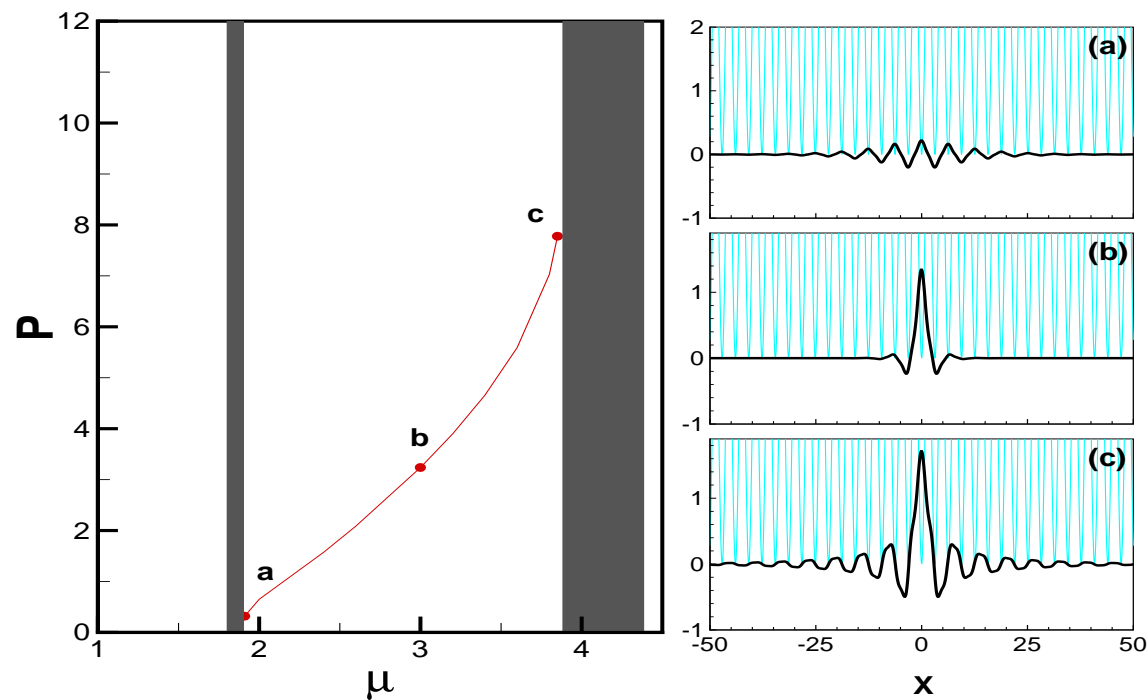


Gap solitons in optical lattices

1-D Gross-Pitaevskii equation with periodic potentials, $V(x) = V_0 \sin^2(k_0 x)$,

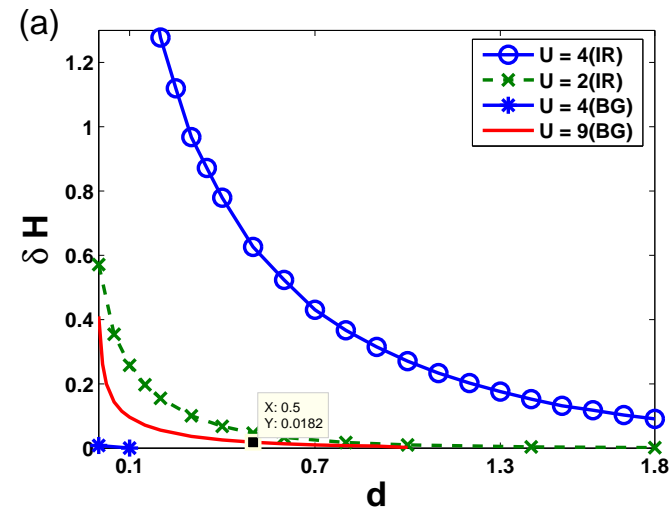
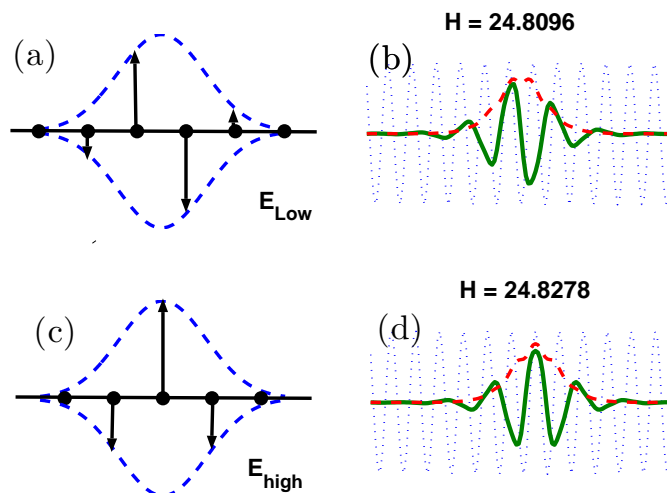
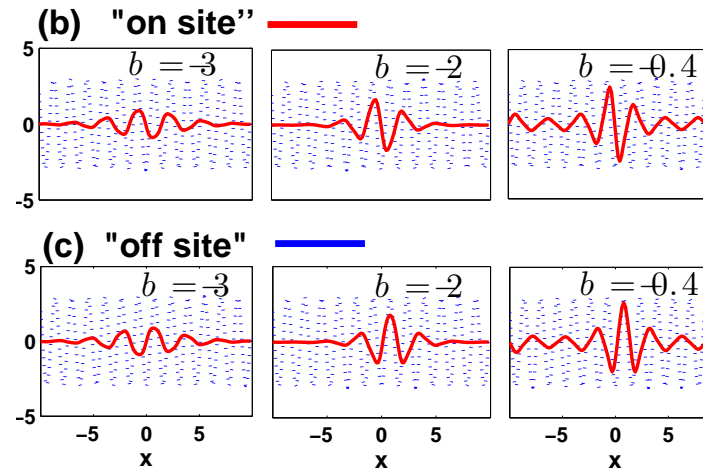
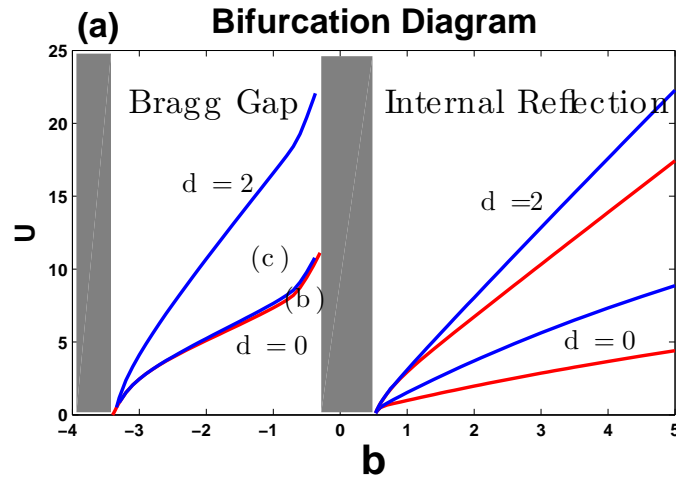
$$i\hbar \frac{\partial}{\partial t} \Phi_0(t, x) = -\frac{1}{2} \frac{\partial^2}{\partial x^2} \Phi_0(t, x) + V(x) \Phi_0(t, x) + g_{1D} |\phi_0(t, x)|^2 \phi_0(t, x)$$

which has gap soliton solutions.



Nonlocal gap solitons in optical lattices, Spring 2006

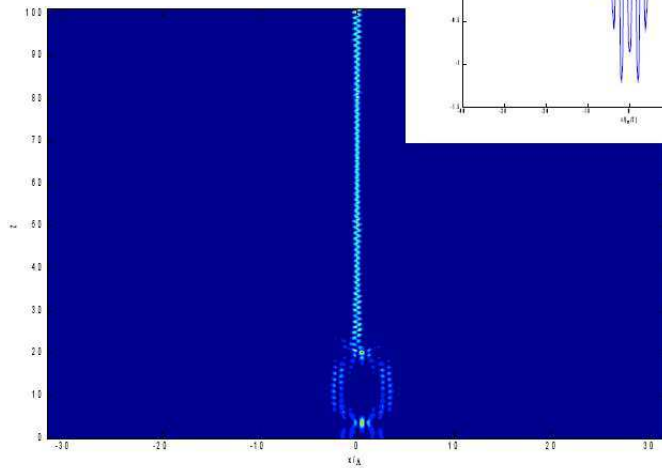
$$i \frac{\partial U}{\partial \xi} + \frac{1}{2} \frac{\partial^2 U}{\partial \eta^2} + V(\eta)U + n(\xi, \eta)U = 0, \quad n - d \frac{\partial^2 n}{\partial \eta^2} = |U|^2,$$



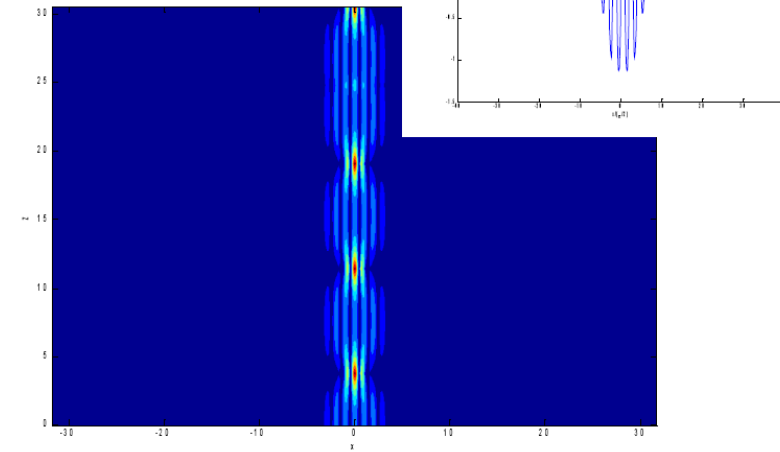
Collisions of Bragg gap soliton within nonlocal lattices

Collisions between two gap solitons in Bragg gap, with/without non-locality.

Local gap soliton,
 $U=3$
Initial separation = 2π



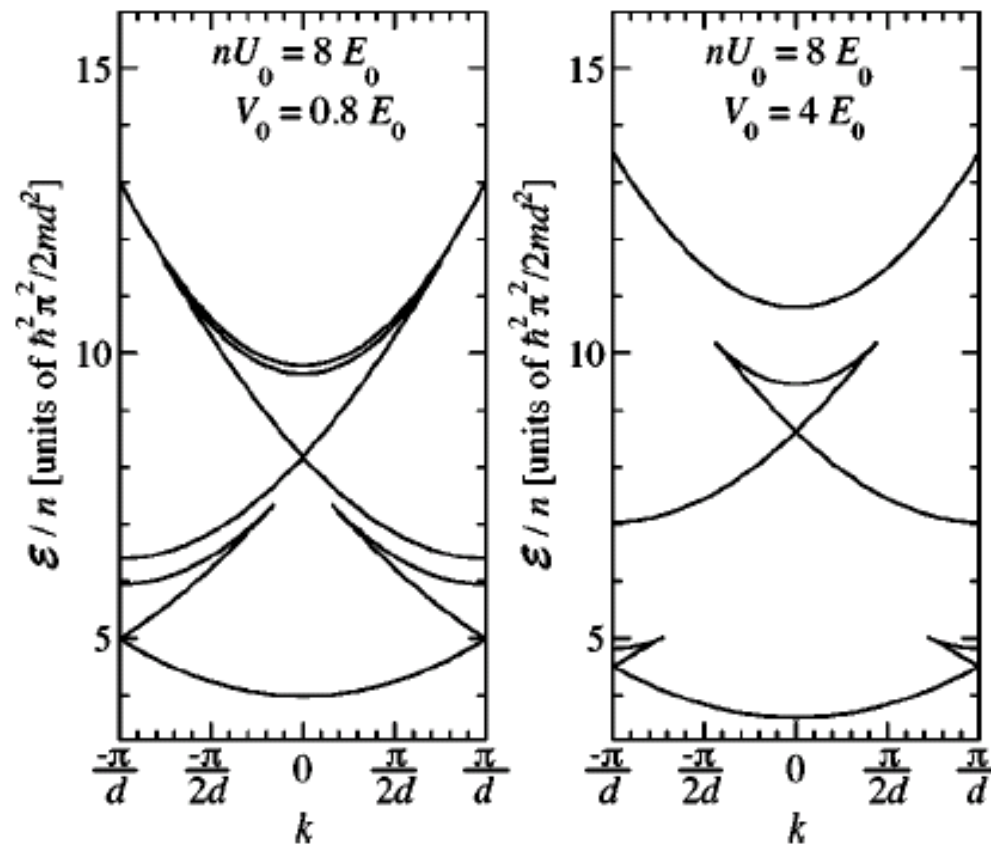
Non-Local gap soliton,
 $U=3, d=0.5$
Initial separation = 2π



- ➔ Only with a small degree of non-locality, solitons in Bragg gaps can be stabilized as well as movable.
- ➔ In such situation, the collision between two solitons behaves like the case without optical lattices.

Nonlinear band structure for BEC in OL

$$\frac{1}{2} \frac{d^2 \phi}{dx^2} - [V_0 \sin^2(Kx)\phi - \mu\phi] - \sigma |\phi|^2 \phi = 0.$$

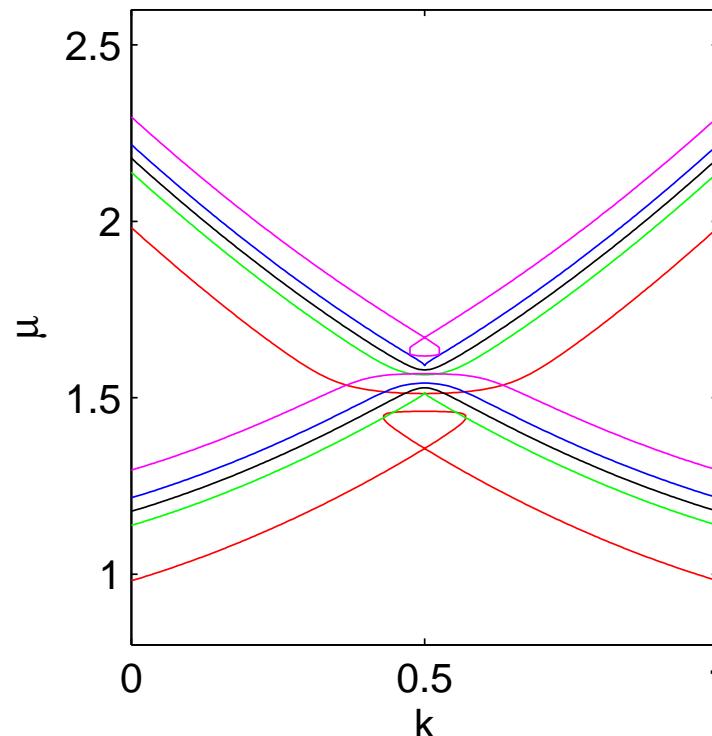


Band structure with dipolar BEC, Spring 2006

$$i \frac{\partial u}{\partial t} = \frac{-1}{2} \frac{\partial^2 u}{\partial x^2} + c|u|^2 u + d^2 n(x, t) u + \nu \cos(x) u,$$

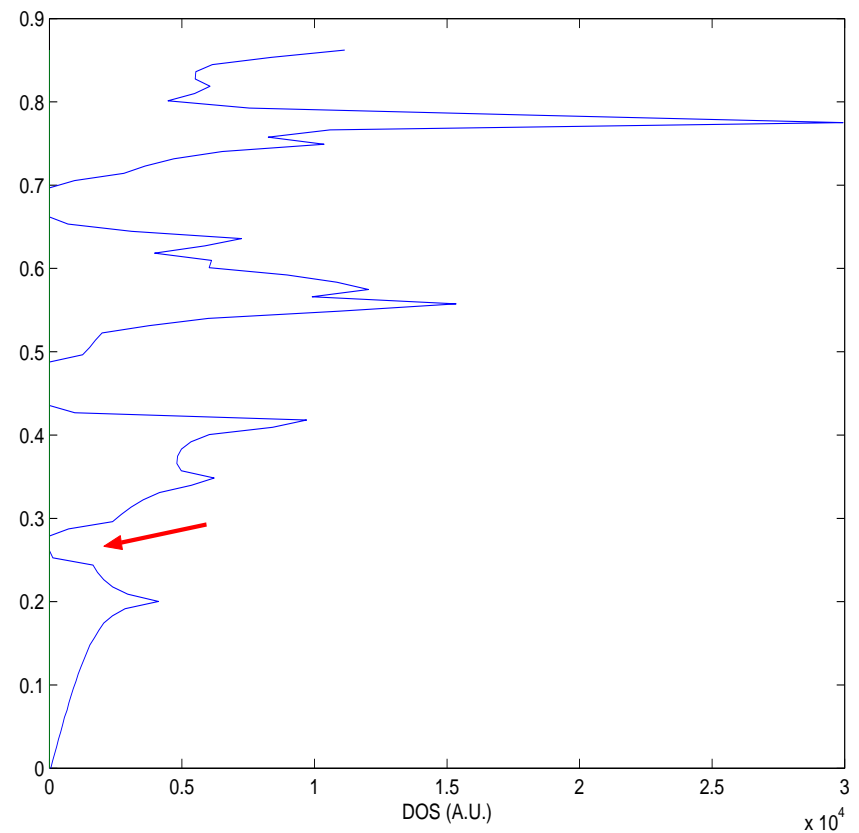
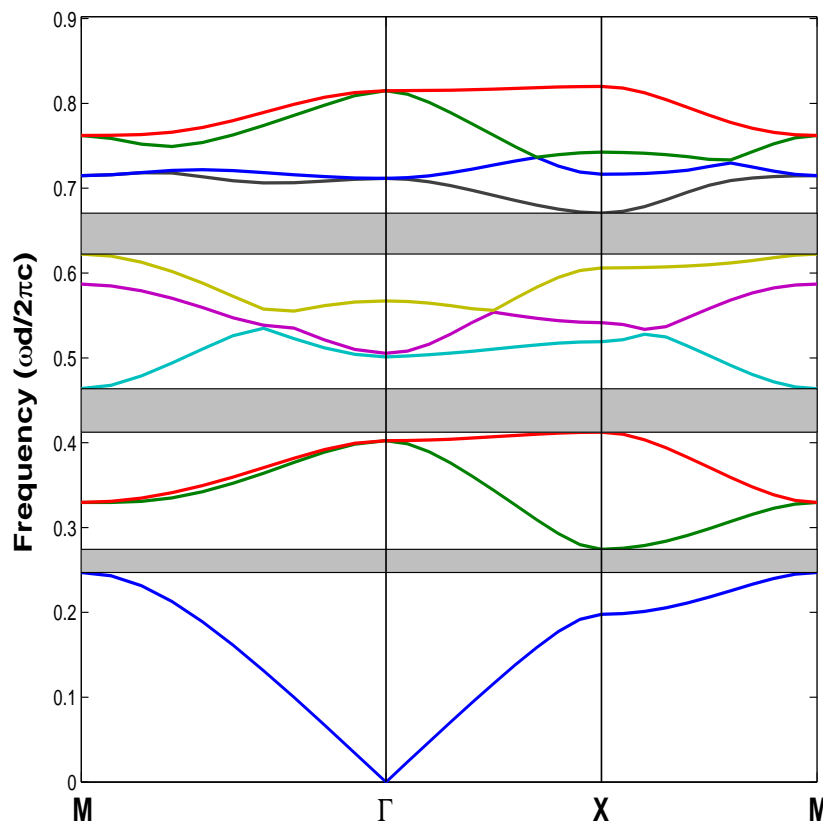
$$n(x, t) = \int_b^\infty d\eta V(x - \eta) |u(\eta)|^2, \quad V(x) = \frac{1 - 3 * \cos^2(r)}{|x|^3}, \text{ for } x > b,$$

where ν : lattice strength, c : mean field effect, d : dipole moment, b : scattering length.

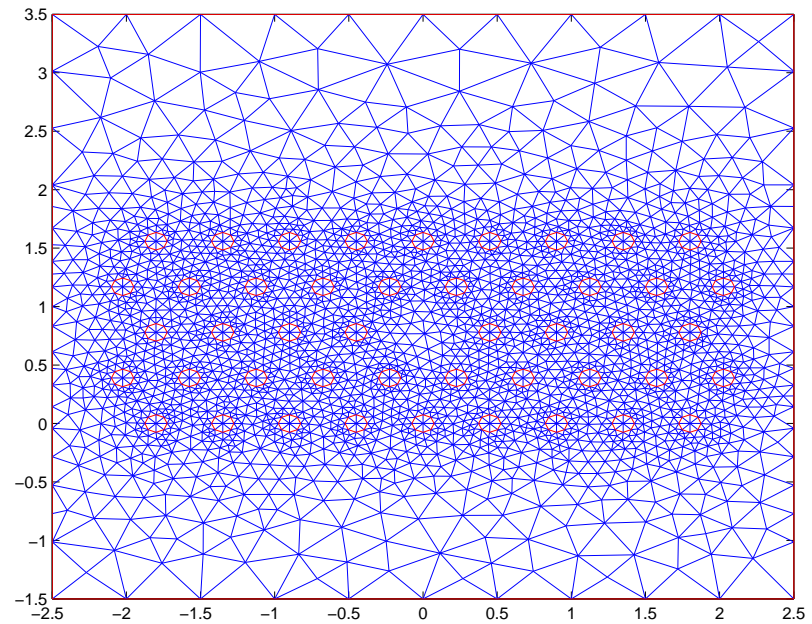
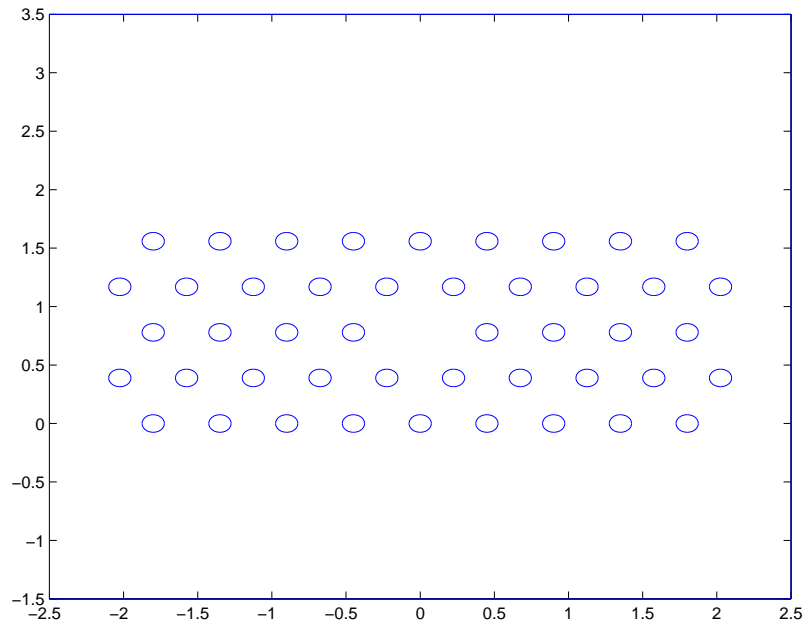


Band diagram and Density of States, Spring 2006

$$\frac{1}{\epsilon(\mathbf{r})} \nabla \times \{ \nabla \times \mathbf{E}(\mathbf{r}) \} = \frac{\omega^2}{c^2} \mathbf{E}(\mathbf{r}),$$



Photonic Crystals



Course Project, Spring 2006

Thermal analysis of Nd:YVO₄ laser crystal in diode pump solid state laser by finite element method

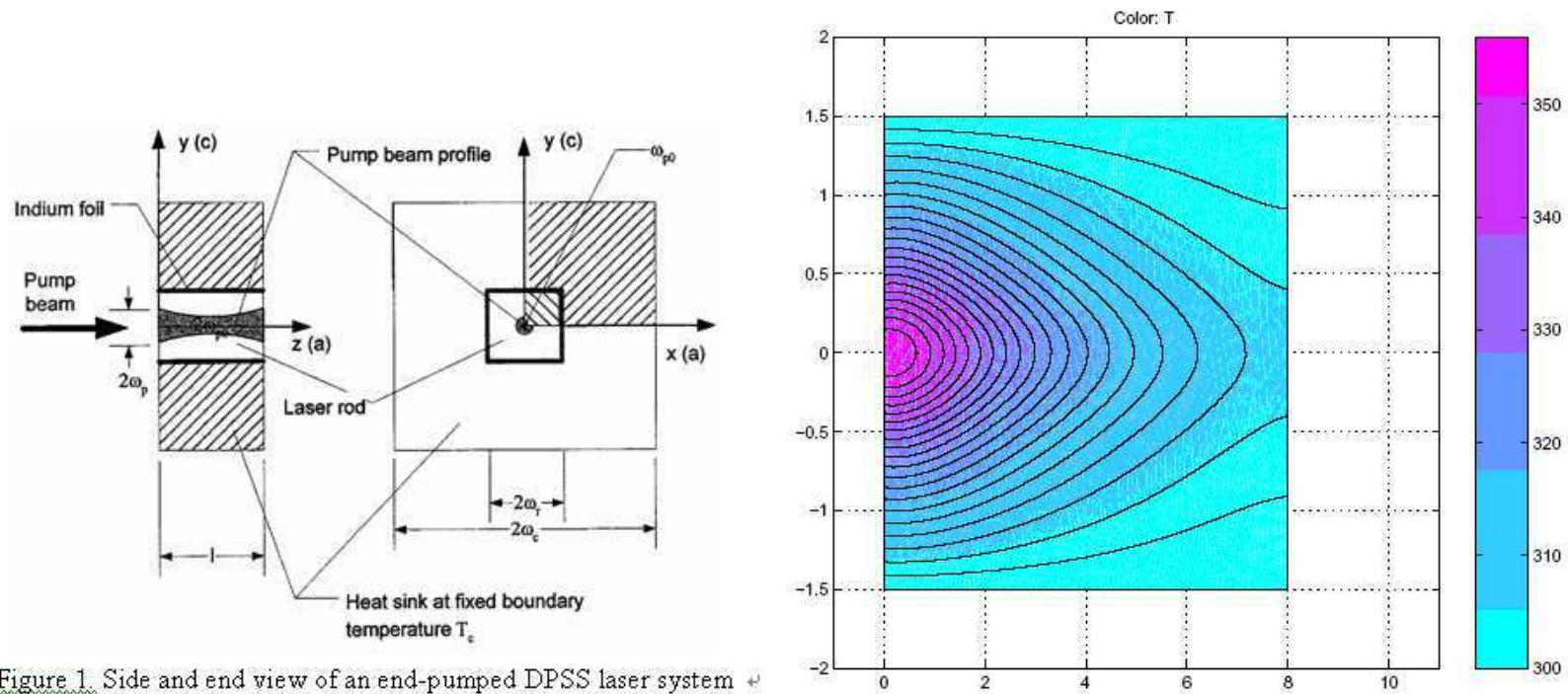


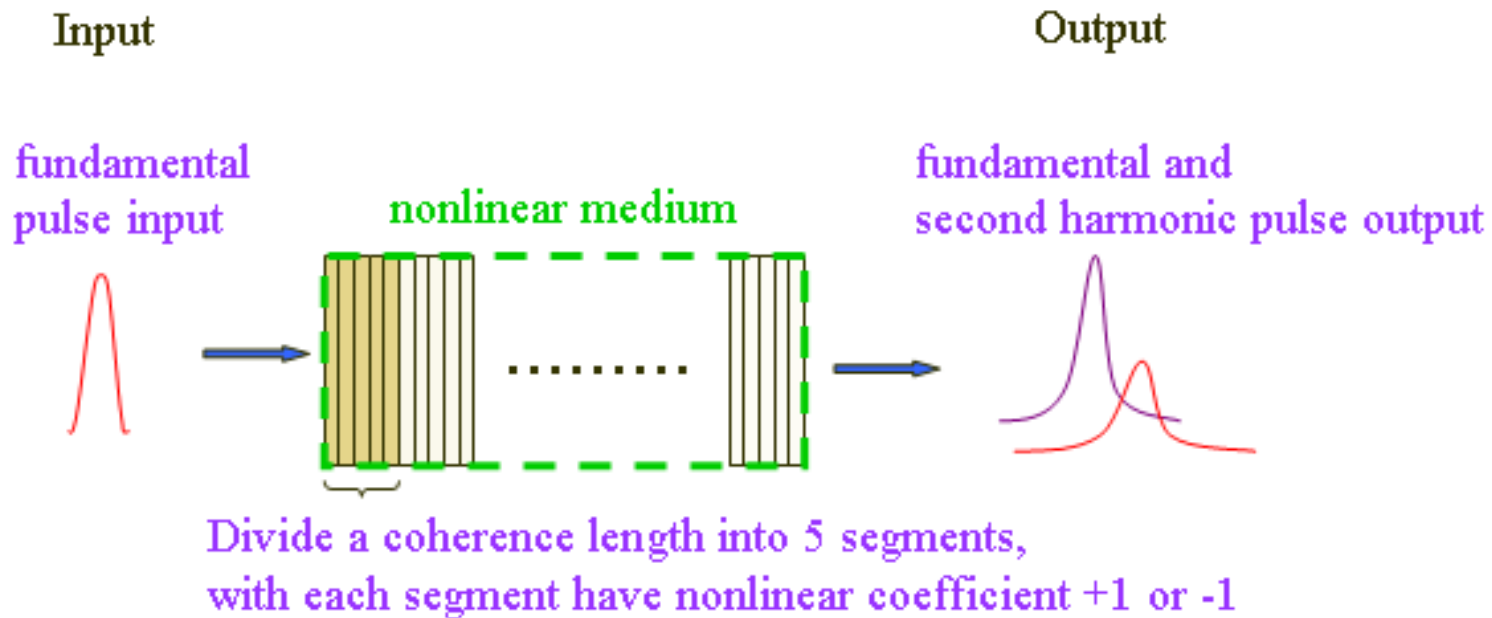
Figure 1. Side and end view of an end-pumped DPSS laser system +

Shoutai Lin, 938106

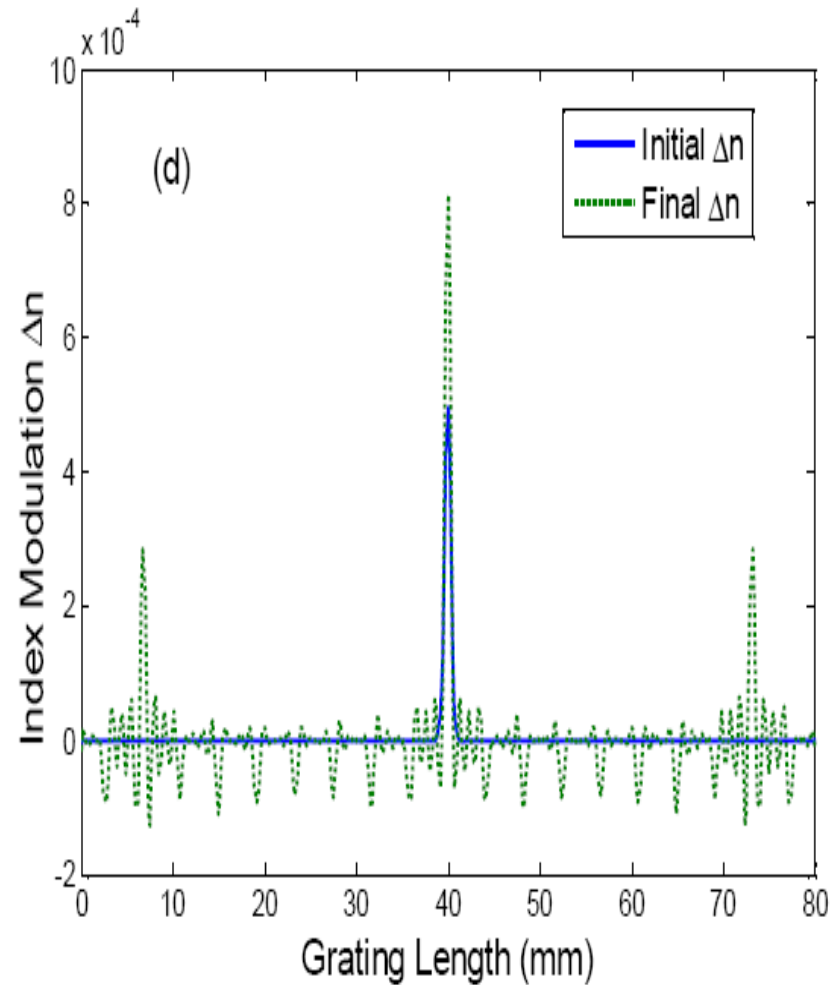
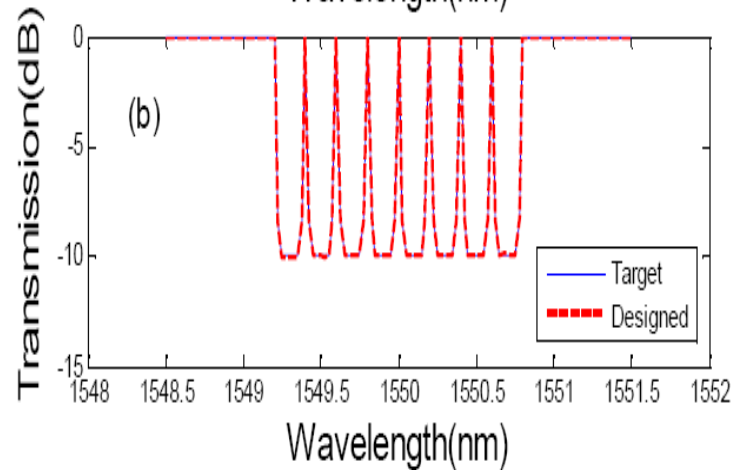
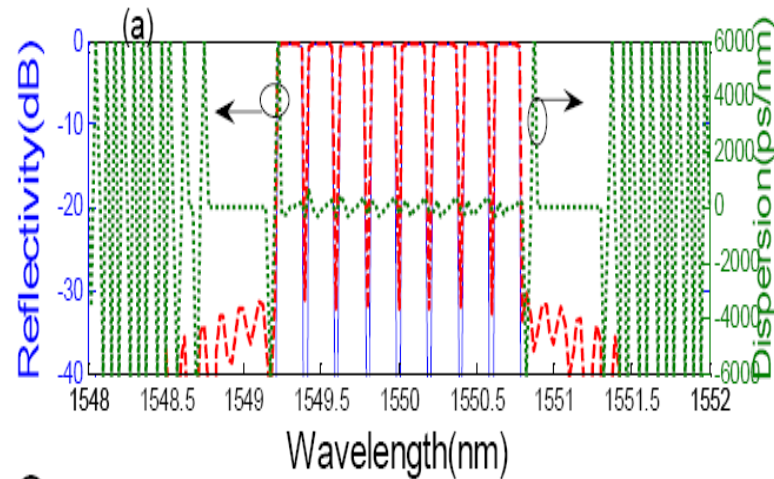
Optimization of SHG pulse

$$\frac{\partial A}{\partial z} = \frac{\eta}{2} \frac{\partial A}{\partial T} + i\xi_1 \frac{\partial^2 A}{\partial T^2} - i\rho_1 A^* B,$$

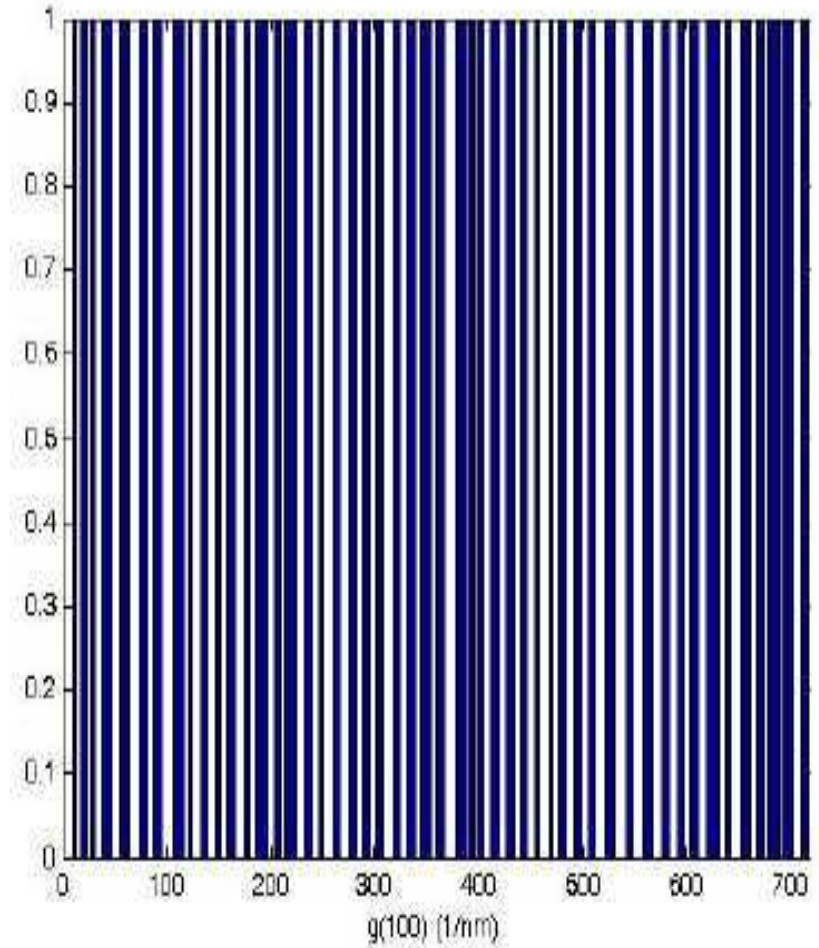
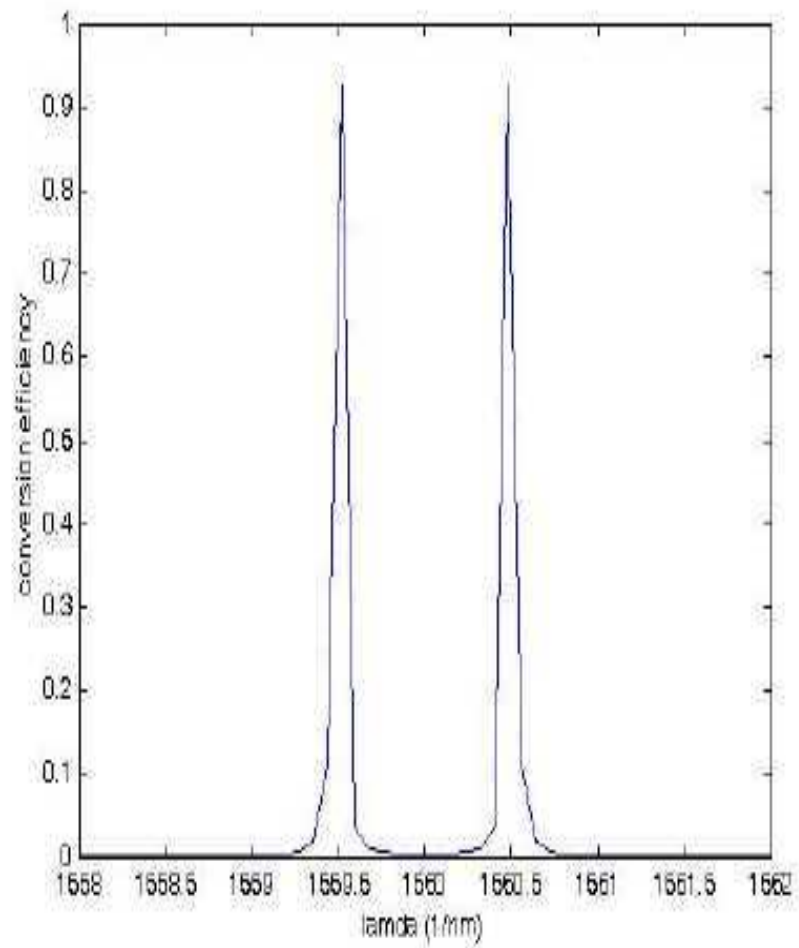
$$\frac{\partial B}{\partial z} = -\frac{\eta}{2} \frac{\partial B}{\partial T} + i\xi_2 \frac{\partial^2 A}{\partial T^2} - i\Delta k B - i\rho_1 A^2,$$



Lagrange Multiplier method for FBG



Course Project, Spring 2006



Ching-Jen Cheng, g936812,

Feedback of VCSEL, Spring 2007

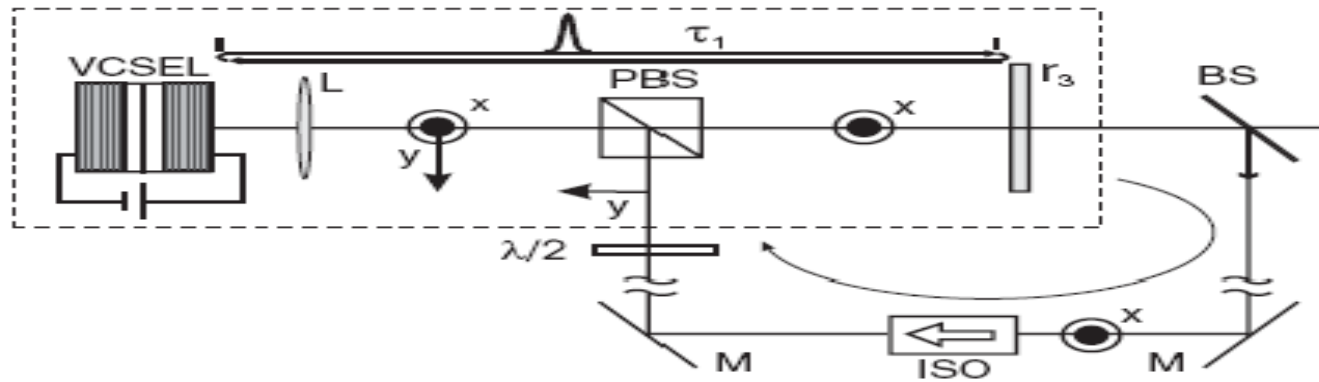
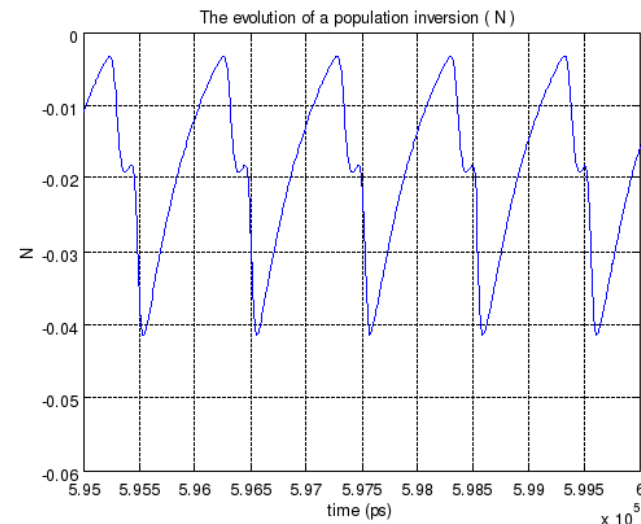
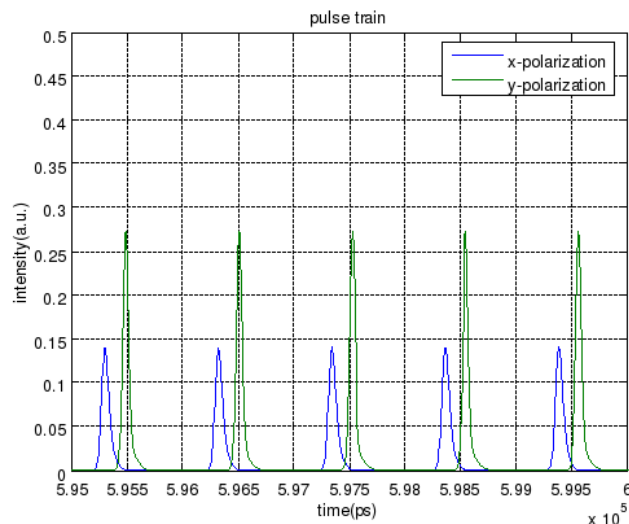
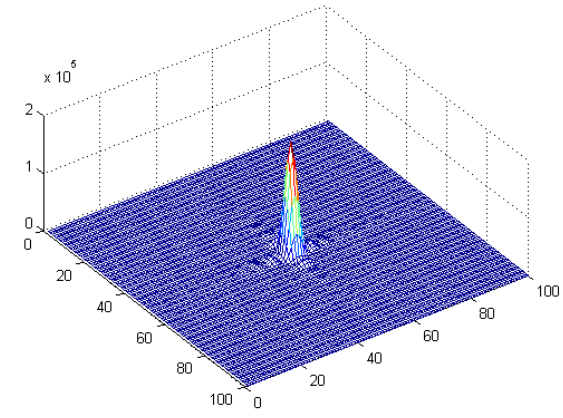
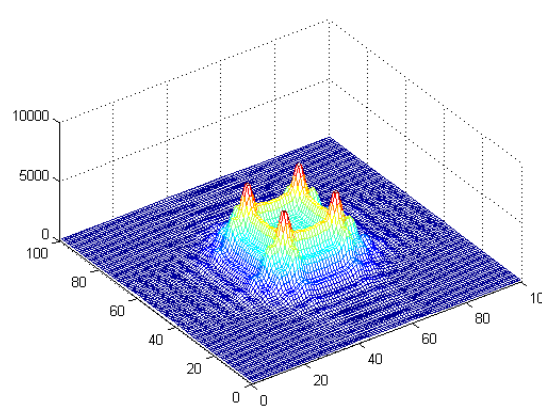
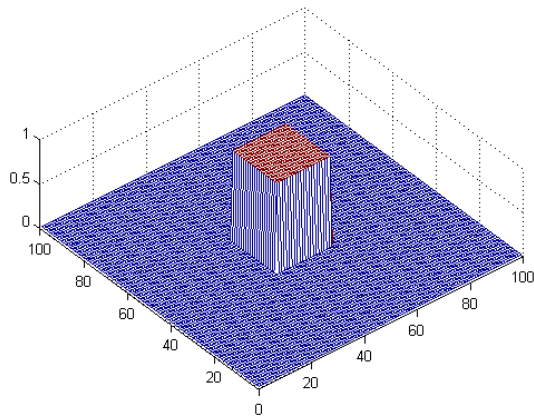
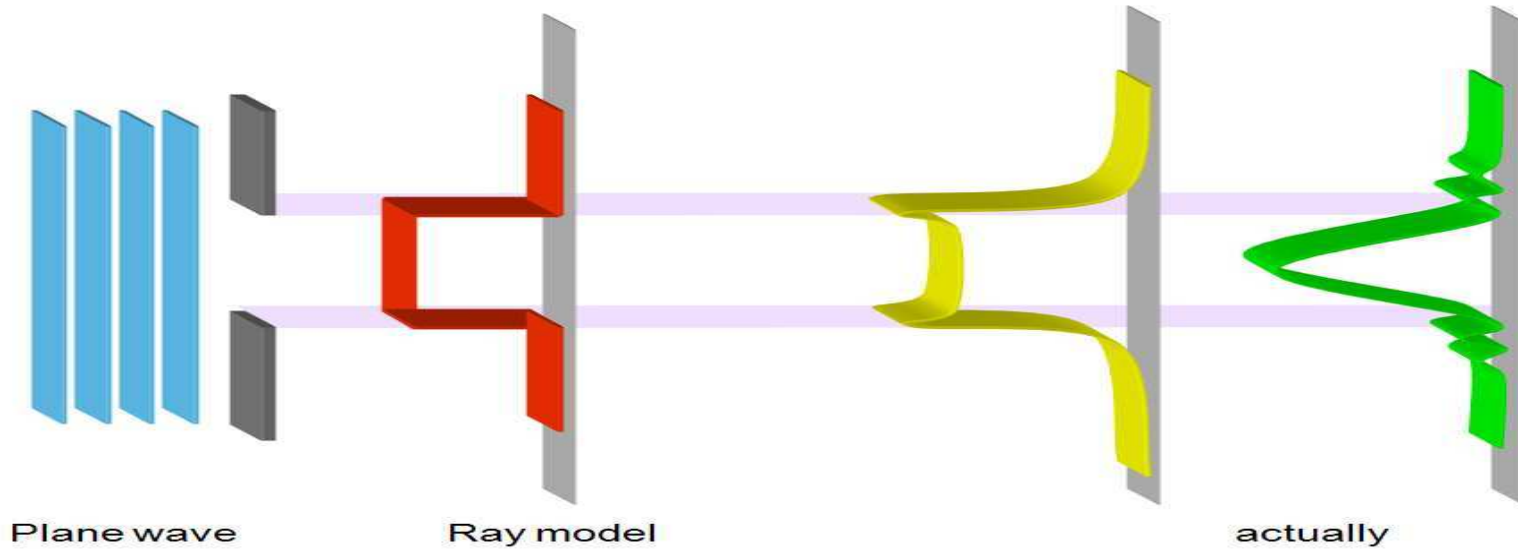


FIG. 1. Mode-locked laser setup comprising an external-cavity VCSEL (dashed box) and a reinjection arm. Symbols: polarizing beam splitter (PBS), optical isolator (ISO), and half-wavelength plate ($\lambda/2$).

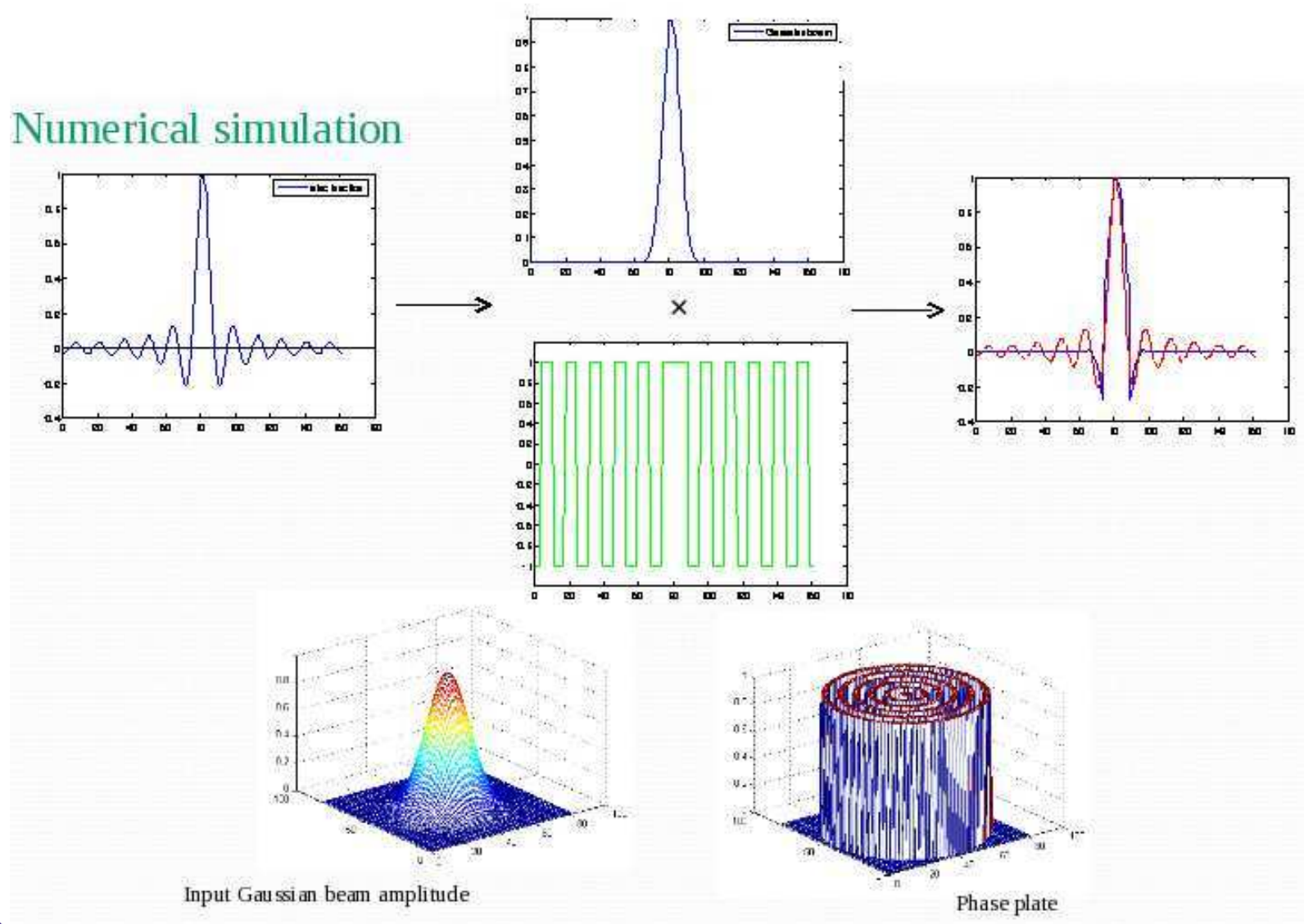


Diffraction Optics, Spring 2007



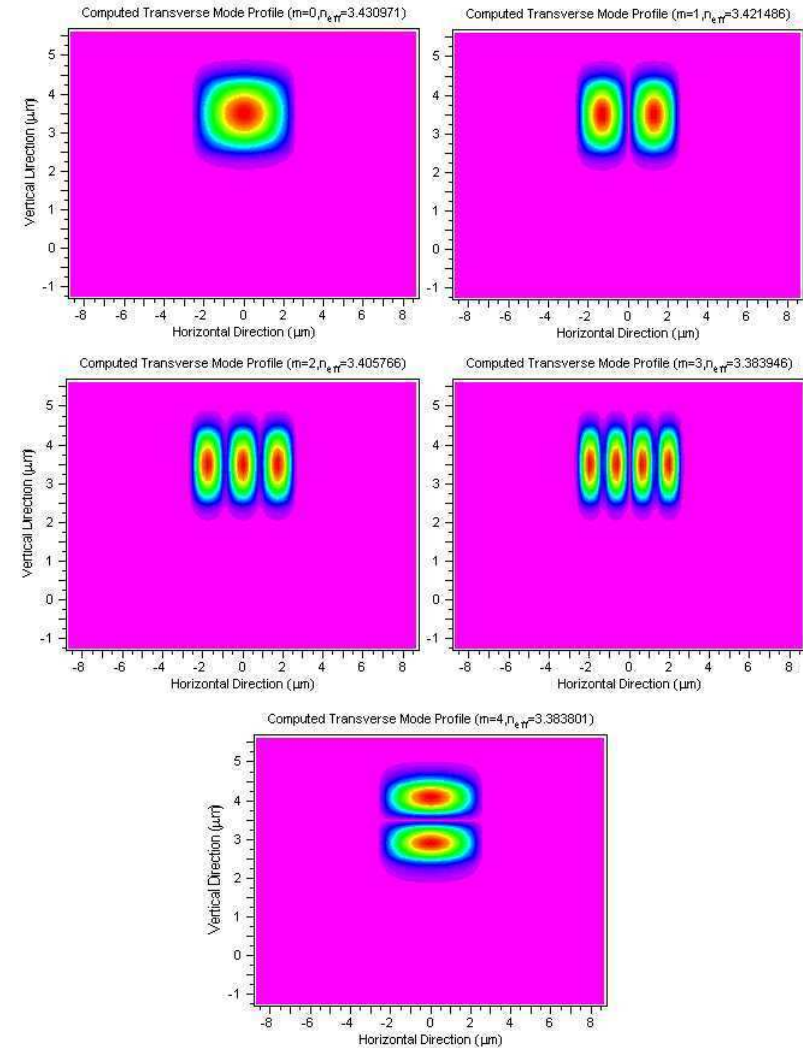
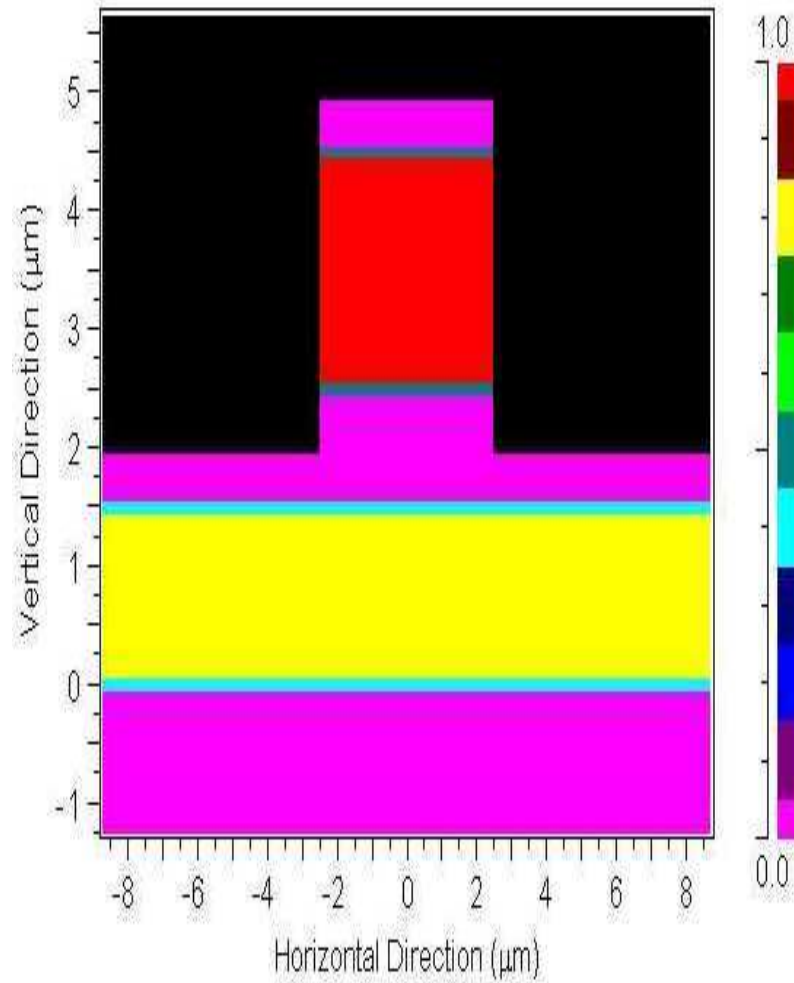
Diffraction Optics, Spring 2007

Numerical simulation



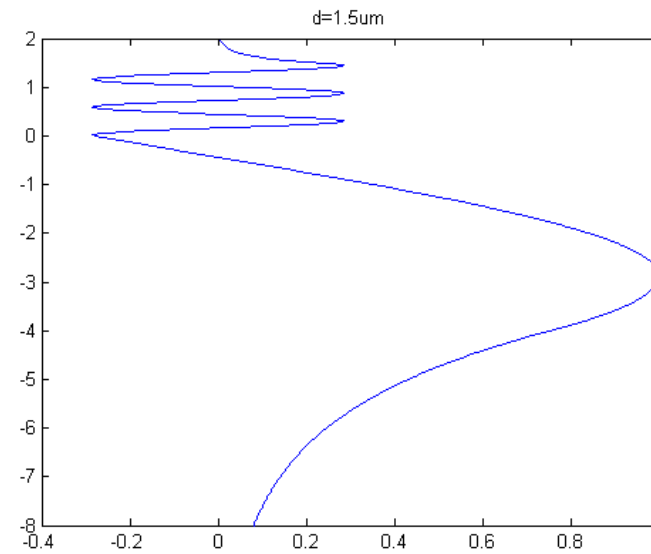
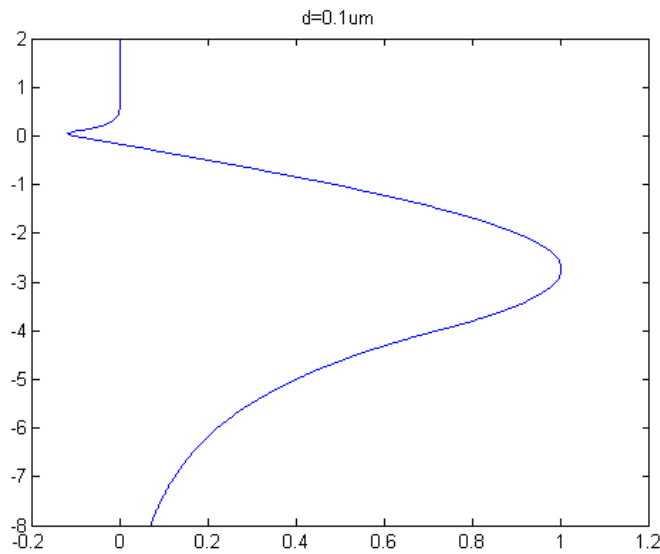
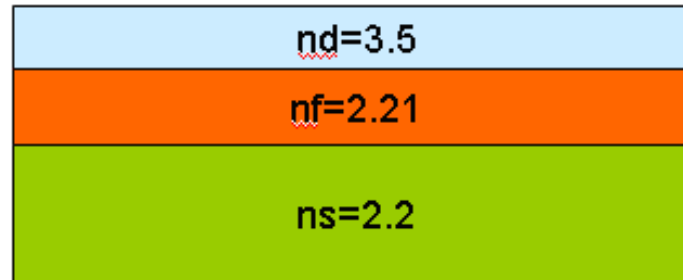
Waveguide structures

Contour Map of Transverse Index Profile at Z=0

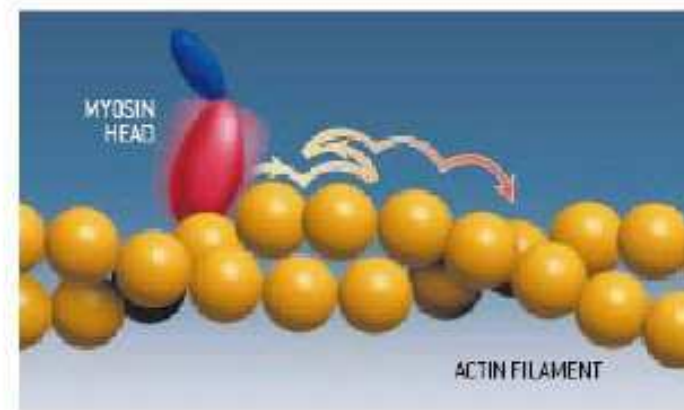
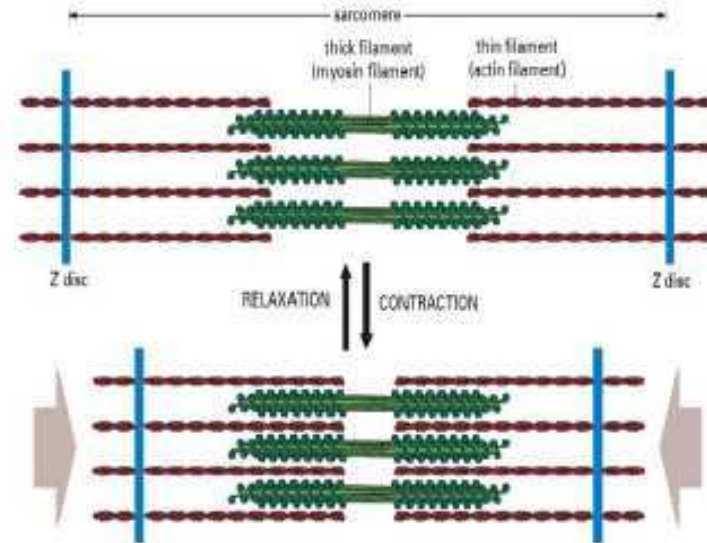
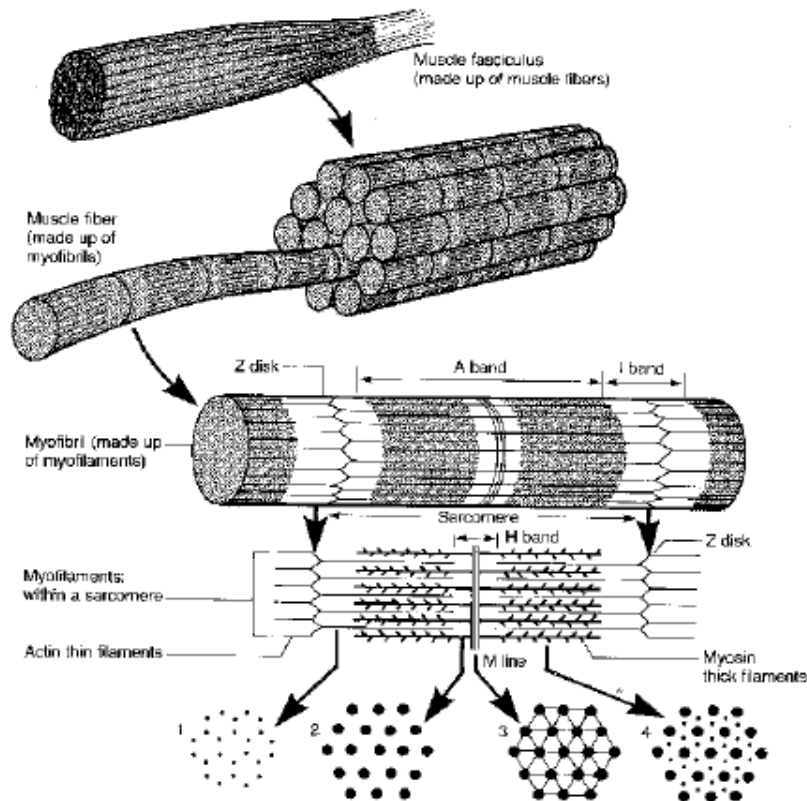


Course Project, Spring 2006

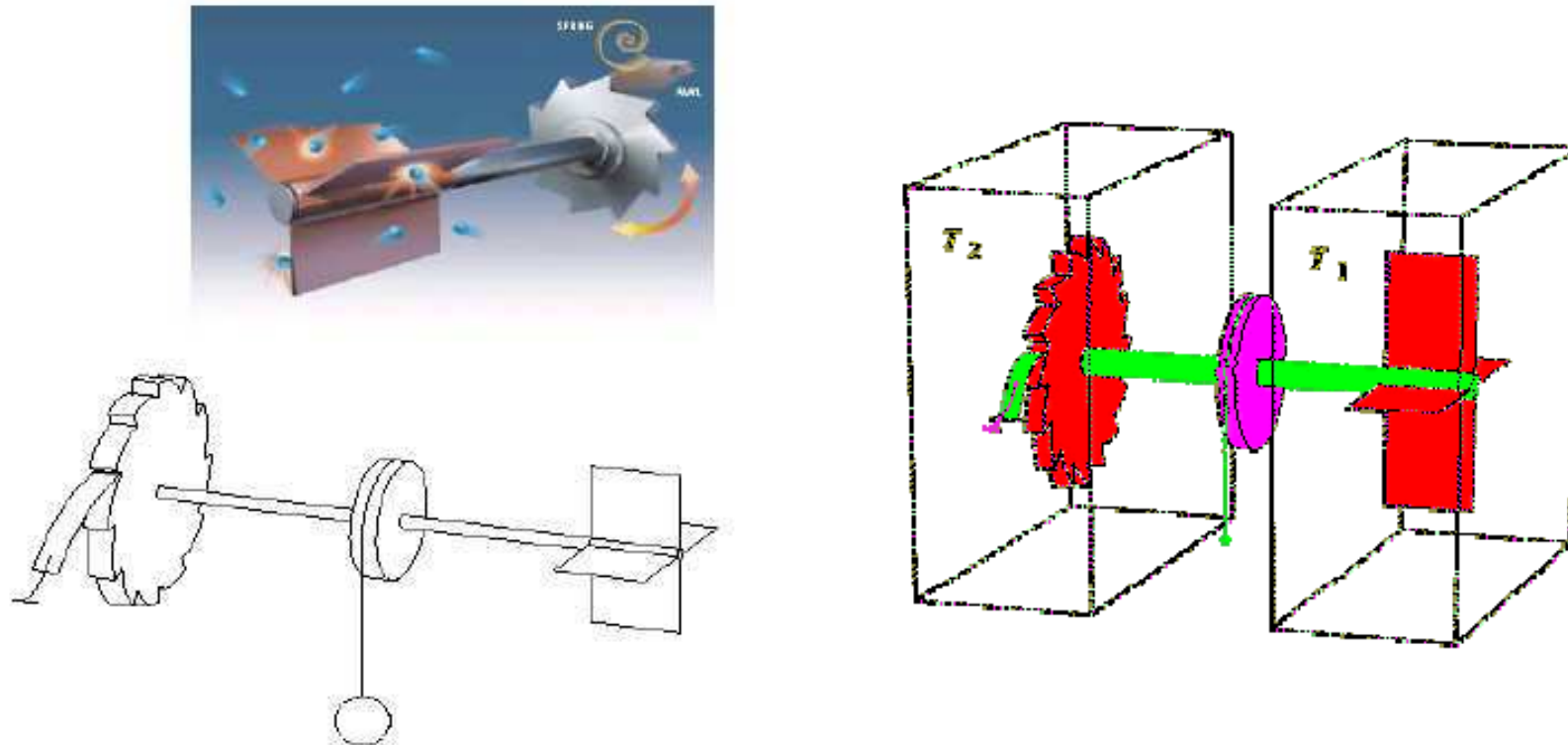
$na=1$



Muscle Contraction: Myosin



Ratchets



thought experiment of a perpetual mobile against the 2nd Law of Thermodynamics

R. P. Feynman, *The Feynman Lectures on Physics*, Vol. 1, Chap. 46 (1963).

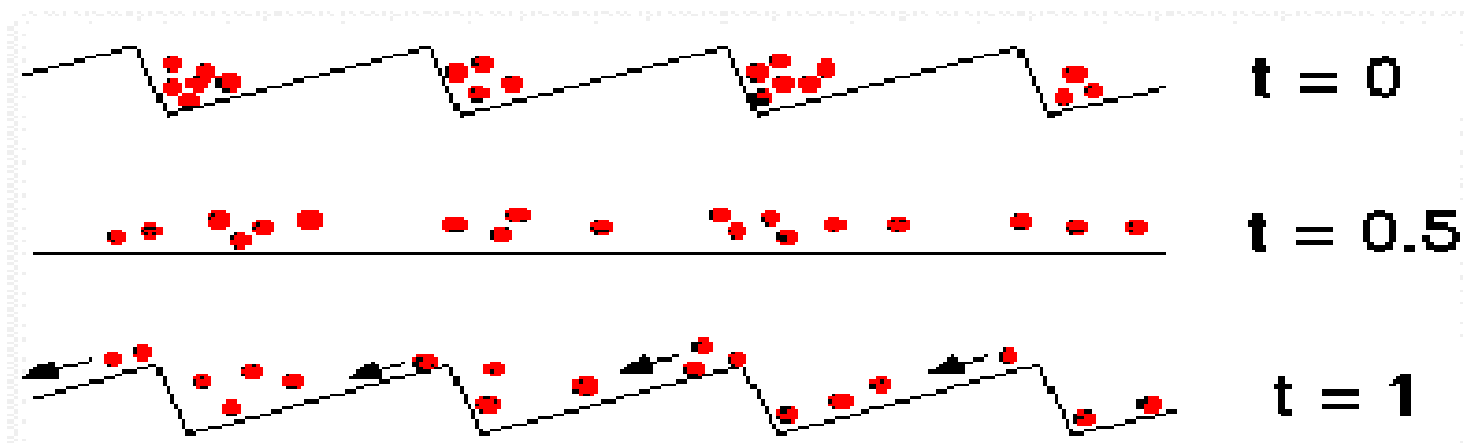
Mathematical Model for Ratchets

Ratchet ingredients:

1. Brownian particle (mass m)
2. periodic asymmetric potential $V(x, t)$
3. zero mean driving forces $f(t)$, i.e. $\langle f(t) \rangle = 0$

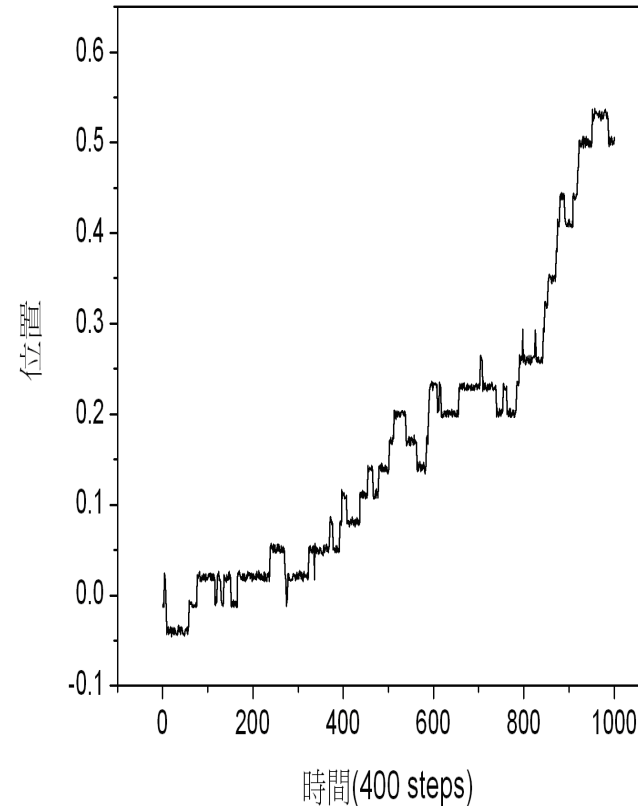
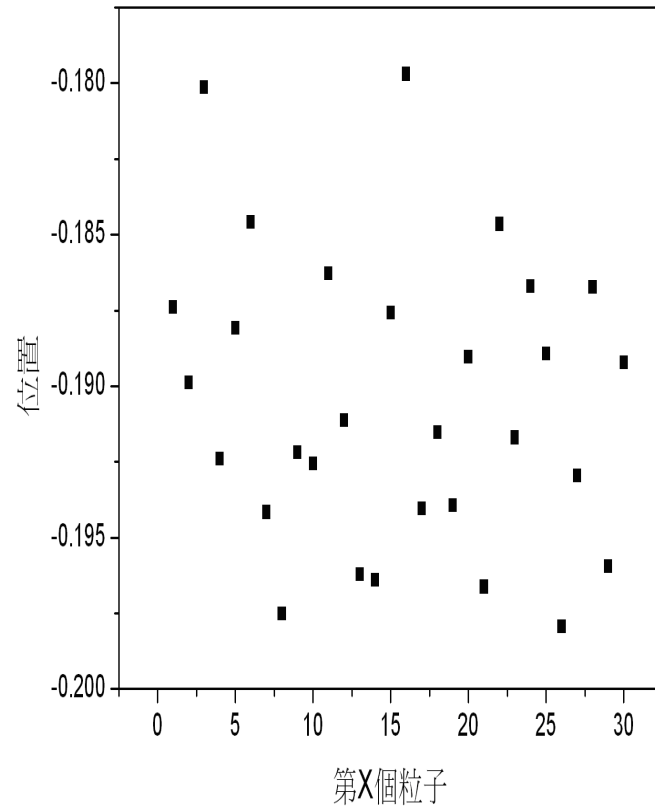
Ratchet model: $m\ddot{x} + \gamma\dot{x} + \frac{d}{dx}V(x, t) = f(t)$ **Interesting**

Behavior: $\langle x(t) \rangle \neq 0$ even when $\langle f(t) \rangle = 0$



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Brownian motor,



Course Project, Spring 2006

Smith-Purcell radiation,

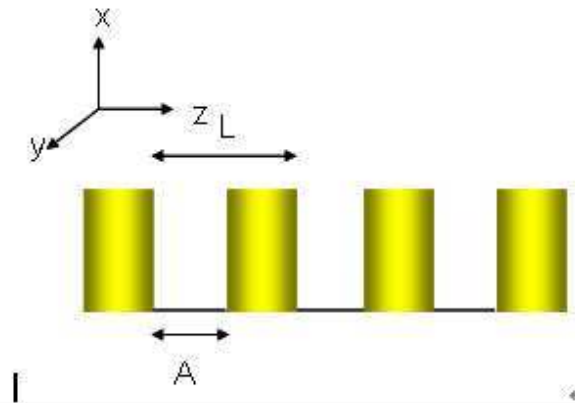


Fig. 1

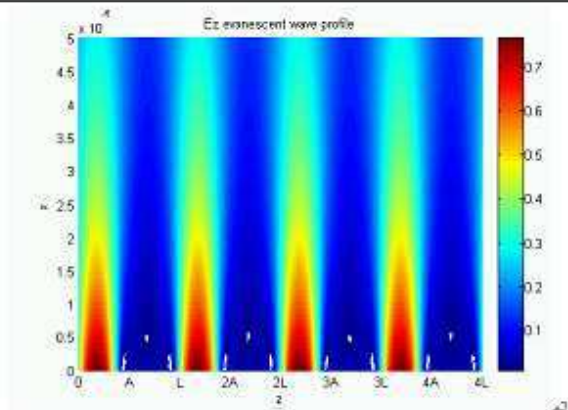


Fig. 2

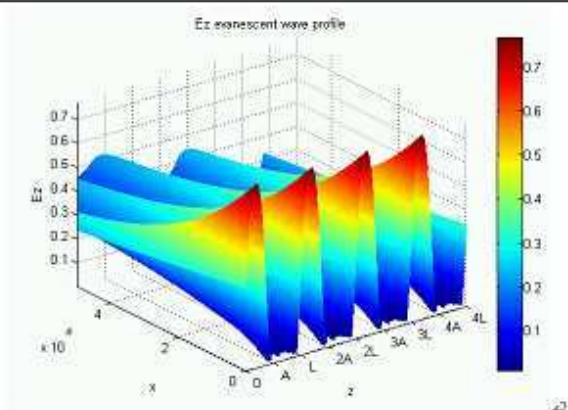
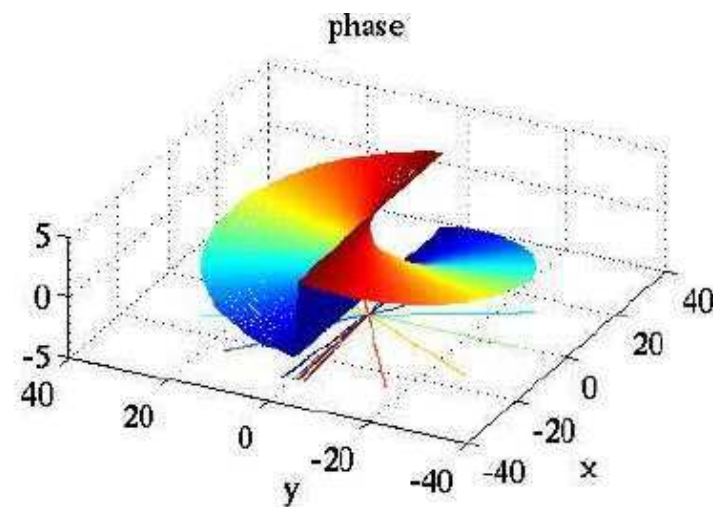
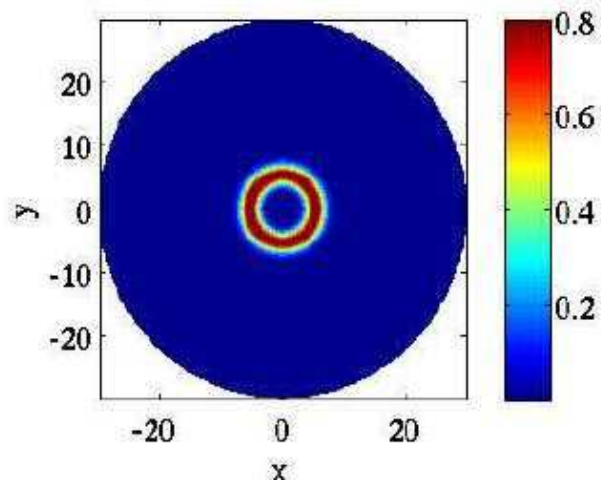
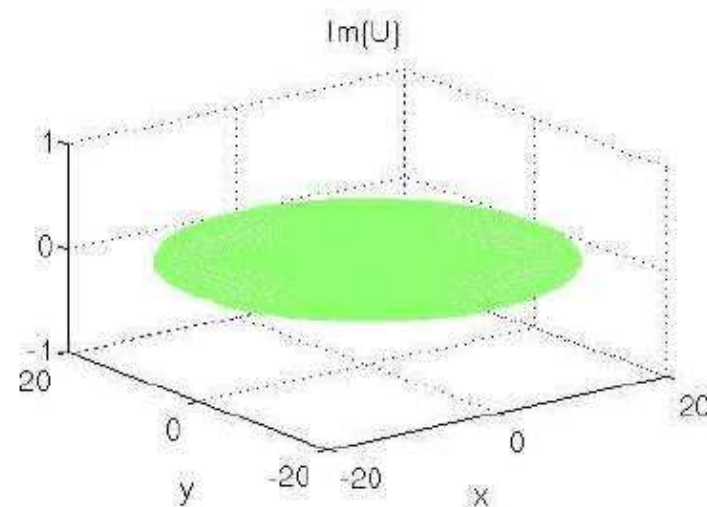
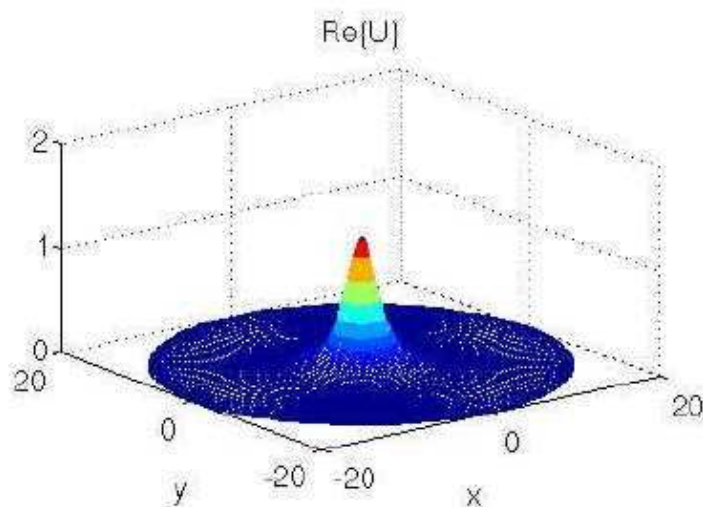
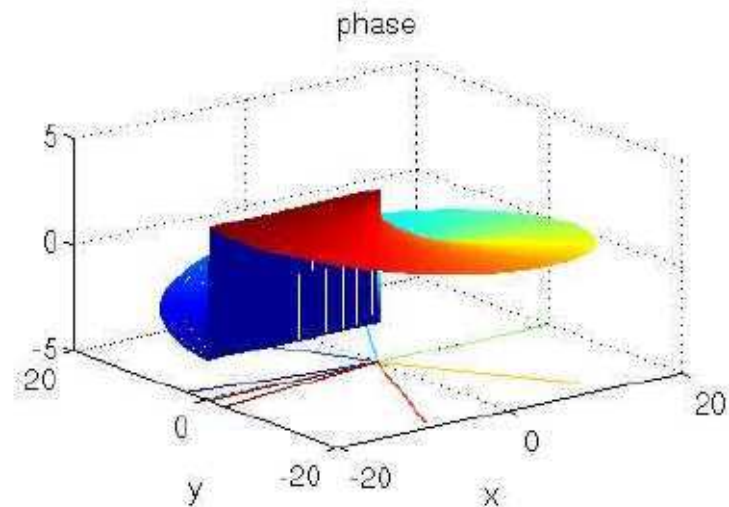
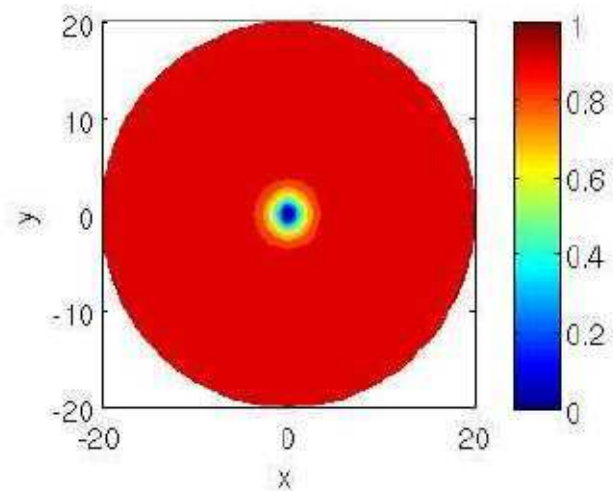
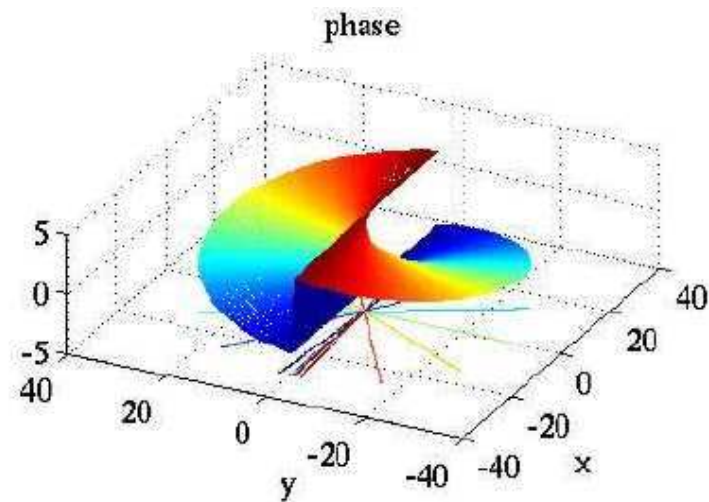
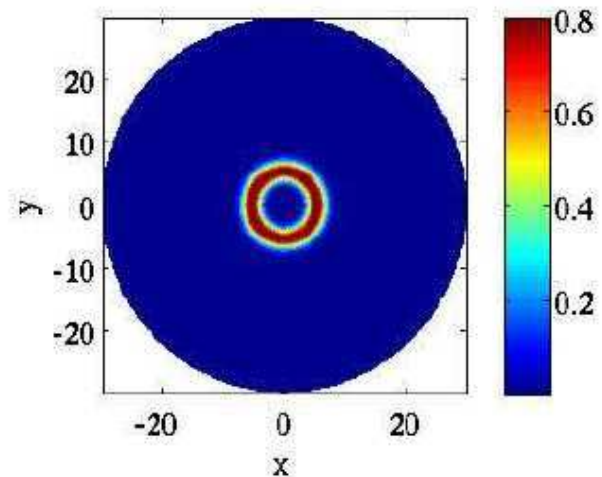


Fig. 3

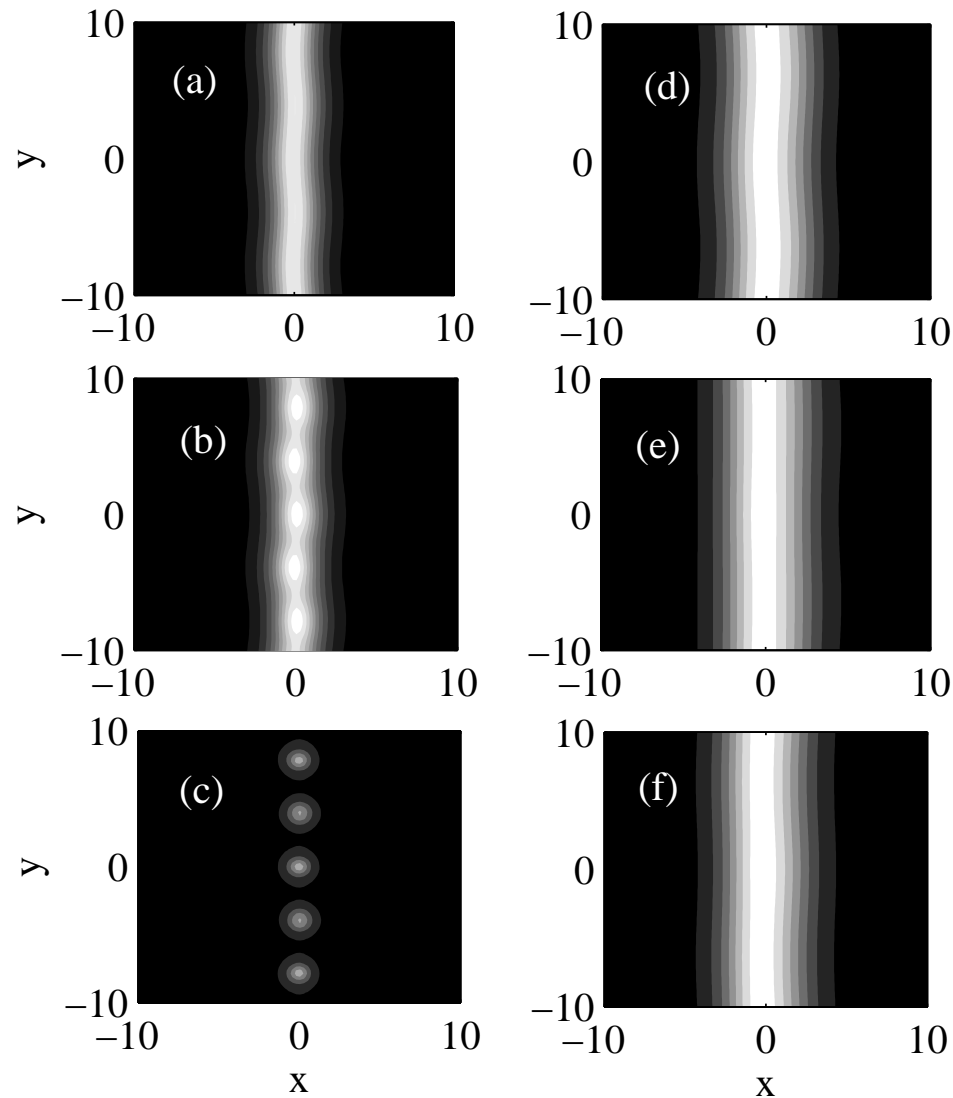
Vortex and Vortex solitons



Vortex and Vortex solitons



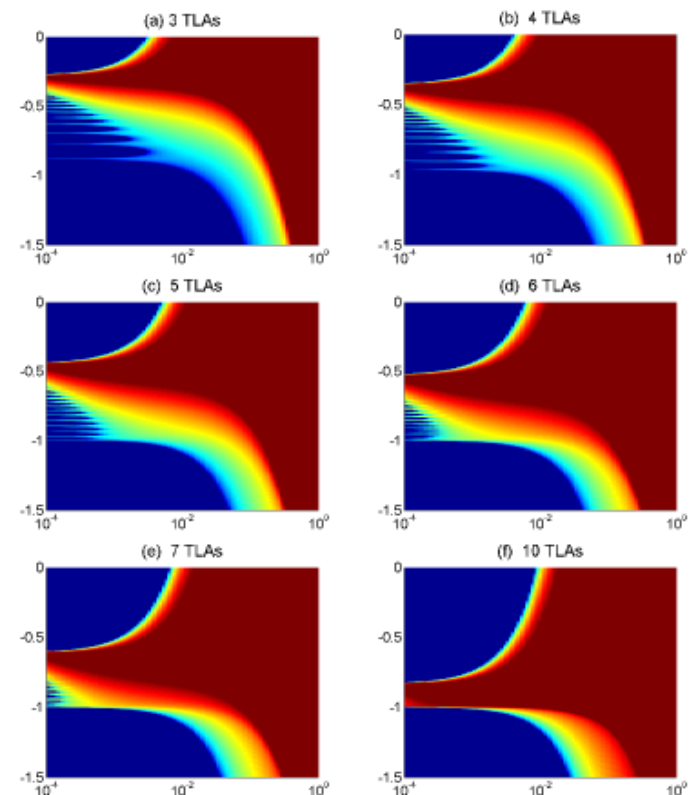
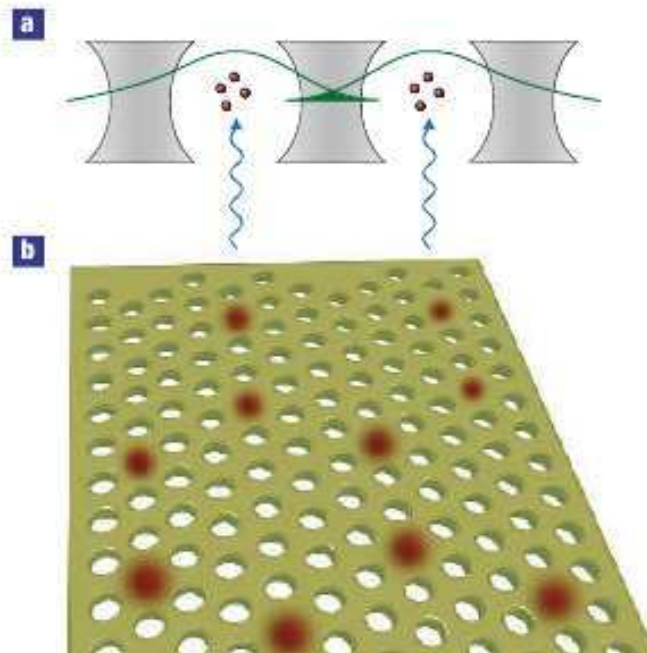
Soliton transverse instabilities in nonlocal nonlinear media



Quantum Phase Transitions of Light in the Dicke-Bose-Hubbard model

$$\hat{H} = \sum_i H_i^{\text{DM}} - \kappa \sum_{ij} a_i^\dagger a_j - \mu \sum_i N_i,$$

$$H_i^{\text{DM}} = \varepsilon J_i^+ J_i^- + \omega a_i^\dagger a_i + \beta(a_i J_i^+ + a_i^\dagger J_i^-),$$



Soi-Chan Lai and Ray-Kuang Lee, *arXiv:0709.1352* (quant-ph) (2007),
to *Phys. Rev. A*.